

Statistical Physics and Thermodynamics (SS 2017)

Problem sheet 8

Hand in: Friday, June 16 during the lecture

<http://www.physik.fu-berlin.de/en/einrichtungen/ag/ag-netz/lehre/>

1 Thermodynamic potentials and Gibbs-Duhem relation (11 points)

The thermodynamic properties of a system are described by a thermodynamic potential. Which potential one uses depends on the physical situation. For example, for a system with constant particle number N and volume V coupled to a heat reservoir with temperature T , the Helmholtz free energy $F(T, V, N)$ is used.

We now allow for particle exchange with a reservoir (with chemical potential μ). To derive the corresponding thermodynamic potential, one performs a Legendre transformation on $F(T, V, N)$ to eliminate N in favor of the chemical potential μ . This is done as follows:

1. One calculates

$$\mu(T, V, N) = \left(\frac{\partial F(T, V, N)}{\partial N} \right)_{T, V}, \quad (1)$$

and solves this equation for N to obtain $N(T, V, \mu)$.

2. One uses the result from step one to calculate the Legendre transform

$$\Omega(T, V, \mu) = F(T, V, N(T, V, \mu)) - \mu N(T, V, \mu). \quad (2)$$

a) In the lecture, we showed that the Helmholtz free energy of an ideal gas is given by

$$F(T, V, N) = Nk_B T \left[\ln \left(\lambda^3 \frac{N}{V} \right) - 1 \right], \quad (3)$$

with the thermal wavelength $\lambda = h/\sqrt{2\pi m k_B T}$, where h is the Planck constant and m the mass of a gas particle. Explicitly perform the Legendre transformation for the ideal gas to show that you recover the grand canonical potential,

$$\Omega(T, V, \mu) = -k_B T \frac{V}{\lambda^3} \exp \left(\frac{\mu}{k_B T} \right), \quad (4)$$

which was derived in the lecture using the grand canonical partition function. **(3 points)**

b) Show that in general (i.e. for a general Helmholtz free energy, and not just for the ideal gas considered in task a)),

$$d\Omega = -SdT - pdV - Nd\mu. \quad (5)$$

(3 points)

Hint: Start from eq. (2), where Ω is expressed as a function of (T, V, μ) . Then use the definition of the total differential, $d\Omega = (\partial\Omega/\partial T)_{V, \mu} dT + (\partial\Omega/\partial V)_{T, \mu} dV + (\partial\Omega/\partial \mu)_{T, V} d\mu$, the chain rule and the fact that because of $dF = -SdT - pdV + \mu dN$ you know the partial derivatives of $F(T, V, N)$, to derive eq. (5).

c) There cannot be a thermodynamic potential which only has the intensive variables (T, p, μ) as independent variables. This follows from the Gibbs-Duhem relation,

$$0 = SdT - Vdp + Nd\mu, \quad (6)$$

which states that the three intensive variables (T, p, μ) of a 1-component system are related. In the lecture, you derived this relation from the grand canonical potential. Derive this relation from the Gibbs free energy $G(T, p, N)$.

Proceed as follows:

1. In analogy to part b), use the definition of the Legendre transform to show that $dG = -SdT + Vdp + \mu dN$. **(2 points)**
2. Show that if a (continuously differentiable) function $f(x)$ is homogeneous, i.e. $f(\alpha x) = \alpha f(x)$ (for all $\alpha \in (0, \infty)$), then it fulfills the Euler relation¹

$$f(x) = x \frac{\partial f(x)}{\partial x}. \quad (7)$$

(1 point)

3. Use the Euler relation to show that $G = \mu N$ (which of the variables (T, p, N) corresponds to x from the Euler relation?), so that $dG = Nd\mu + \mu dN$. **(1 point)**
4. Put 1. and 3. together to obtain the Gibbs-Duhem relation. **(1 point)**

2 Law of mass action for hydrogen (9 points)

In this exercise you will derive the law of mass action for the reaction



Consider H-atoms (mass m) and H_2 -molecules (mass $2m$) in thermodynamic equilibrium in a volume V . Both molecule species can be treated as an ideal gas.

a) Since the individual particle numbers are not conserved, we work in the grand canonical ensemble. Write down the grand-canonical potential $\Omega_{\text{H}}(T, V, \mu_{\text{H}})$ of the H-atoms and the grand-canonical potential $\Omega_{\text{H}_2}(T, V, \mu_{\text{H}_2})$ of the H_2 -molecules. How are the thermal wavelengths λ_{H} , λ_{H_2} related? **(1 point)**

b) For given T , V , μ_{H} , μ_{H_2} , calculate the respective number of particles $\langle N_{\text{H}} \rangle \equiv N_{\text{H}}$, $\langle N_{\text{H}_2} \rangle \equiv N_{\text{H}_2}$. **(1 point)**

c) Assuming that in a single reaction $2\text{H} \rightarrow \text{H}_2$, the energy $2\mu_{\text{H}} - \mu_{\text{H}_2} = \Delta\mu$ is released, calculate the equilibrium fraction $(N_{\text{H}}/V)^2/(N_{\text{H}_2}/V)$ in terms of $\Delta\mu$ and $\lambda_{\text{H}} \equiv \lambda$ to obtain the law of mass action. **(1 point)**

d) Using your result from c) and assuming that the total number of particles is conserved, $N = N_{\text{H}} + 2N_{\text{H}_2}$, express the equilibrium density of hydrogen in atomic form, $c_{\text{H}} = N_{\text{H}}/V$, in terms of the total atomic density $c = N/V$. **(1 point)**

e) Is the relative concentration of atomic hydrogen, c_{H}/c , a monotonic function of c or does it have extrema? What is the value of c_{H}/c in the limits of low and high total atomic density? Draw a schematic plot of c_{H}/c as a function of c . Furthermore, calculate the total atomic density c at which exactly half of the hydrogen is dissociated, given by the condition $c_{\text{H}}/c = 1/2$. **(4 points)**

f) The interstellar medium consists largely of hydrogen at low density $c = 10^7 \text{ m}^{-3}$ at $T = 100 \text{ K}$, interspersed with clouds of $c = 10^{12} \text{ m}^{-3}$ at $T = 10 \text{ K}$. Based on the present calculation, in which form do you expect the hydrogen to be in the two different regions? Use the following constants:

$$\begin{aligned} h &= 6.63 \cdot 10^{-34} \text{ J s} \\ m &= 1.67 \cdot 10^{-27} \text{ J s}^2/\text{m}^2 \\ k_{\text{B}} &= 1.38 \cdot 10^{-23} \text{ J/K} \\ \Delta\mu &= 7.24 \cdot 10^{-19} \text{ J.} \end{aligned}$$

¹The Euler relation is actually more general, it also holds for functions with vector arguments: If $f(\alpha\vec{x}) = \alpha f(\vec{x})$ (for all $\alpha \in (0, \infty)$), then $f(\vec{x}) = \sum_i x_i (\partial f(\vec{x})/\partial x_i)$. But we will only need the one-dimensional special case here.

Comment: In fact, the hydrogen in the low-density regions of the interstellar space exists mainly in atomic form due to photodissociation under the influence of UV light. For details, see: Stecher and Williams, The Astrophysical Journal, Vol. 149, L29 (1967). (1 point)