

Visualizing a one-way protein encounter complex by ultrafast single-molecule mixing

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We combined rapid microfluidic mixing with single-molecule fluorescence resonance energy transfer to study the folding kinetics of the intrinsically disordered human protein α -synuclein. The time-resolution of 0.2 ms revealed initial collapse of the unfolded protein induced by binding with lipid mimics and subsequent rapid formation of transient structures in the encounter complex. The method also enabled analysis of rapid dissociation and unfolding of weakly bound complexes triggered by massive dilution.

α -synuclein is an intrinsically disordered protein whose folding is not solely encoded by its amino acid sequence but depends on its binding to other proteins or membranes in the cell. A general open question for such intrinsically disordered proteins is whether binding is initiated by formation of transient structures in the unfolded protein or whether folding is only induced after contact with binding partners¹. These two scenarios can be distinguished more directly using single-molecule methods, which avoid averaging the measurements of properties over many proteins. In particular, single-molecule fluorescence resonance energy transfer (smFRET) is a tool of choice for direct observation of folding pathways². When used in combination with a kinetic mixing apparatus to rapidly trigger transitions, smFRET can reveal the sequence of binding and folding steps, and uncover transiently formed intermediates.

Advances in microfluidic mixing technology during the last decade^{3,4} have resulted in devices with enhanced capabilities for studies of early events in protein folding⁵⁻⁷. However, accessing early steps during non-equilibrium folding with smFRET is a technical challenge because existing microfluidic setups rely on high flow velocities to rapidly trigger conformational changes. For example, to reach the fastest mixing times (<10 μ s), laminar sheath-flow microfluidic mixers typically require 1–10 m s⁻¹

flow velocities^{3,4}, which are incompatible with smFRET because a fluorescently labeled protein needs to spend enough time in the focal volume for sufficient photons to be collected. In recent years, several microfluidic setups have been specially designed for smFRET measurements at low flow velocities with improved time resolution⁸⁻¹⁰. Nevertheless, the earliest protein folding events reported so far have been at relatively long times on the order of ~10 ms^{11,12}, far from the current limit of ~100 μ s for smFRET time resolution.

To detect folding events on the 200- μ s to 100-ms timescale, we built a microfluidic device that creates an abrupt flow deceleration, allowing high-speed laminar mixing to be followed by smFRET detection in a slow flow (Fig. 1 and Supplementary Figs. 1–13). In the mixing region of the device (Fig. 1c), two high-speed buffer streams squeeze the protein stream into a layer as narrow as ~20 nm (Supplementary Note 1). Ligands and other small molecules diffuse rapidly across this layer and trigger protein transitions. Immediately downstream of the mixing region, >99.5% of the flow is abruptly diverted into two side outlets at a four-way junction (Fig. 1c). This diversion caused the protein flow velocity to decrease by about three orders of magnitude from a peak value of ~2 m s⁻¹ within ~10 μ m, resulting in an only ~200 μ s dead time between the medium exchange and earliest smFRET measurements (Fig. 1d). We optimized the four-way junction by extensive three-dimensional flow simulations to prevent vortices and flow instabilities. The protein flow greatly decelerated and expanded, and some of its streamlines overshot the junction and made >90° turns curling back toward the side channels. The inertia of the protein stream made it diverge in the z-direction (Fig. 1g). Hence, nearly all protein molecules interrogated traveled along the channel centerline and had uniform passage times through the mixing region.

We used this device with improved mixing time resolution to investigate the coupled binding and folding of α -synuclein using smFRET. Previous studies showed that interaction with negatively charged lipid bilayers can transform α -synuclein from a disordered state (U state) to a folded extended helix state (F state) and that addition of the lipid mimic SDS can reproduce this behavior¹³. Here we used SDS monomers as binding partners to explore the kinetic folding and unfolding paths between U and F states.

To test the conformations of α -synuclein with smFRET, we labeled α -synuclein with donor (Alexa Fluor 488) and acceptor (Alexa Fluor 594) dyes at residue positions 7 and 84 (ref. 14). We mounted the microfluidic device on a home-built confocal setup with two detectors that simultaneously recorded bursts of fluorescence as individual labeled proteins passed through the

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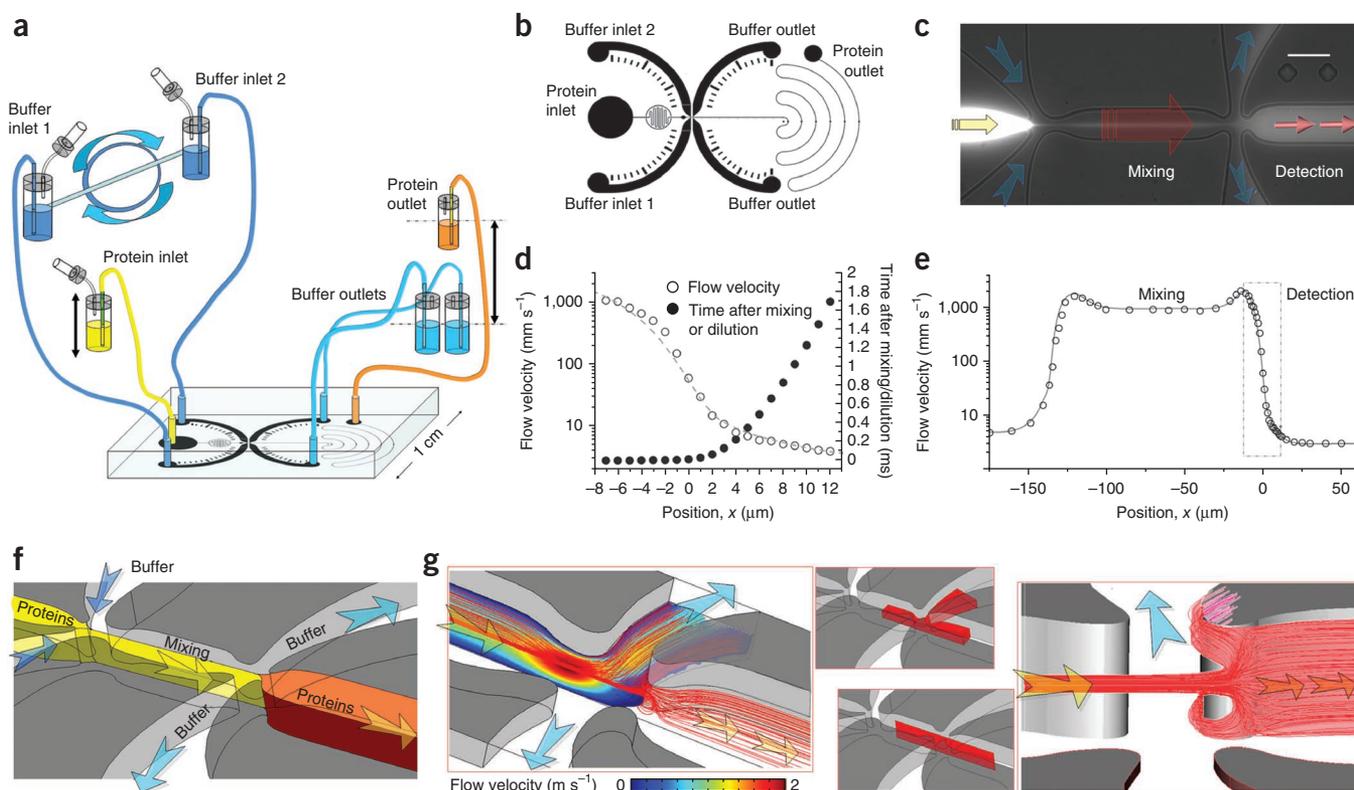


Figure 1 | Microfluidic setup for kinetic smFRET measurements. **(a)** After passing through a mixing region, the protein stream and two buffer streams (fed from three separate inlets) are directed to three outlets, which are connected to separate reservoirs, whose heights are adjusted to tune mixing and dilution. **(b)** Diagram of the device showing inlets and outlets. **(c)** Micrograph of the functional region with a fluorescent solution fed to the protein inlet. Arrows indicate buffer flow (blue), protein flow before mixing (yellow) and protein flow in mixing and detection regions (red). Scale bar, 25 μm . **(d)** Flow velocity as measured with fluorescence correlation spectroscopy and time after mixing or dilution, both as functions of the position along the channel in the deceleration region; dashed line is from three-dimensional flow simulations (Comsol). **(e)** Flow velocity (as in **d**) in the focusing, mixing and deceleration regions. **(f)** Flow diagram illustrating how the protein stream is squeezed horizontally by two buffer streams for medium exchange and then directed to the smFRET detection region. **(g)** Simulations of flow velocities and streamlines in two device regions highlighted in insets.

confocal volume. We analyzed donor and acceptor photon counts to generate histograms of FRET efficiency (E_{FRET}), which translated into structural distributions^{2,14}. By placing the confocal volume at different points along the detection region (Fig. 1c,d), we followed the evolution of α -synuclein E_{FRET} distributions from 200 μs to >100 ms after mixing (Fig. 2a–c).

We first used the microfluidic device to rapidly change the buffer of U-state α -synuclein to a solution promoting formation of the F state. This solution contained 1.2 mM SDS binding partner and 20 μM unlabeled α -synuclein, as in a previous study¹⁴. We observed multiple folding events en route from U state to F state (Fig. 2a). The earliest smFRET measurement, at only 200 μs after the addition of SDS, already showed a shift of the histogram peak from an E_{FRET} value of 0.42 (U state in the absence of SDS) to a higher E_{FRET} value, indicating an initial rapid collapse (Supplementary Figs. 14 and 15). After this initial compaction (E_{FRET} of ~ 0.6 at 500 μs), we observed a break in the shift of the FRET peak and a transition toward a high- E_{FRET} intermediate state (I) with E_{FRET} of ~ 0.8 , fully formed after 1 ms. In the next 11 ms, this intermediate was converted via a two-state process to the expected final F-state structure characterized by a low- E_{FRET} peak (Fig. 2c and Supplementary Figs. 16–21).

Observation of these multiple folding steps raises an interesting question: are similar intermediates observed when protein-SDS

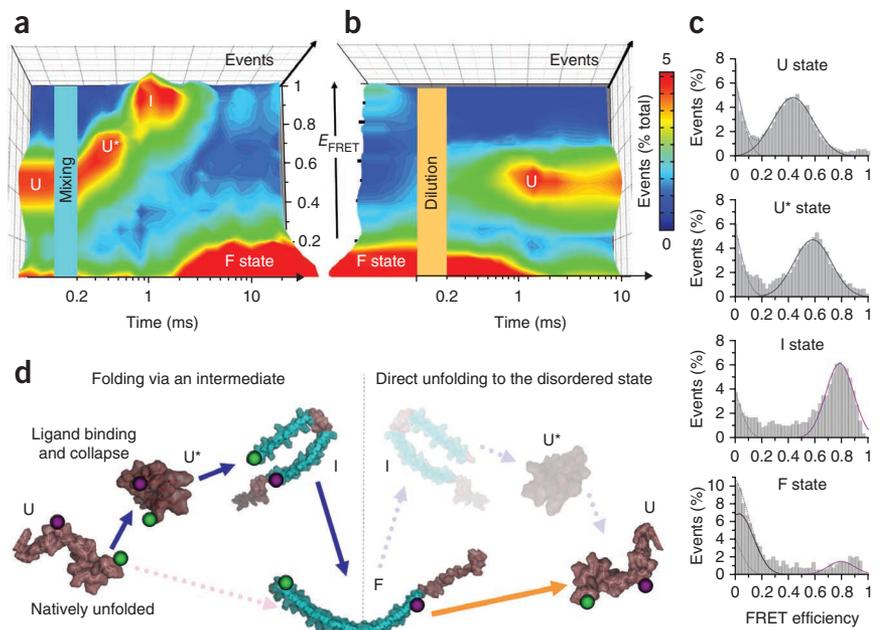
complexes dissociate? To address this question, we used another powerful feature of the microfluidic device: the ability to rapidly dilute the protein sample. The technique takes advantage of the lateral diffusion (dilution) of proteins out of their focused, narrow stream while they travel through the mixing region. By applying slightly asymmetric flow rates of the two buffer streams, we could steer the middle of the focused stream away from the detection region and direct there a solution with greatly reduced protein concentration, thus diluting a protein sample by a factor as high as 10^6 within ~ 200 μs (Supplementary Fig. 11 and Supplementary Note 1).

For this unfolding experiment (Fig. 2b), we used this asymmetric focusing to dilute the SDS-synuclein complexes in the F state by >100 -fold, thus triggering rapid dissociation of binding partners and protein unfolding. The low- E_{FRET} F-state peak present at 200 μs disappeared after ~ 1 ms with concomitant formation of a U-state peak, indicating that the weakly bound complexes dissociated rapidly, and a simple two-state transition occurred. We observed no folding intermediates with high E_{FRET} along the unfolding pathway (Fig. 2b,c and Supplementary Figs. 22 and 23).

The smFRET kinetic data suggest that at early stages, binding of SDS molecules to the unfolded α -synuclein continuously collapsed the protein in an ‘encounter complex’ that preceded formation of large-scale structures (Fig. 2d). The folding of

Figure 2 | Folding and unfolding of α -synuclein.

(a,b) Histograms of E_{FRET} for the folding (a) and unfolding (b) reactions, obtained at different time points, are used to generate a three-dimensional histogram (~50,000 events) in coordinates of time and E_{FRET} , with the percentage of total events color-coded as indicated. (c) Representative E_{FRET} histograms for various states: intrinsically disordered (U state, $E_{\text{FRET}} \approx 0.42$, obtained before mixing), collapsed unfolded (U^* state, $E_{\text{FRET}} \approx 0.6$, 490 μs after mixing), intermediate (I state, $E_{\text{FRET}} \approx 0.8$, 1.2 ms after mixing) and extended structures (F state, $E_{\text{FRET}} \approx 0.1$, >10 ms after mixing). (d) Model of the α -synuclein conformational transitions; the mirror representation emphasizes the asymmetry between the folding and unfolding pathways. Shown are random coil (brown), α -helix (turquoise), as well as donor (green sphere) and acceptor (purple sphere) dye molecules.



α -synuclein was induced in the complex and proceeded by way of a partly folded intermediate state. The smFRET peak of this intermediate closely resembled the smFRET peak observed previously for the protein bound to SDS micelles in a helix-kink-helix structure^{13,14}. This 'broken helix' conformation has been previously believed to be an artifact of highly curved micelles, with the elongated helix being the relevant membrane-stabilized state. Our data (Fig. 2a,c) showed that the broken helix structure may be relevant as an essential kinetic step in the process of folding toward the membrane-bound extended helix.

Another interesting observation was that the folding from the U state to the F state and unfolding from the F state to the U state followed radically different pathways (Fig. 2d). On a molecular level, it opens an intriguing possibility: the intermediate folding step suggests that not all 'binding sites' are readily accessible in the U-state α -synuclein, potentially owing to a nonrandom native state conformation. These sites would only be exposed after the formation of the intermediate folded (I) state, allowing partners to bind to the protein, thus inducing the final conformational switch. In contrast, when the binding partners are rapidly removed from the F state, the different SDS sites may exhibit similar high dissociation rates, resulting in rapid and random loss of overall structure and a fast collapse with no detectable intermediates.

The single-molecule microfluidic mixer described here enabled detailed mapping of the binding-induced folding kinetics of α -synuclein. It revealed transient conformations and steps that would have remained hidden in both bulk and single-molecule experiments using existing instrumentation. The device brings the time resolution of smFRET kinetics close to the physical limit imposed by current fluorescent dyes, as the latency of 200 μs is nearly the same as the minimal time required to collect sufficient photons¹⁵. Moreover, the capability of adjustable rapid and massive dilution in the device provides the ability to study weakly bound complexes and transient oligomers, which usually dissociate quickly at single-molecule concentrations.

METHODS

Methods and any associated references are available in the online version of the paper at <http://www.nature.com/naturemethods/>.

Note: Supplementary information is available on the Nature Methods website.

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AUTHOR CONTRIBUTIONS

Y.G., A.G. and A.A.D. designed research; Y.G. performed smFRET experiments; Y.G. and V.V. characterized the device; A.C.M.F. provided α -synuclein expertise and samples; E.A.L. provided instrumentation support; and all authors contributed to writing the paper.

COMPETING FINANCIAL INTERESTS

The authors declare no competing financial interests.

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ONLINE METHODS

Microfluidic device design. The microfluidic device was designed to provide fast and efficient medium exchange in a laminar sheath flow in its mixing region and abrupt deceleration of the sample stream at a four-way junction, which is situated immediately upstream of its detection region. The shape of channels in the device and flow conditions in the detection region are optimized for smFRET spectroscopy (**Supplementary Fig. 1**).

Similarly to the previous microfluidic hydrodynamic focusing mixers with hydrodynamic focusing in a laminar sheath flow, the buffer exchange in the mixing region occurs by molecular diffusion between a thin sample stream and two wide buffer streams on its sides. Unlike previous designs, the mixing region is separated from the detection region, where smFRET interrogation is performed, and the flow speeds in these two regions can be adjusted independently by varying the part of flow diverted into two side channels at the four-way junction between the mixing and detection regions. Typically, ~99.8% of flow in the mixing region is diverted to the side channels, leading to about three orders of magnitude reduction of flow velocity in the detection region as compared with the mixing region. This mode of operation makes it possible to combine rapid buffer exchange in a fast flow in the mixing region (velocities up to 2 m s^{-1}) with smFRET detection in a slow flow in the detection region (velocities of a few millimeters per second). The deceleration of the sample flow from the high velocities in the mixing region to the low velocities in the detection region occurs over a short distance, leading to a short deceleration time (in the Lagrangian reference frame of the moving liquid) and small dead time between mixing and smFRET detection. Small dead times provided by the device made it possible to access sub-millisecond timescales (comparable to those in stopped-flow instruments) in a system that requires a very small amount of sample and is compatible with high numerical aperture (NA) confocal optical setup and high photon efficiency smFRET detection.

The device also makes it possible to take advantage of the lateral diffusion of the sample molecules in the mixing region that has a characteristic distance of $\sim 0.1 \mu\text{m}$ during $\sim 100 \mu\text{s}$ long passage through the mixing region. It is worth noting that with the typical initial sample stream width of $\sim 20 \text{ nm}$, this diffusive broadening causes a major change in the distribution of concentration across the stream, bringing a large part of the sample molecules outside of geometrical boundaries of the stream. (Those boundaries can be defined as boundaries in the absence of diffusion.) The nearly three orders of magnitude deceleration of the sample stream at the transition from the mixing to the detection region leads to greater than two orders of magnitude magnification of distances across the stream in the lateral direction. (There is also some expansion of the stream and concomitant increase of distances in the out-of-plane direction; **Fig. 1** and **Supplementary Fig. 12**.) As a result, in the detection region, the characteristic length scale of sample concentration variations across the stream becomes $>10 \mu\text{m}$, making the sample concentration nearly uniform within the focal volume of the smFRET setup ($\sim 0.3 \mu\text{m}$ in diameter). This feature makes it possible to selectively interrogate parts of the sample stream with molecules that diffused particularly far from the geometrical boundaries of the sample stream in the mixing region and with sample concentration dramatically reduced compared with the original sample solution. A part of the sample

stream with a desired sample concentration is directed into the detection region by adjusting the flow rates of the two focusing buffer streams to steer the center of the sample stream somewhat away from the detection region. With this technique, a protein sample can be diluted by factors as high as 10^6 within $<1 \text{ ms}$. The flow toward the two side channels at the four-way junction is kept symmetric to prevent flow instabilities.

Device measurement timescales. The device has three smFRET measurement areas. Before entering the mixing region, the sample solution passes through a wide chamber where the flow is sufficiently slow for smFRET to be essentially diffusion-limited. After exiting the mixing region, the sample flows through a 5-mm-long detection region (**Supplementary Fig. 1c**). At normal device operation conditions, the most upstream and the most downstream smFRET measurement points in this region correspond, respectively, to $\sim 200 \mu\text{s}$ and $\sim 1 \text{ s}$ from the time of buffer exchange, making the entire range from $200 \mu\text{s}$ to $\sim 1 \text{ s}$ accessible in a single experiment. Special markers in the microfluidic chip positioned every $100 \mu\text{m}$ along the detection region help to determine the precise distance of a given measurement point from the detection region entrance (**Supplementary Fig. 1c**). The detection region is connected to the protein outlet by a long serpentine channel, which makes it possible to access protein conformations over an extended time range (up to 5 min, at reduced flow rates) by measuring smFRET at various points along the serpentine channel (**Supplementary Fig. 1c**). After $\sim 100 \text{ ms}$, substantial uncertainty is introduced into the distance-time conversion due to the combined effect of diffusion of sample molecules across the channel and parabolic profile of velocity of the laminar flow in the detection region and in the serpentine channel. Nevertheless, measurements at large reaction times still provide information on 'equilibrium' protein conformations and semiquantitative data on the kinetics of slow folding processes.

Device channel dimensions. The choice of the channel depth was dictated by multiple considerations. Channels with small diameters are preferable for reducing the time of deceleration and the measurement dead time, giving access to better time resolution. In addition, small channel dimensions lead to small values of the Reynolds number, Re , at a given flow speed, making the flow more stable and less prone to vortex buildup near corners. In contrast, channels with small diameters result in high flow resistances, making it difficult to combine high flow velocity in the mixing region (for example, 2 m s^{-1}) with low driving pressure (preferably $\leq 30 \text{ kPa}$ to prevent large deformations of soft PDMS walls and separation of the PDMS chip from the substrate). Small channels are also more difficult and costly to fabricate and more prone to clogging, especially when their diameters are less than a few micrometers. Lastly, performing measurements with the confocal volume close to walls may be detrimental to smFRET signal quality. In particular, in early device designs, we noticed high autofluorescence near channel bottoms, apparently because of the proximity to the borosilicate cover glass substrate. This autofluorescence significantly contributed to the smFRET background for distances up to $20 \mu\text{m}$ from the glass surface. We found a uniform channel depth of $15 \mu\text{m}$ to be a reasonable choice, making it possible to reach sufficiently short reaction dead times and suppress vortices without using excessively high driving pressure.

The problem of cover glass autofluorescence was resolved by bonding the microfluidic chips to cover glasses coated with a 50 μm layer of PDMS. The autofluorescence of PDMS at the wavelength used in the experiments was much lower than that of glass (2 counts ms^{-1} versus 20 counts ms^{-1}). Therefore, the PDMS layer on the cover glass considerably reduced the noise and allowed clean smFRET measurements in channels as shallow as 10 μm without using expensive fused silica cover glasses as substrates.

If not adjusted for, the addition of a PDMS layer on a cover glass can lead to spherical aberrations in the optical setup, reducing the confocality and the efficiency of light focusing and collection. When, as in the present study, a water immersion objective is used, the spherical aberrations originate from the difference between the refractive index of PDMS ($n_p = 1.41$) and water ($n_w = 1.33$), $\Delta n_p = n_p - n_w = 0.08$. These aberrations are of the same nature as the spherical aberrations resulting from the difference between the refractive index of glass ($n_p = 1.51$) and water, $\Delta n_g = n_p - n_w = 0.18$ and a mismatch between the actual thickness of the cover glass and the cover glass thickness the objective is optimized for. In terms of spherical aberrations, the addition of a PDMS layer with a thickness $d = 50 \mu\text{m}$ is nearly equivalent to the addition of $d' = d \times \Delta n_p / \Delta n_g = 22 \mu\text{m}$ of cover glass thickness. Therefore, to correct for the PDMS layer, we set the correction collar to a cover glass thickness 0.02 mm greater than the actual thickness of the cover glass. Specifically, for #1 cover glasses used in the experiments, the correction collar was set to a thickness of 0.17 mm, corresponding to a #1.5 cover glass.

To check whether the smFRET data were not affected by the possible changes in focusing and light collection, we compared histograms of FRET efficiencies for two different folded states, in the device with PDMS-coated cover glass and in a cuvette with a plain glass bottom (Supplementary Fig. 13). No data degradation was found in the device versus cuvette, which was consistent with the absence of appreciable optical aberrations in the experiment in the device.

The width of the channel in the detection region was chosen at 25 μm , so the beams focused and collected by the high-NA objective used in the experiments (63 \times , 1.2 NA water immersion) would not be distorted by passing through the PDMS side-walls of the channel, which had a refractive index of ~ 1.41 as compared to 1.33 for water. The lengths and widths of other channels were selected to obtain a desired range of flow velocities in the mixing and detection regions within a convenient range of driving pressures applied to the device inlets.

Device fabrication. The microfluidic chip is a poly(dimethylsiloxane) (PDMS; Sylgard 184 (Dow Corning)) cast of a UV lithography–microfabricated master mold. To make the master mold, a 15 μm layer of UV light–curable epoxy SU8 2015 (MicroChem) was spin-coated onto a 5-inch silicon wafer, exposed to UV-light through a specially designed chrome photomask (Photo Sciences; 0.5 μm resolution), cured on a hot plate, and developed. The resulting master mold was silanized by exposure to a vapor of chlorotrimethylsilane (Sigma-Aldrich) at room temperature for 90 s to facilitate release of the PDMS replicas. Next, PDMS pre-polymer was poured onto the master mold and baked at 80 $^\circ\text{C}$ for 1 h. The PDMS casts ($\sim 5 \text{ mm}$ thick) were peeled off the master mold and cut into individual chips, which were trimmed to size and punched with luer stubs to make inlet and outlet holes. PDMS-coated cover glasses were made by spin-coating a 50 μm

thick layer of PDMS pre-polymer on #1 microscope cover glasses and baking the cover glasses at 80 $^\circ\text{C}$ for 1 h. The microfabricated PDMS chips were then reversibly bonded to the PDMS-coated cover glasses by overnight baking in an 80 $^\circ\text{C}$ oven. The PDMS–PDMS bond withstood pressures as high as 50 kPa in the device without delamination (higher pressures were not tested since 50 kPa was more than sufficient for our experiments). The master mold had four copies of the device, and a single master mold was used for making all the PDMS casts. Devices were cleaned by boiling in 0.1 M HCl for 1 h and then bonded to fresh PDMS-coated cover glasses for repeated use. We are willing to assist with technical advice to facilitate utilization of the device and technique by other research groups, and to provide devices for testing in specific cases as feasible.

Protein expression, purification and labeling. Mutagenesis, expression and purification of mutant proteins were carried out following the procedure described previously¹³. The plasmid construct for wildtype α -synuclein was kindly provided by R.L. Nussbaum (US National Institutes of Health). Preparation of labeled protein was carried out as reported previously¹⁴. Briefly, the G7C,G84C mutant was sequentially reacted with Alexa Fluor 488 maleimide (donor) and Alexa Fluor 594 maleimide (acceptor) dyes (Molecular Probes). Singly labeled protein was purified after the first step of donor labeling, and this was followed by labeling with an excess of acceptor dye, and another round of purification. Protein identity and purity were ascertained by high-performance liquid chromatography and mass spectrometry analyses.

Single-molecule FRET measurements. Single-molecule FRET measurements were performed as previously described¹⁴. Briefly, experiments were carried out on a laser confocal microscope system home-built around an Axiovert 200 microscope (Zeiss). Excitation of single molecules was achieved by focusing a beam derived from the 488 nm line of a 543-AP-A01 tunable argon-ion laser (Melles Griot) inside the sample solution, 7.5 μm from the surface of the PDMS layer on a #1 cover glass, using a water immersion objective (63 \times , 1.2 NA, Zeiss). Donor and acceptor single-molecule fluorescence signals were collected through the same objective, were separated from the excitation light by using a dichroic mirror (Q495LP, Chroma), and spatially filtered using a 100 μm pinhole. A second dichroic mirror (560 DCXR, Chroma) was then used to separate donor and acceptor emission components, which were further filtered using an HQ 525/50M band-pass filter (donor, Chroma) and a 590 LPV2 long-pass filter (acceptor, Chroma), respectively. These photon streams were then detected using two SPCM-AQR-14 avalanche photodiode (APD) photon counting modules (Perkin-Elmer Optoelectronics). Photon counts were recorded with a counter card (PCI 6602, National Instruments) interfaced with a computer using a code in LabView (National Instruments).

The FRET efficiency histograms reported in the paper were generated by using a two-channel data collection mode to record donor and acceptor light intensities (photon counts) as functions of time. To ensure that virtually all of the detected photon bursts were from single α -synuclein molecules, the concentration of donor-acceptor solutions in the detection region was $\sim 100 \text{ pM}$ (in fluorophore concentration). The protein sample solution also contained a higher concentration of unlabeled α -synuclein to avoid

the strong effect of even very low concentrations of SDS micelles on labeled synuclein conformation by rebalancing the stoichiometries of the SDS and synuclein, as discussed previously^{13,14}. We note that control experiments in a previous paper¹³ using the concentration dependence of synuclein-SDS binding-folding excluded measurable association/aggregation. Hence, the observations reflect intrinsic monomer folding/unfolding properties. In addition, the molecules interrogated by smFRET travel near the centerline of the channel, minimizing the effect of their interaction with channel walls. The background counts, leakage of donor into acceptor channel (~8%) and the direct excitation of acceptor (~5%) were used to correct the signals before calculating FRET efficiencies.

A threshold of 40 counts (the sum of the photon counts from the two channels) was used to separate the photon burst events due to fluorescently labeled protein molecules from background. Ratiometrically measured E_{FRET} are defined on the basis of the measured donor (I_{D}) and acceptor (I_{A}) fluorescence intensities as:

$$E_{\text{FRET}} = \frac{I_{\text{A}}}{I_{\text{A}} + \gamma I_{\text{D}}}$$

γ is a correction factor dependent on the donor (Φ_{D}) and acceptor (Φ_{A}) quantum yields, and the donor channel (η_{D}) and acceptor channel (η_{A}) detection efficiencies, defined as:

$$\gamma = \frac{\eta_{\text{A}}\Phi_{\text{A}}}{\eta_{\text{D}}\Phi_{\text{D}}}$$

For the dual-labeled α -synuclein molecules used in the study, γ was previously found to be close to 1; this value was used here. FRET efficiencies were calculated for each validated event and plotted in the form of a histogram.

At least ten measurements were made for each position, x , along the detection region of the microfluidic device to construct a FRET histogram representing at least 2,000 different molecules; the smFRET three-dimensional map displayed in **Figure 2** was compiled from 23 separate histograms taken at different reaction time points (different positions along the detection region), representing a total of more than 50,000 individual events. The histograms were fitted with Gaussian functions using Origin (OriginLab), IGOR PRO (WaveMetrics) and Excel (Microsoft), and the peak positions and areas under the peaks were obtained from the fitting parameters.

We verified in previous work that the F state results in a low E_{FRET} peak, which can be distinguished from the zero-peak using alternating excitation of the donor and acceptor dyes¹⁴. Here for the cases with F state, the zero-peak parameters were fixed, allowing the F state to be fitted.