

Laser-Ion Acceleration

R. Sauerbrey

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Laser-Particle-Acceleration

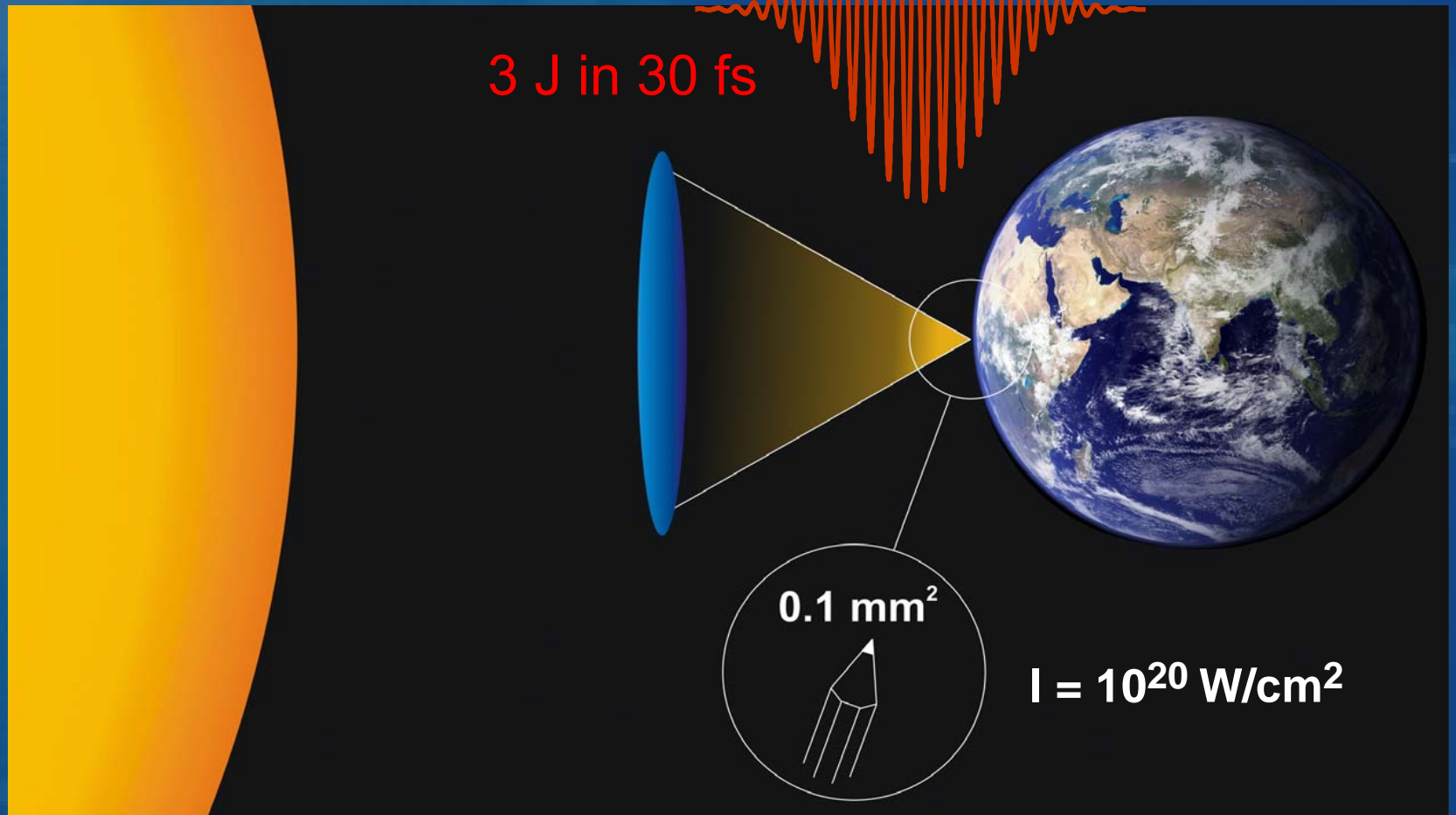
R. Sauerbrey

- What happens to one electron in the laser field ?
- Laser electron acceleration
- Laser ion acceleration
- Applications

100 TW Laser

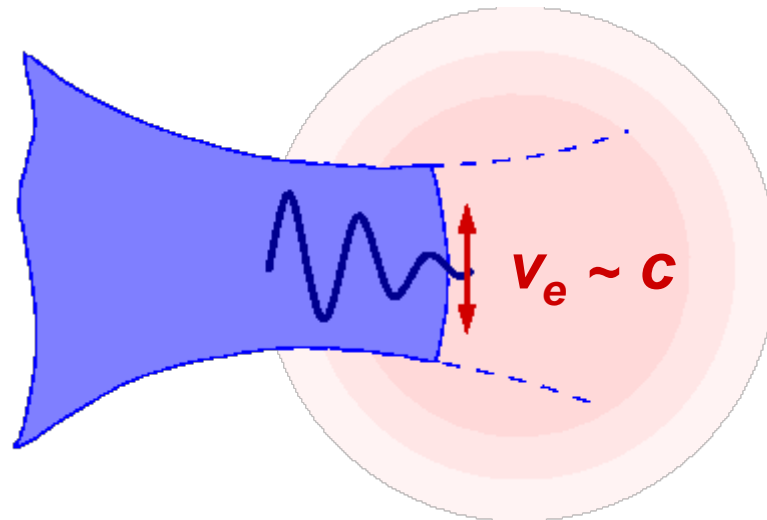
$$I = 10^{20} \text{ W/cm}^2$$

$$E_0 = 10^{12} \text{ V/m}$$



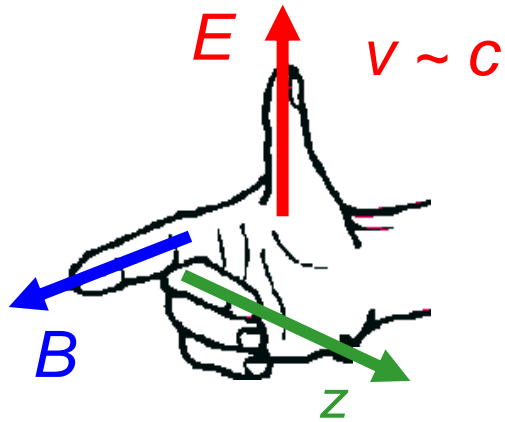
For laser intensities exceeding $I \sim 10^{18} \text{ W/cm}^2$, the electron quiver motion becomes relativistic within half a period

target: one electron

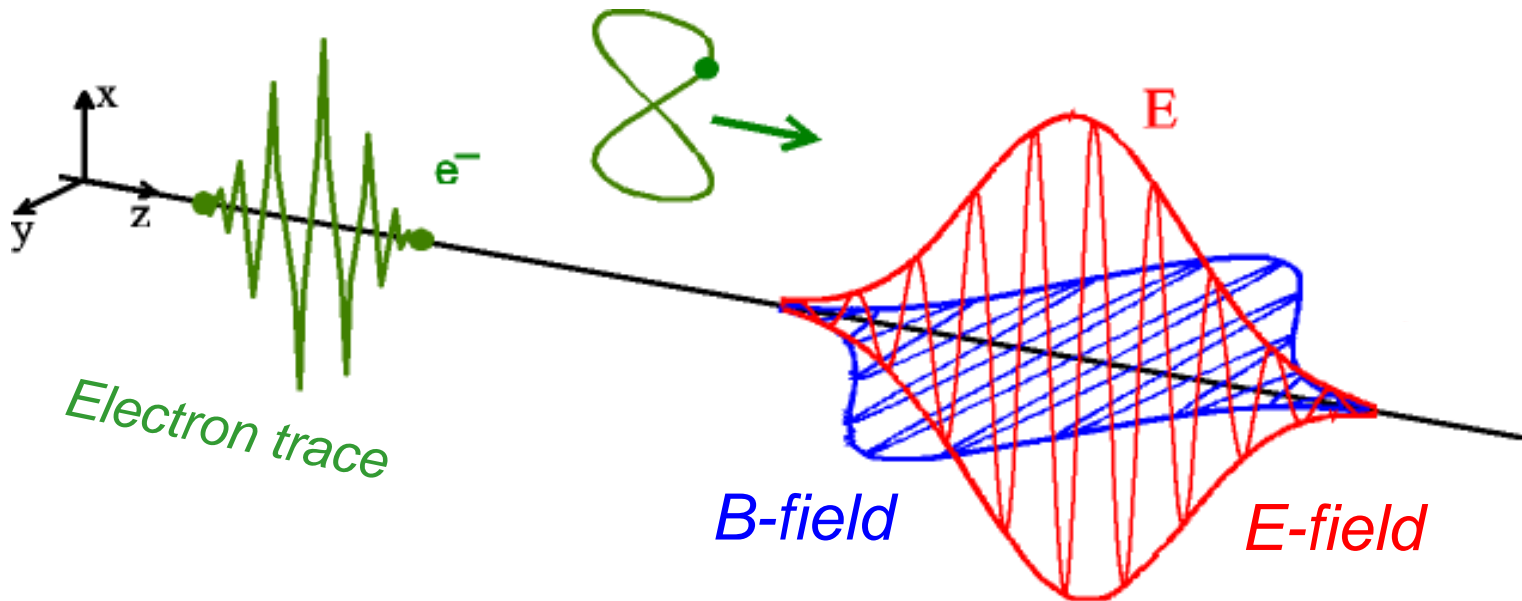


- ▶ *mass increase*
- ▶ *forward acceleration due to Lorentz force*
- ▶ *anharmonic osc.*

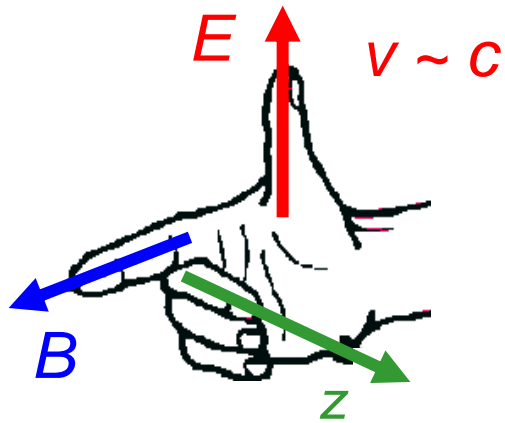
$$a_0 = \frac{eE_0}{\omega m_e c} \quad I = \frac{E_0 B_0}{\mu_0} = \frac{E_0^2}{\mu_0 c} = \frac{a_0^2}{\lambda^2 [\mu\text{m}]} \cdot 1.4 \cdot 10^{18} \frac{\text{W}}{\text{cm}^2}$$



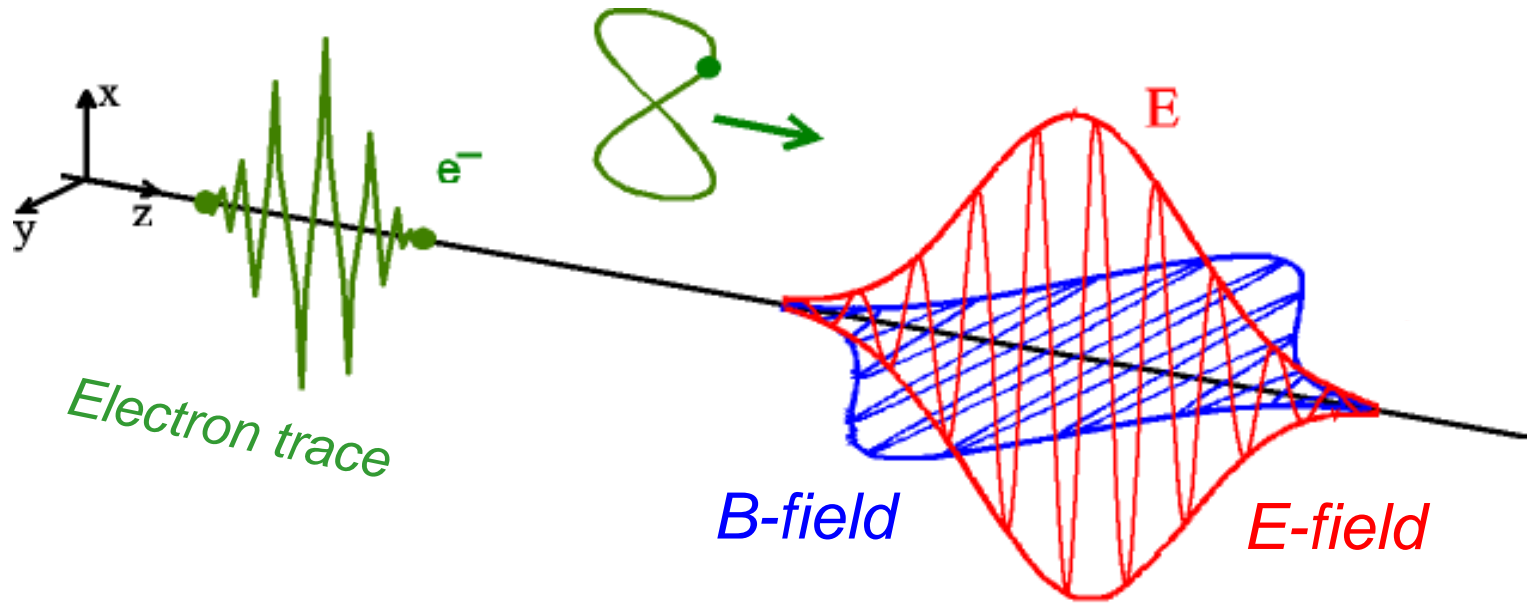
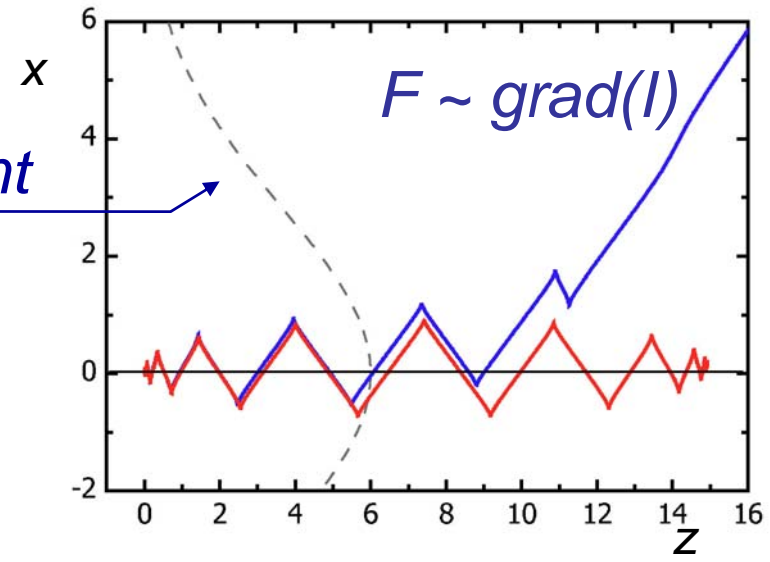
$$\vec{F} = e\vec{E} + e\vec{v} \times \vec{B} \quad (B_0 = E_0/c)$$



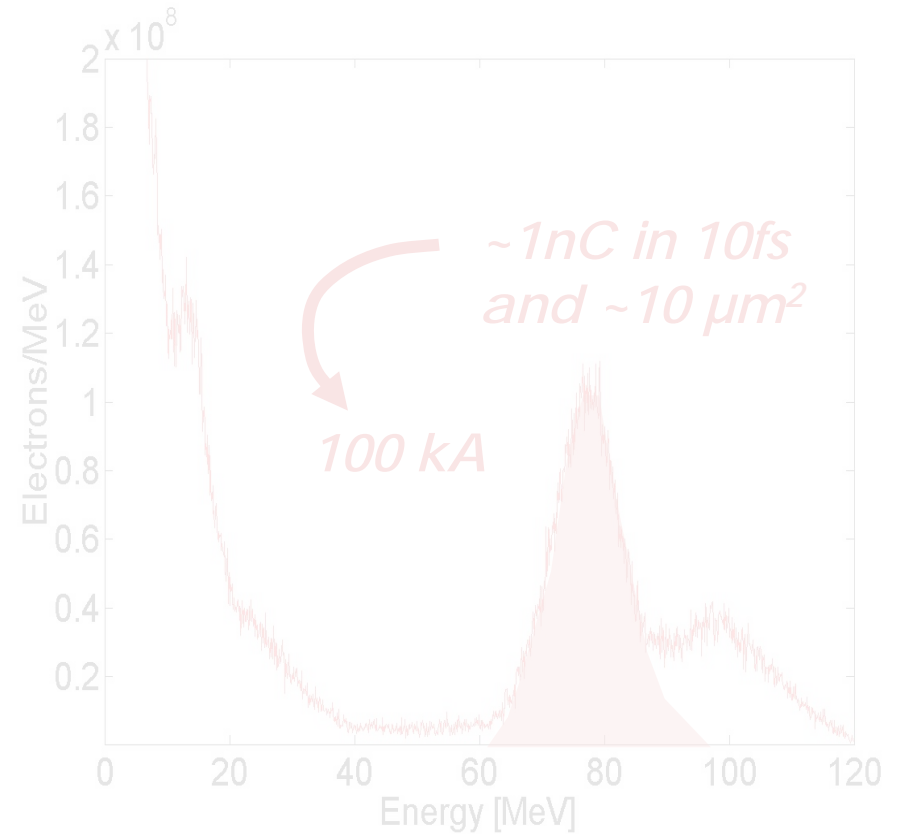
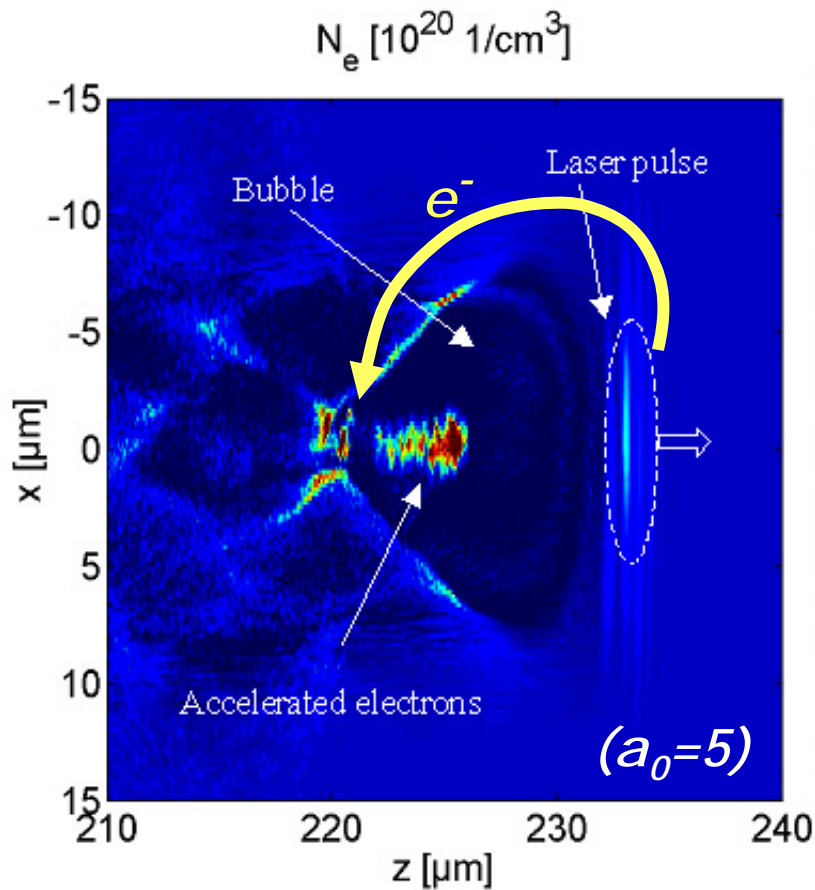
single electron dynamics



*intensity gradient
ponderomotive
acceleration*



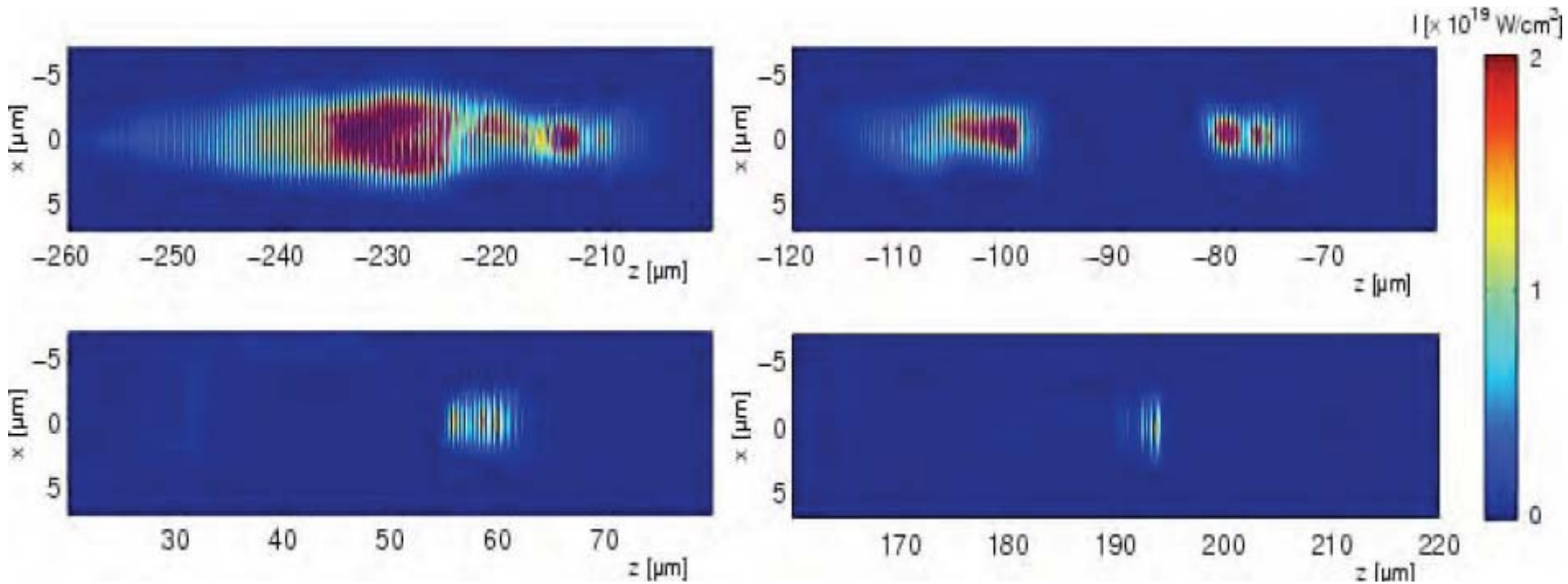
relativistic bubble regime



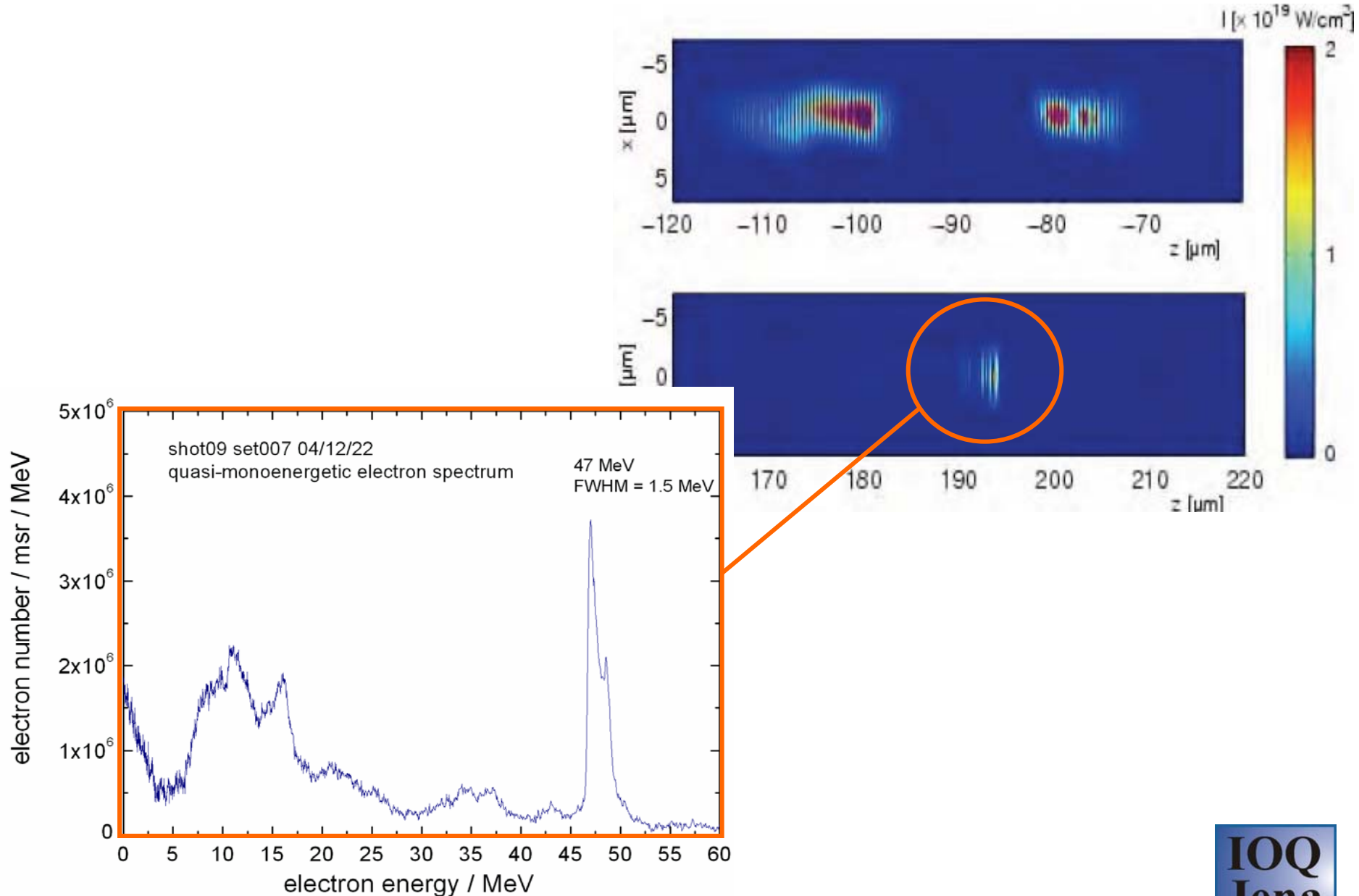
plasma wavelength < pulse length

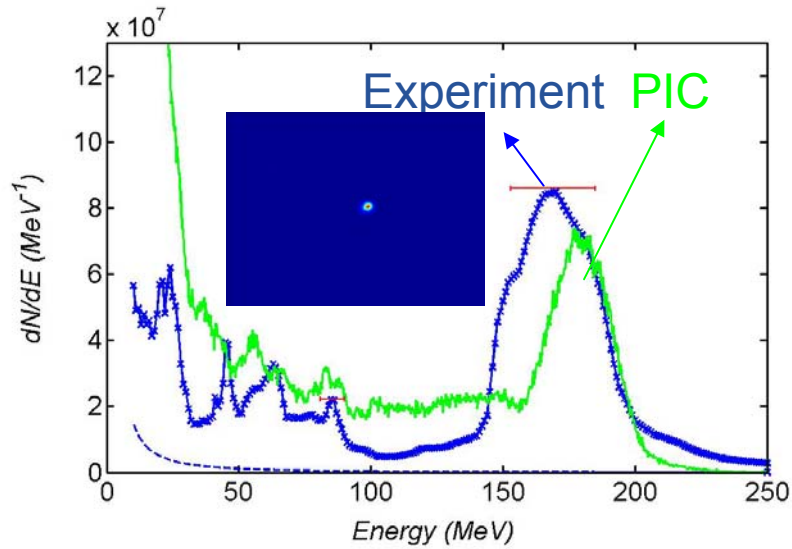
[M. Geissler, NJP 8 (2006) 186]

relativistic pulse shortening ...



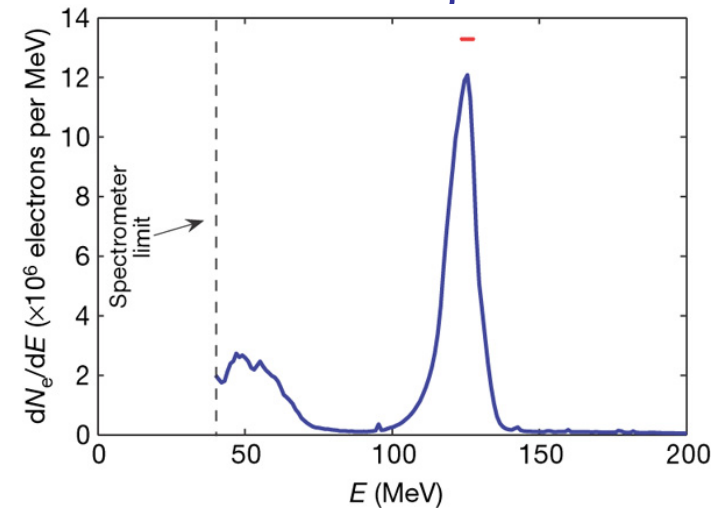
relativistic pulse shortening ...





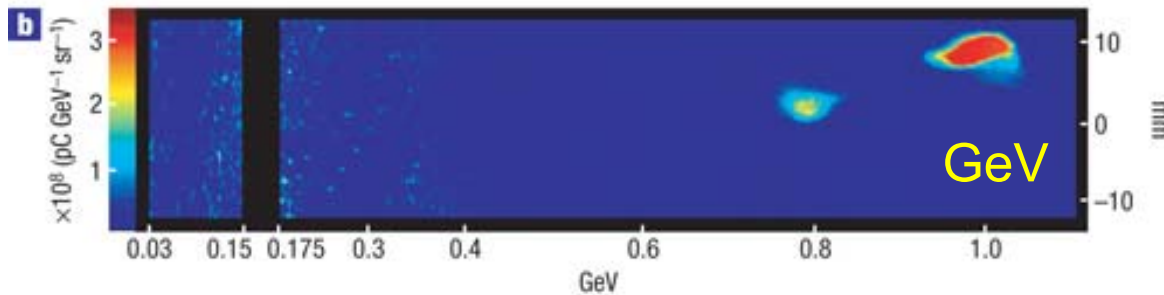
S. Mangles et al., C. Geddes et al., J. Faure et al.,
in Nature 431 (2004)

controlled injection



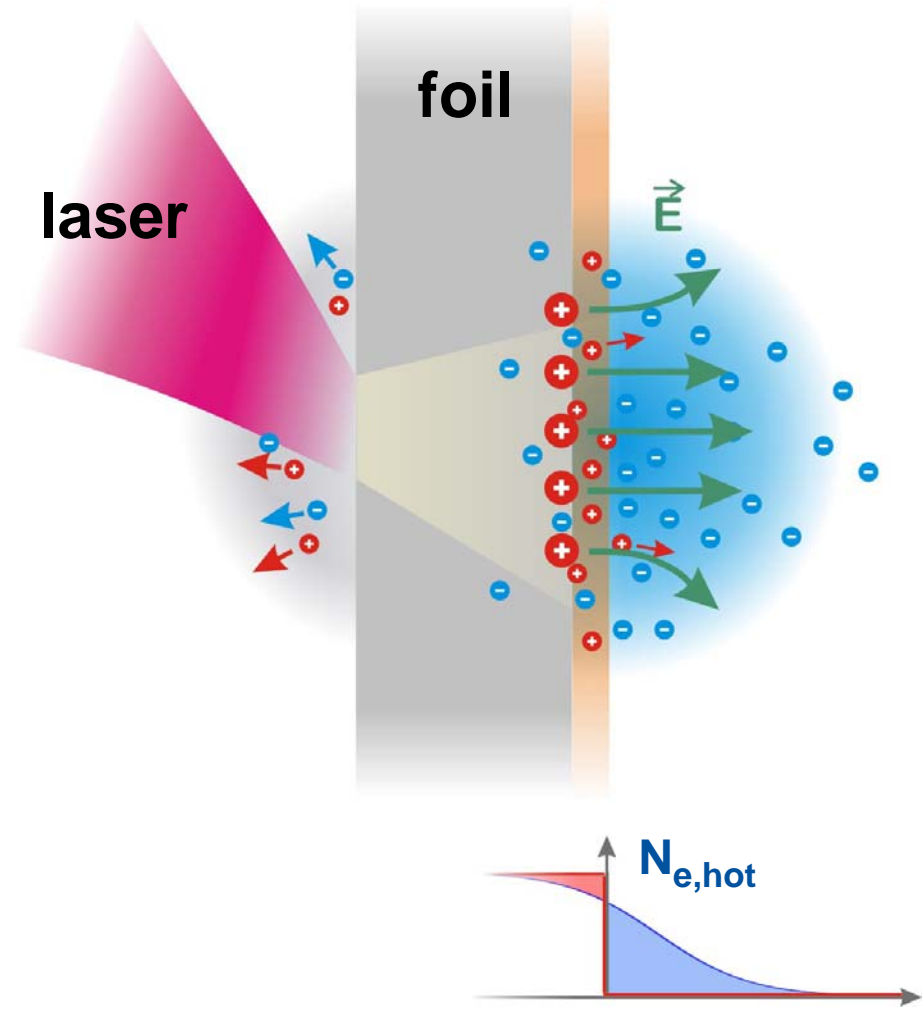
J. Faure, et al., Nature 444 (2006) 737

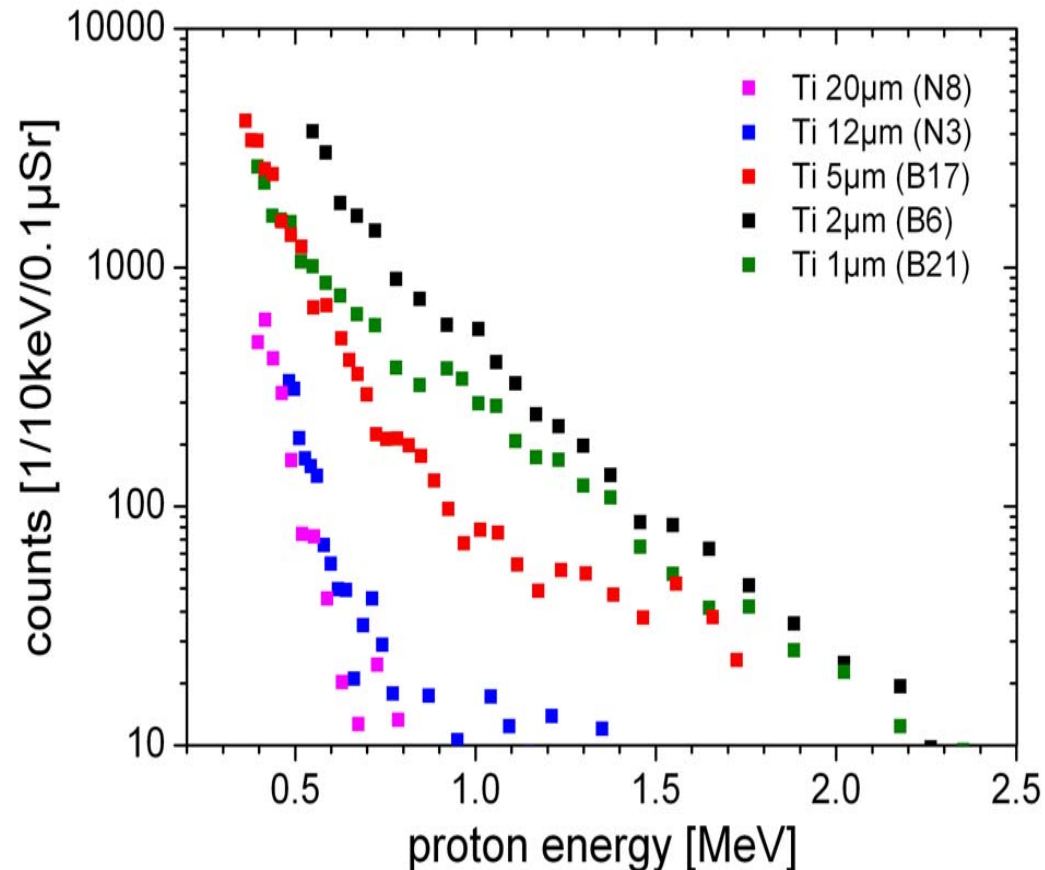
matched discharge capillary



W. Leemans, et al.,
Nature physics 2 (2006) 696

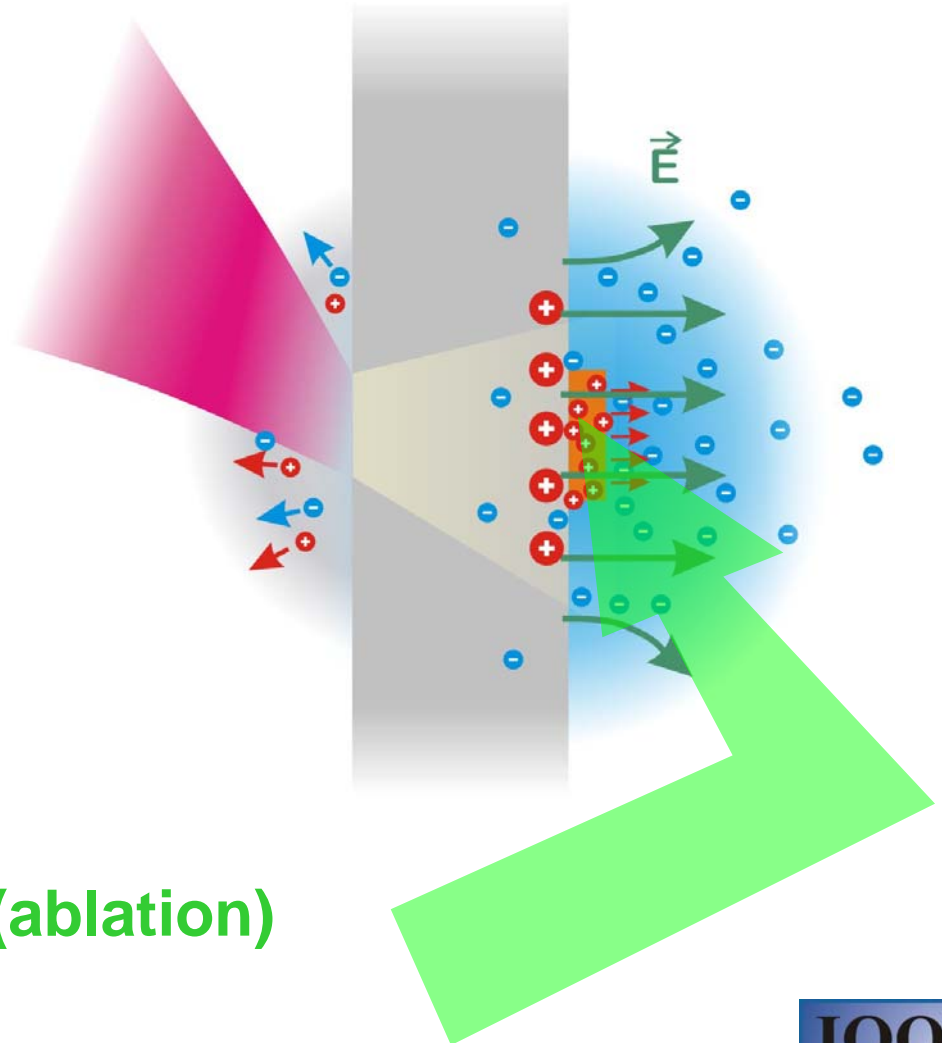
- **electron acceleration**
- **hot (MeV) electrons penetrate the (μm) foil**
- **quasi static field forms normal to target surface, source size \gg laser spot**

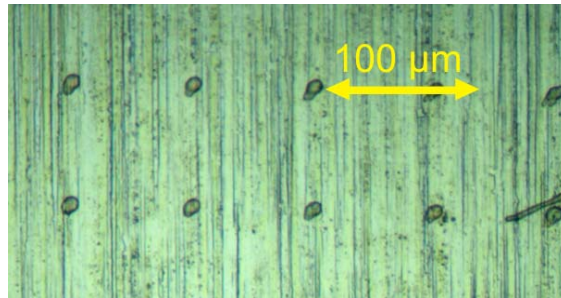
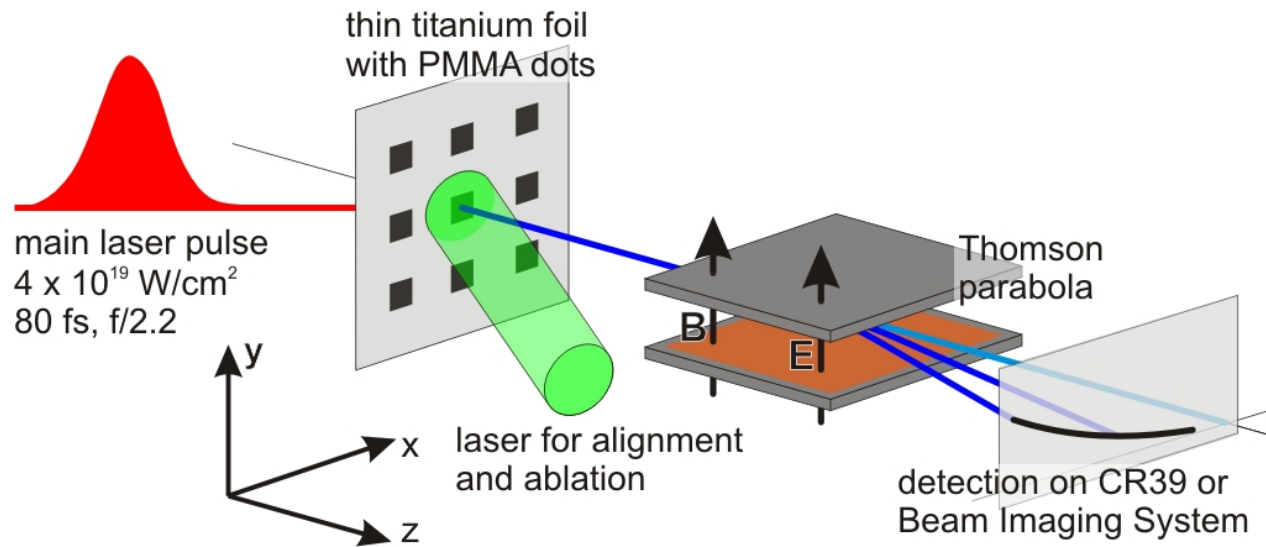




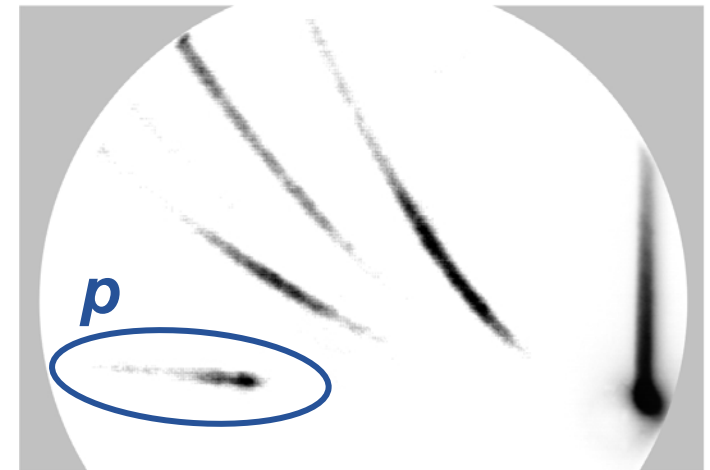
quasi-neutral pulse with exponential energy distribution (with max. energy depending on laser pulse duration, energy, and target thickness)

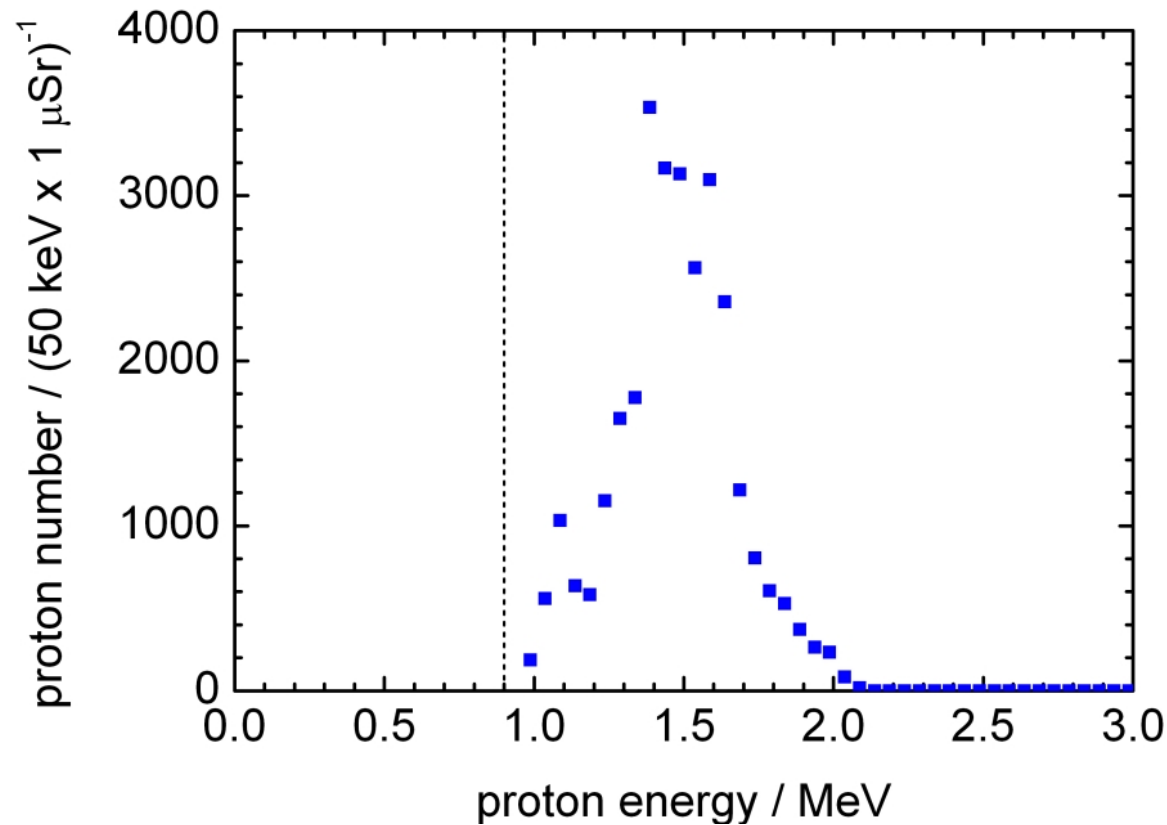
- enhance yield in the central, homogeneous region by applying a proton rich „dot“
- use thin dot (to avoid temporal field depletion and shielding)
- Careful backside cleaning (ablation) increases the fidelity





Lith. polymer coating on Ti-foil
 $8 \times 8 - 20 \times 20 \mu\text{m}^2$ base area
 $0.15 - 1.0 \mu\text{m}$ height

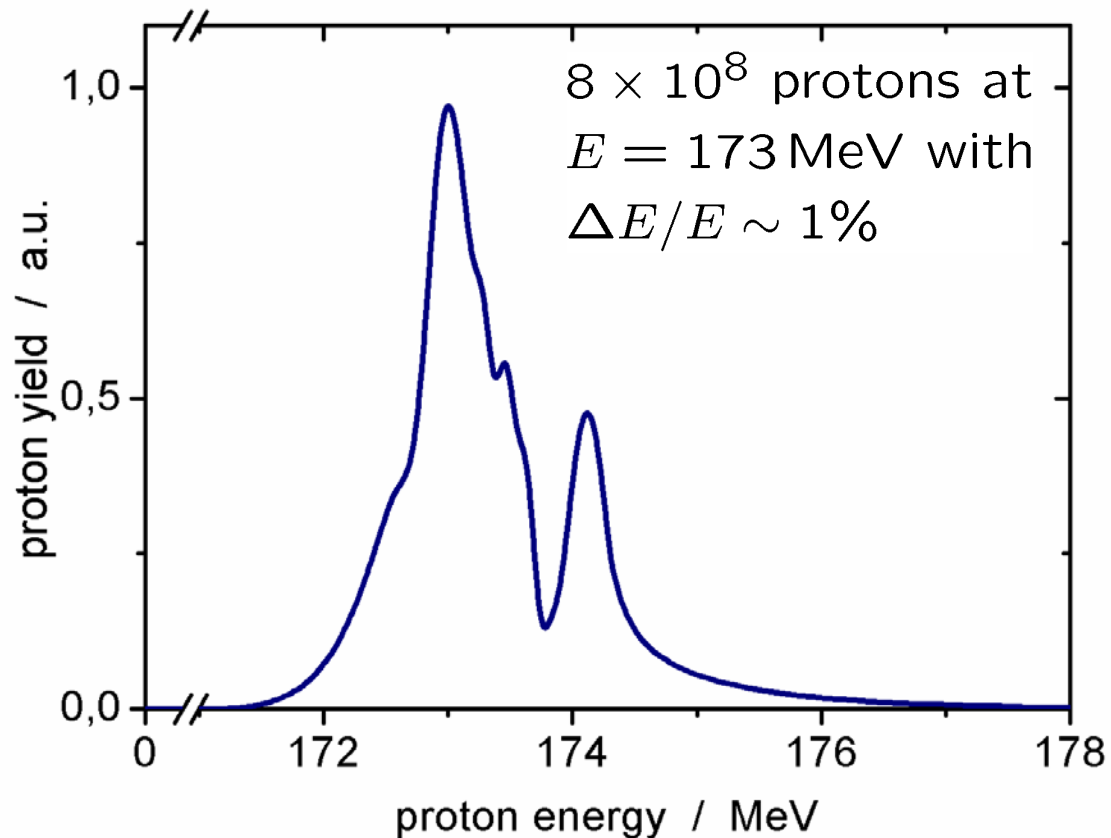




- overall number of ions about 10^8 in 20msr
- **80% fidelity with online target cleaning (ablation)**

[H. Schworer et al., Nature 439 (2006) 445]

2D-PIC simulation by T. Esirkepov for next laser generation (POLARIS):
100 J in 100 fs, $I_L = 10^{21}$ W/cm², 5 μ m Ti-foil + 0.1 μ m PMMA dot (\varnothing 2.5 μ m)



Features of laser accelerated beams

high charge (up to nC)
short pulses (down to 10fs)



high peak current
(up to 100kA)



space charge



„compact“ accelerators (up to GeV electrons)



low rep-rate (<10Hz)



excellent emittance, yet poor divergence

Applications of laser-accelerated particle beams

- Ultra-short X-ray pulses
 - Free electron lasers, Thomson back scattering
- Nuclear and particle physics
- Everything you can do with conventional accelerators ?
- Medical applications
 - Radiation therapy and imaging

Laser accelerators vs conventional accelerators

Available av. laser power

Average power (on target / stored)

CO₂ 100 kW

fs PW class 100 W

Diode pumped PW ...few kW

Acc. Efficiency ~10%

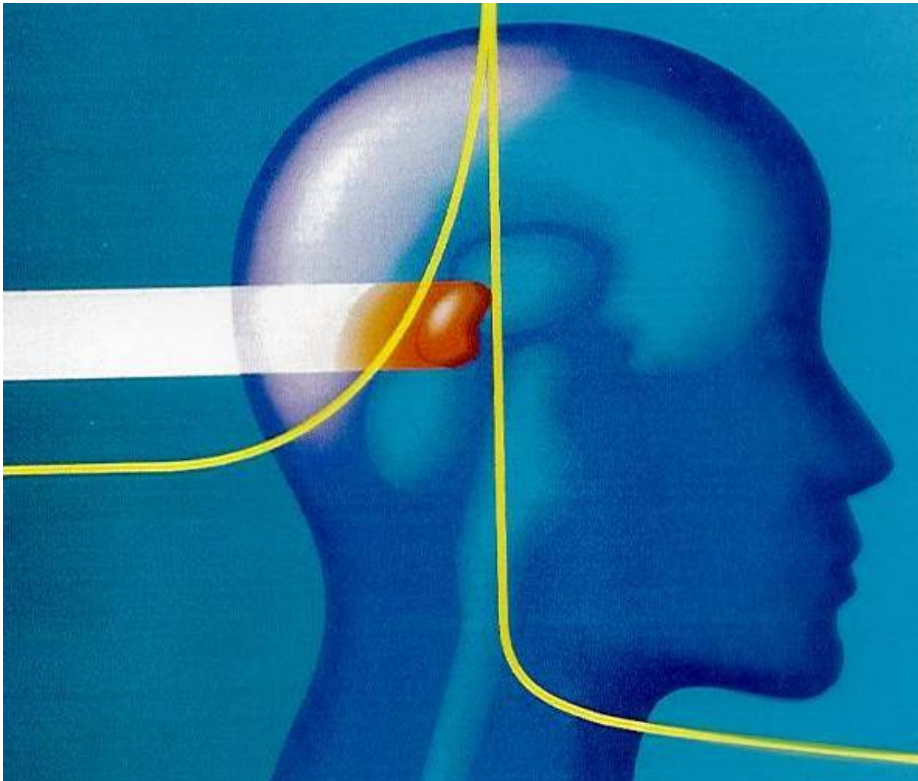
ELBE 10 kW (40 MeV e)

SNS 1.4 MW (GeV p)

LHC 350 MJ (7 TeV p)

SIS100 50 MJ (20 AGeV U)

Laser driven ion (proton) beam therapy ?



© GSI Darmstadt

requirements for ion beam therapy

Dose: 40-80 Gray distributed over 10-20 fractions

-> 10^9 - 10^{10} ions per fraction and few minutes

Spatial control: mm-scale @ 20cm depth

-> 200 MeV @ percent level control

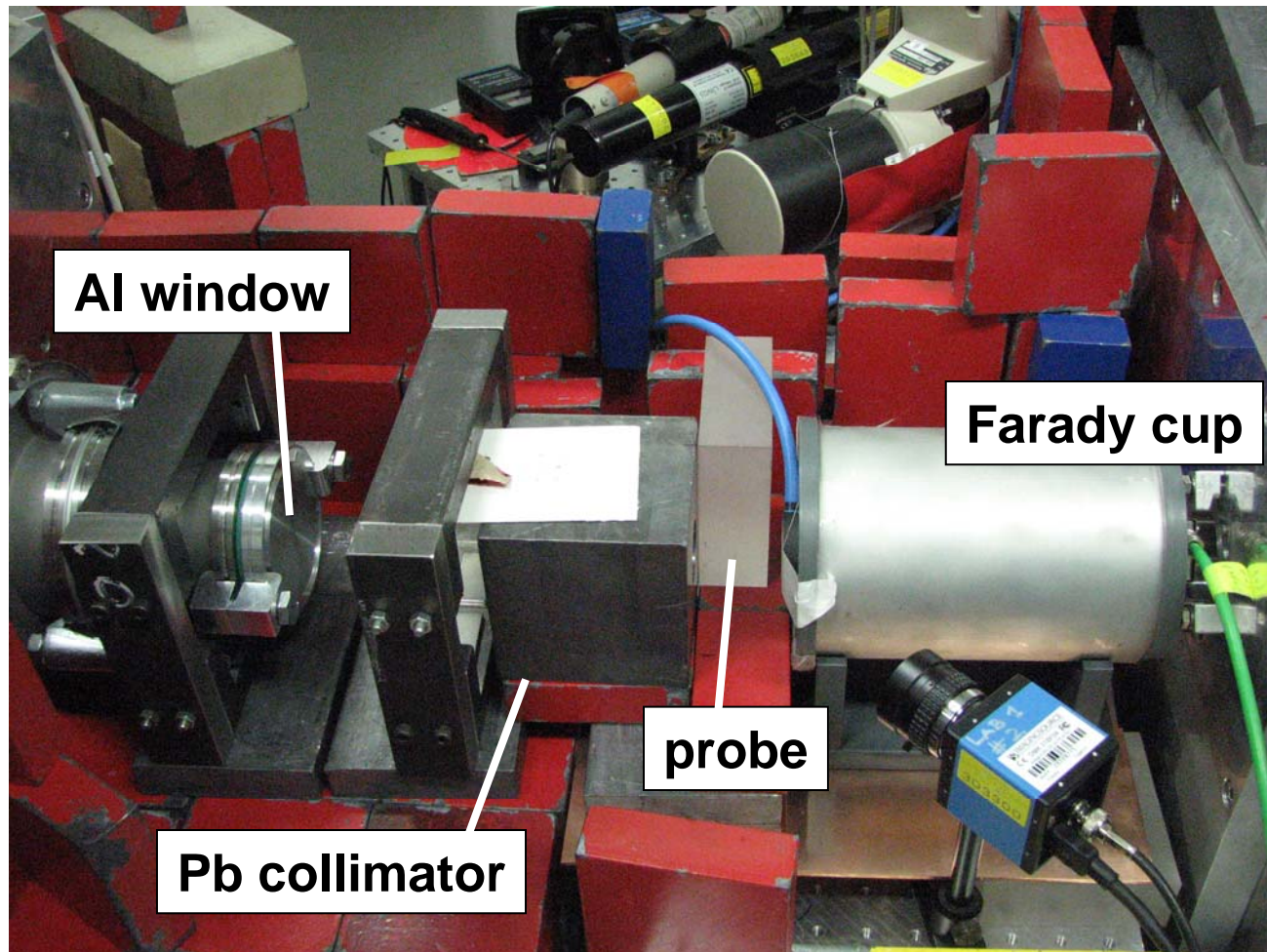
-> mm pointing (contour shaping)

-> 5% position dependent dose control

Complete (nondestructive shot-to-shot) monitoring

Clean beam (no other species, X-rays...)

Two cell lines were irradiated with doses of ~ 3 Gy of laser accelerated electrons (undefined spectrum...)

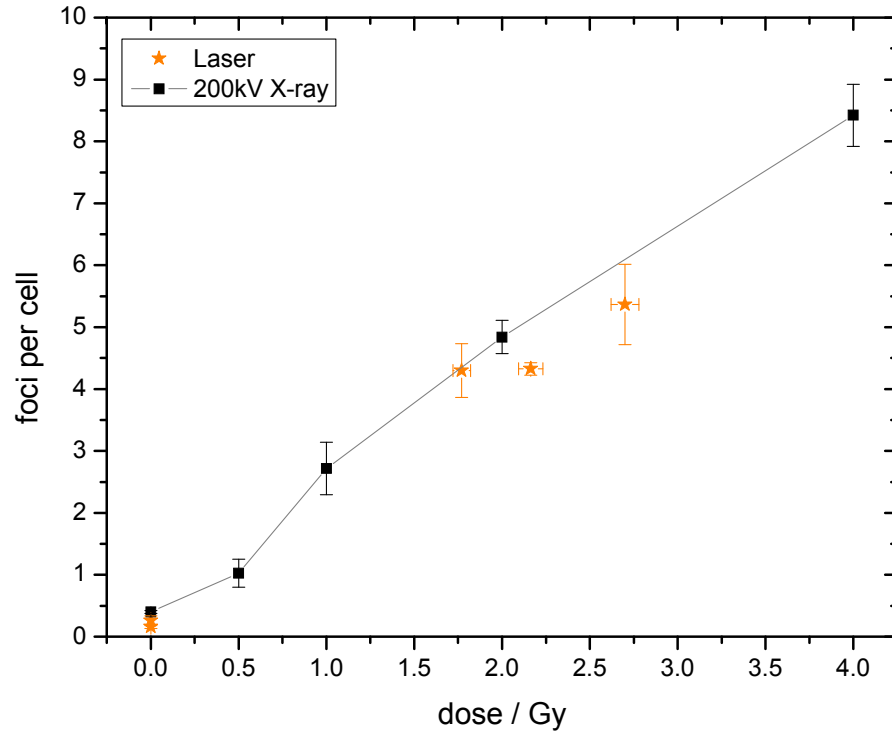


Detection of DNA damage (off line)

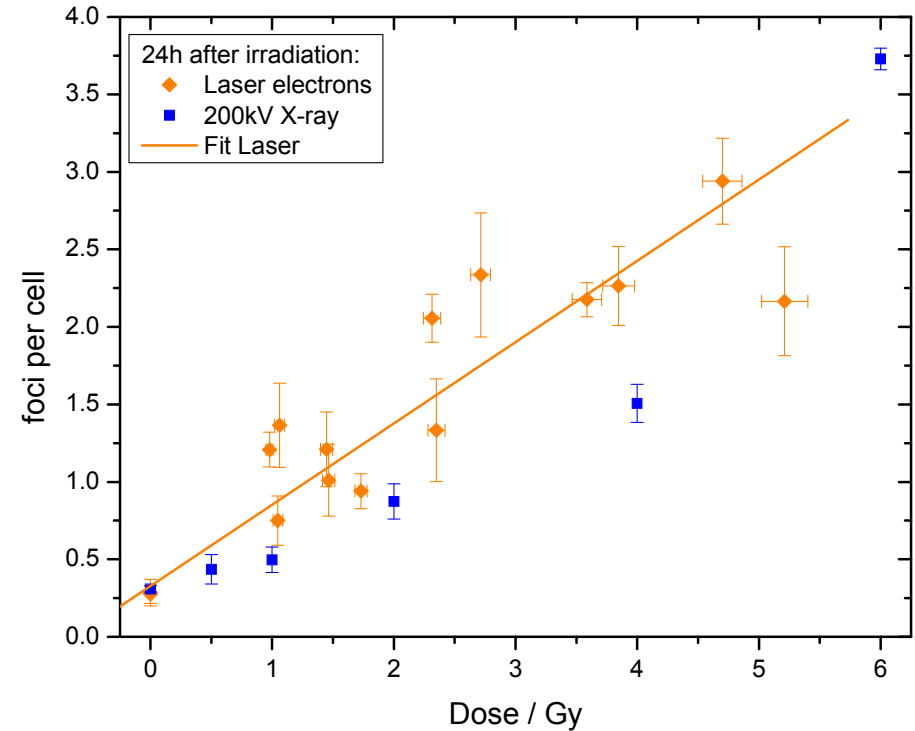


Probe 1, 184A1

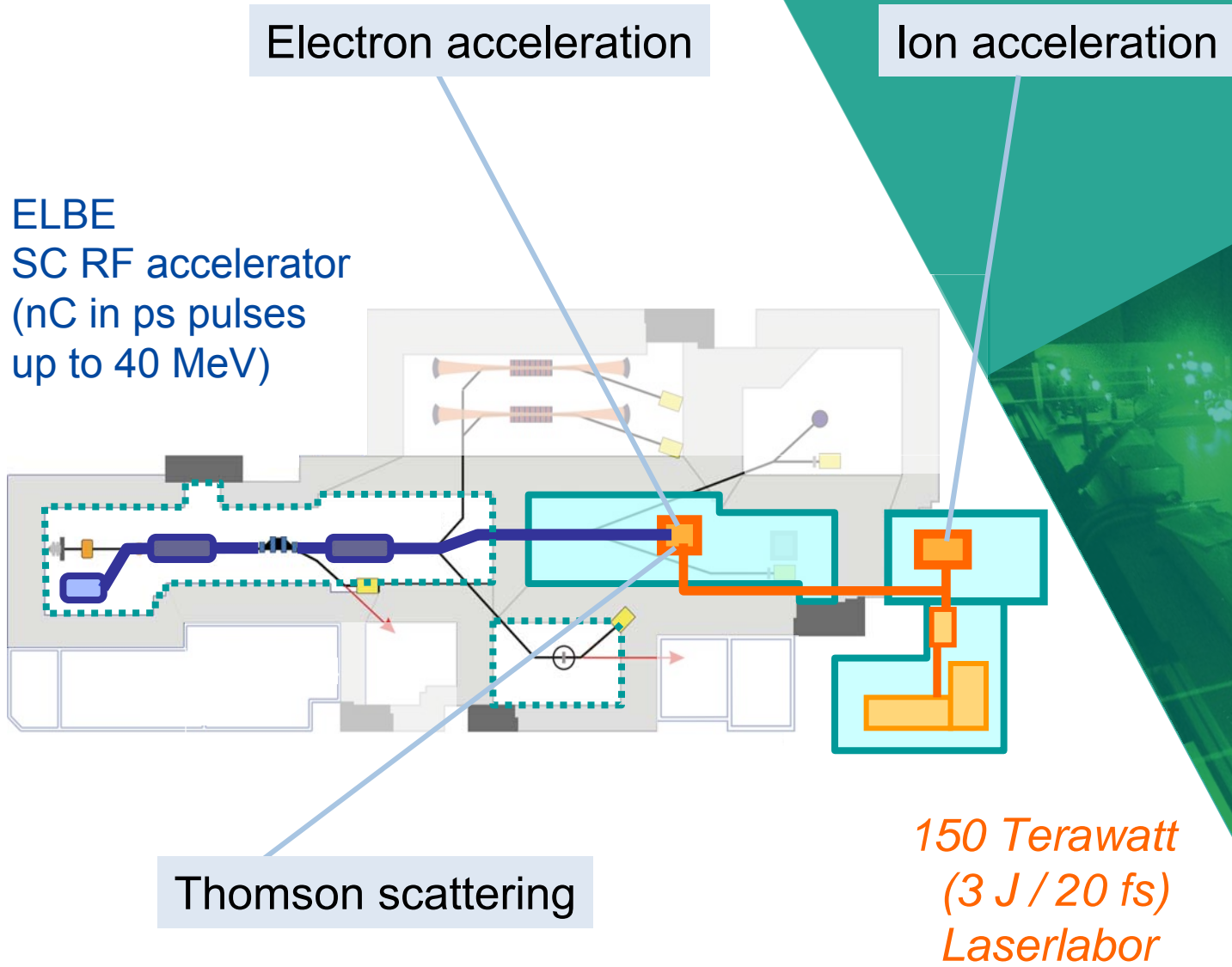
2h after irradiation



24h after irradiation



184A1 (humane Brustdrüsenepithelzellen, normal)





U. Schramm, A. Debus, T. Kluge, S. Kraft, K. Zeil, S. Bock



K. Ledingham



H. Schworer, B. Liesfeld, K.-U. Amthor, W. Ziegler, O. Jäckel,
S. Pfotenhauer, S. Podleska, R. Bödefeld, J. Hein, J. Polz, F.
Ronneberger, H.-P. Schlenvoigt, B. Beleites

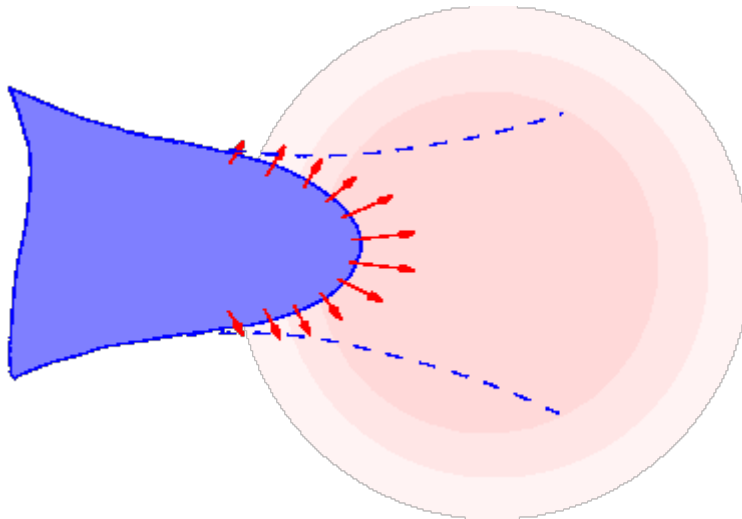


E. Beyreuther, L. Karsch, J. Pawelke, W. Enghardt, M. Baumann



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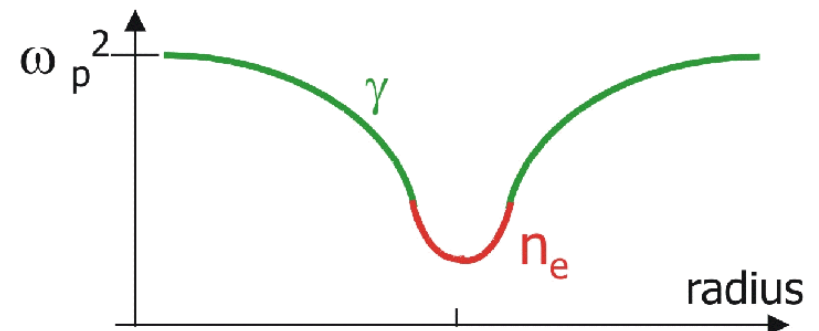
target: underdense (=transparent) plasma



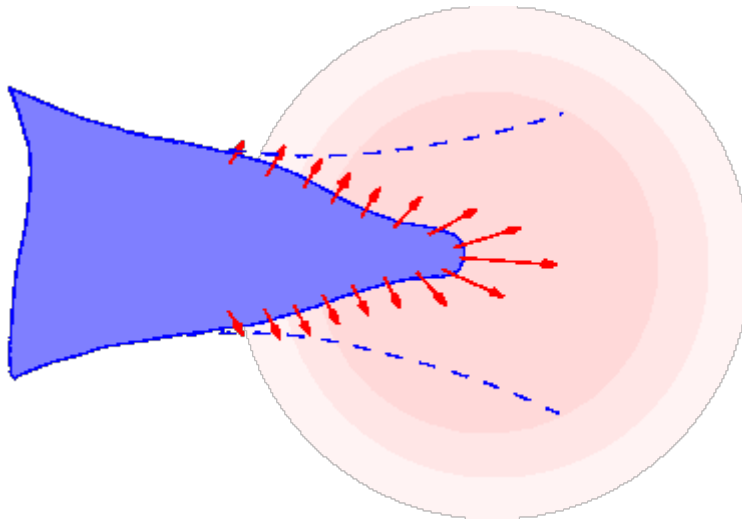
plasma frequency

$$\omega_p^2 = \frac{en_e}{\epsilon_0 \gamma m_e}$$

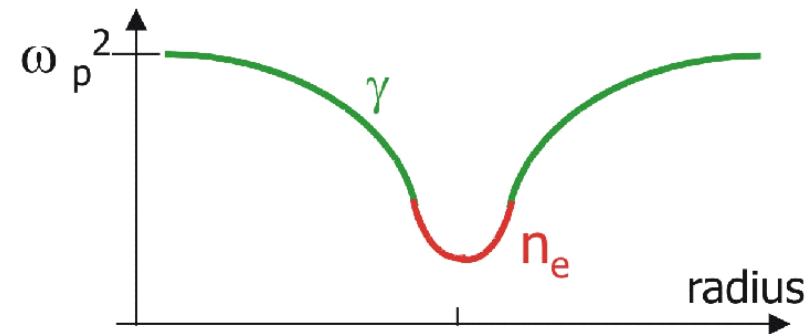
- **mass increase**
- **density reduction**



Index of refraction n locally increases -> **relativistic self focusing**

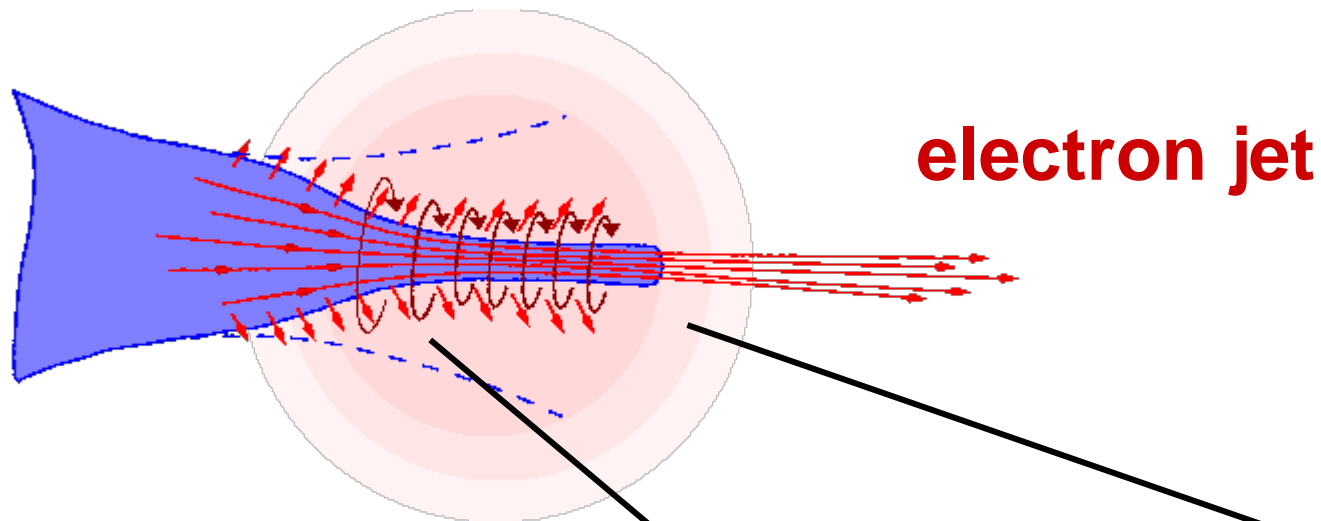


$$n = \sqrt{1 - \frac{\omega_p^2}{\omega^2}}$$

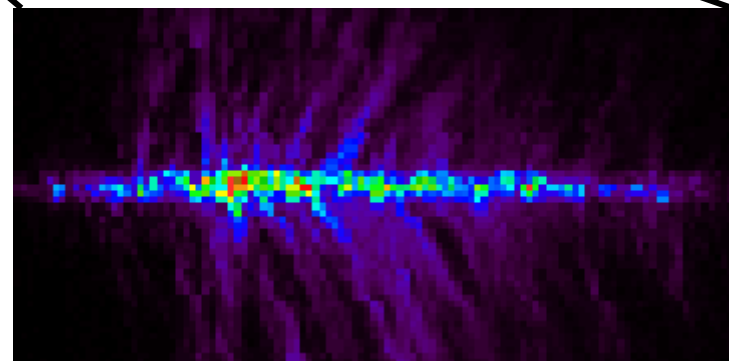


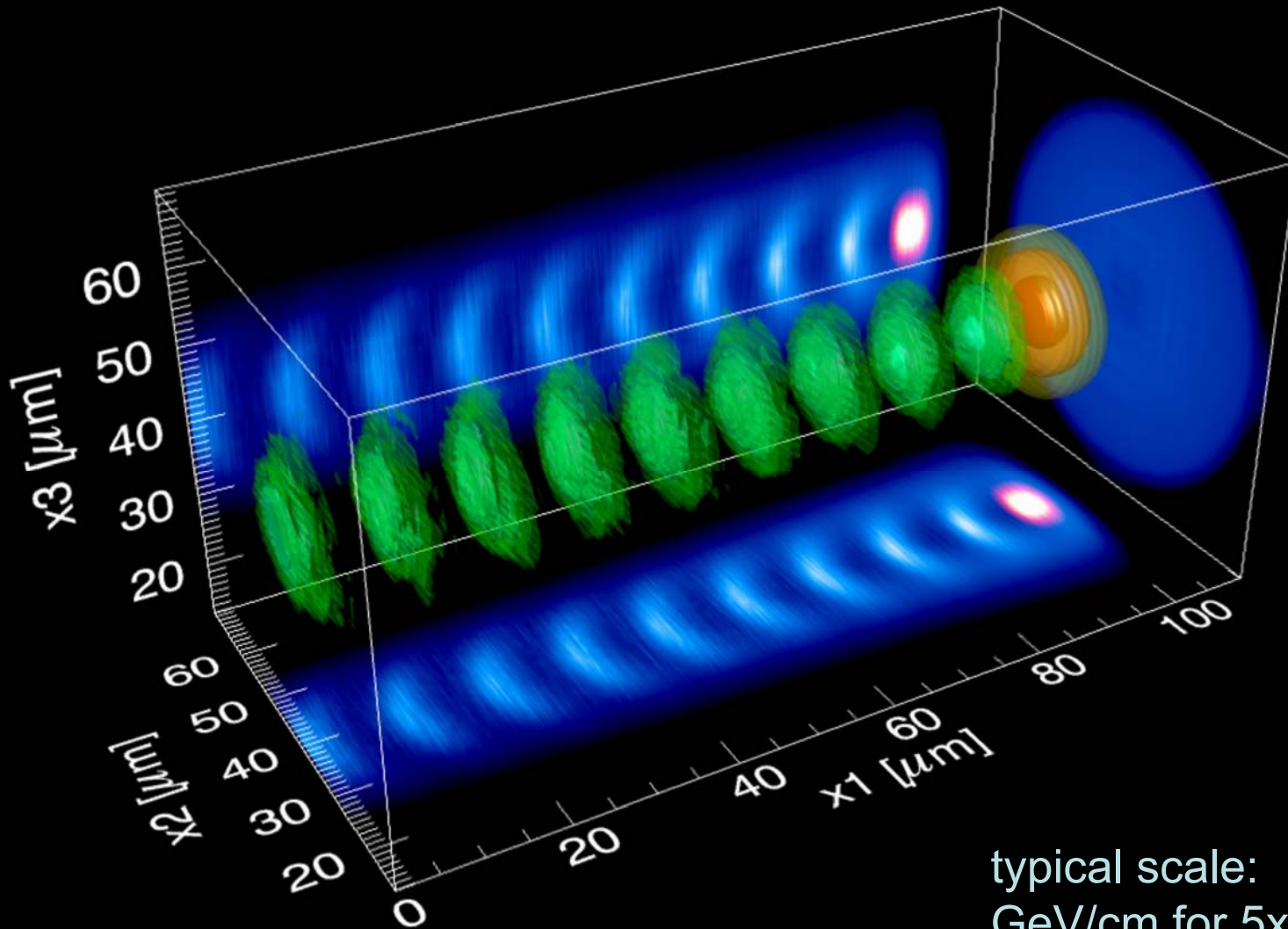
$$v_{ph} = c/n \quad v_{gr} = c \cdot n \quad n_{19} = 0.999 \quad n_0 = 0.995$$

channel formation



Thomson scattering @ 2ω
(length $0.4\text{mm} = 400 \lambda$)

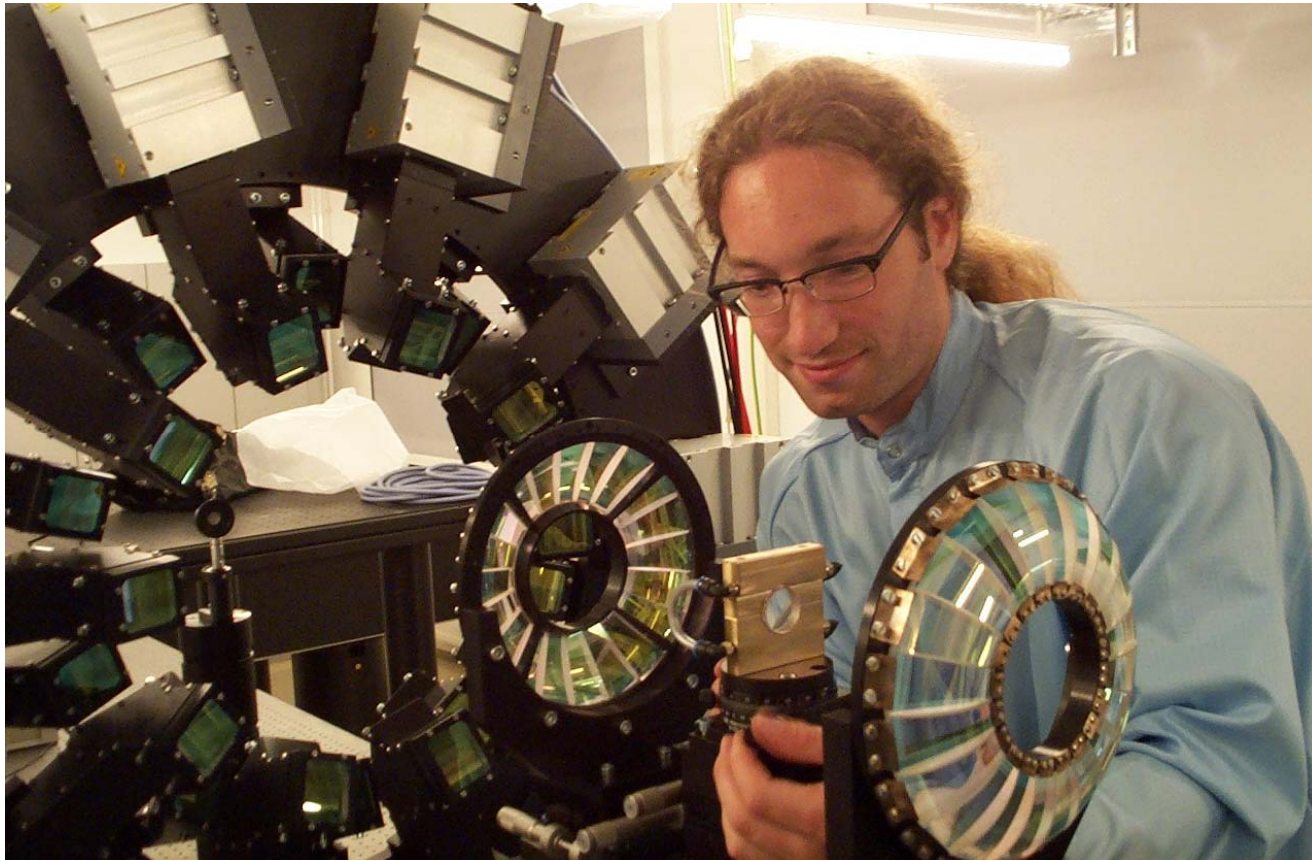


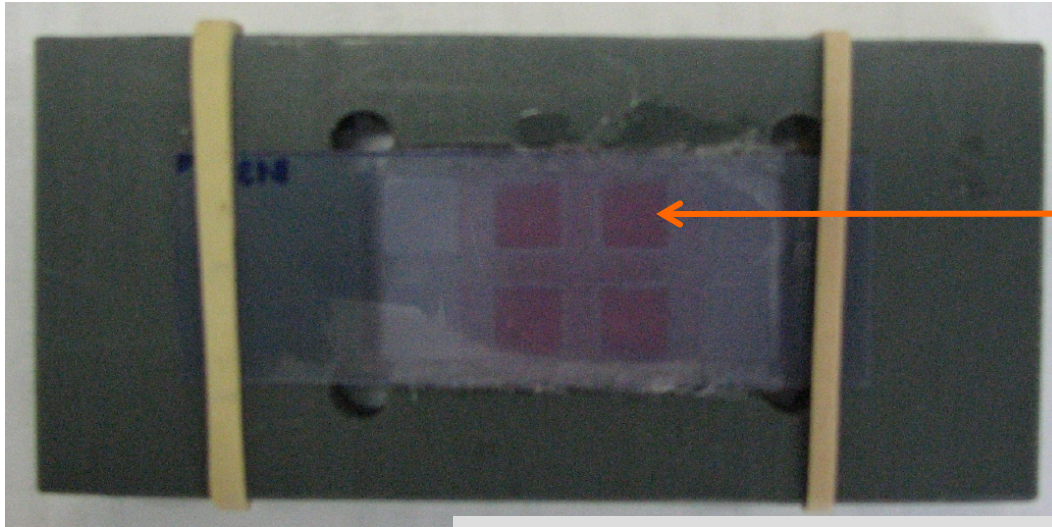


typical scale:
GeV/cm for 5×10^{18} e/cm³

3D PIC Simulation courtesy L. Silva, W. Mori

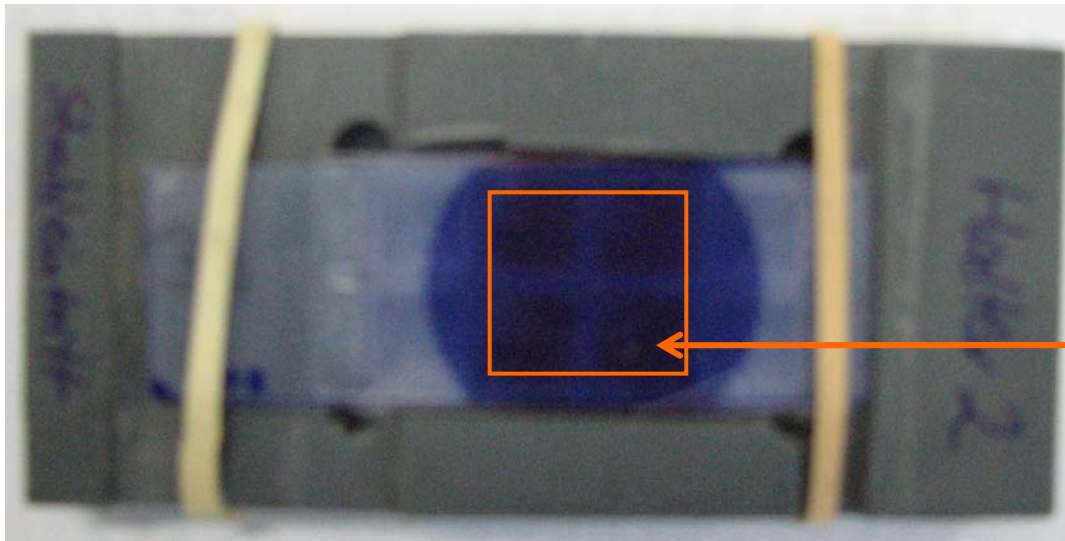
2D-PIC simulation by T. Esirkepov for next laser generation (POLARIS):
100 J in 100 fs, $I_L = 10^{21}$ W/cm², 5 μ m Ti-foil + 0.1 μ m PMMA dot (\varnothing 2.5 μ m)





Rückseite, vor Bestrahlung

Mittlere 4 Kammern des Objektträgers mit Zellen bedeckt

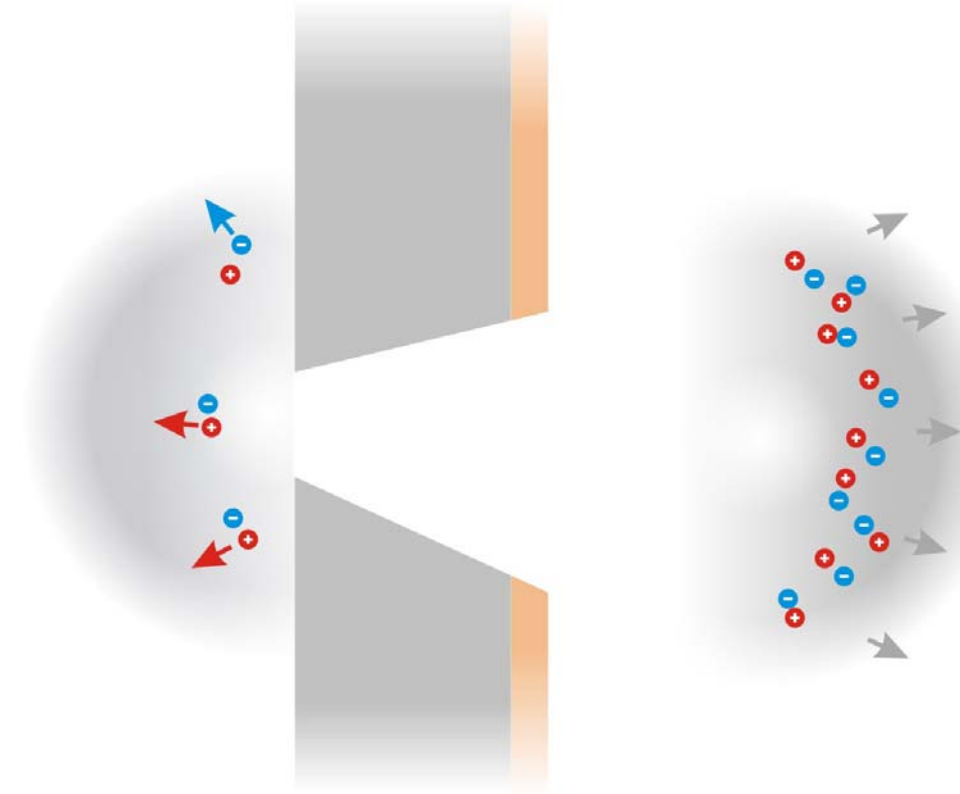


Vorderseite, nach Bestrahlung

Strahlfleck: 35mm

Ausgewerteter Bereich:
4 Kammern à 50 Zellen

- **electron acceleration**
- **hot (MeV) electrons penetrate the foil**
- **quasi static field forms normal to target surface, source size \gg laser spot**



quasi-neutral pulse with exponential energy distribution (with max. energy depending on laser pulse duration, energy, and target thickness)

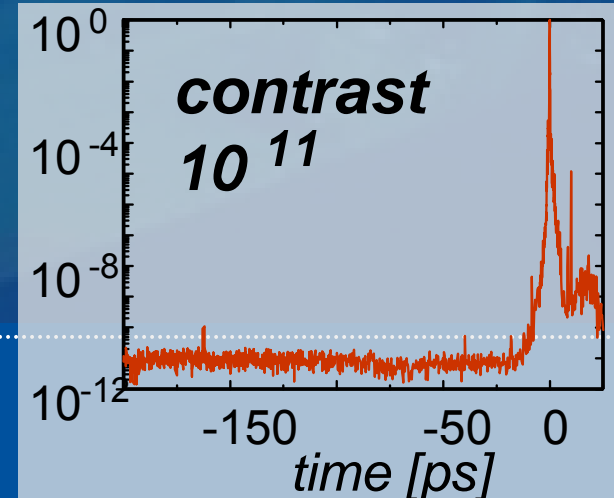
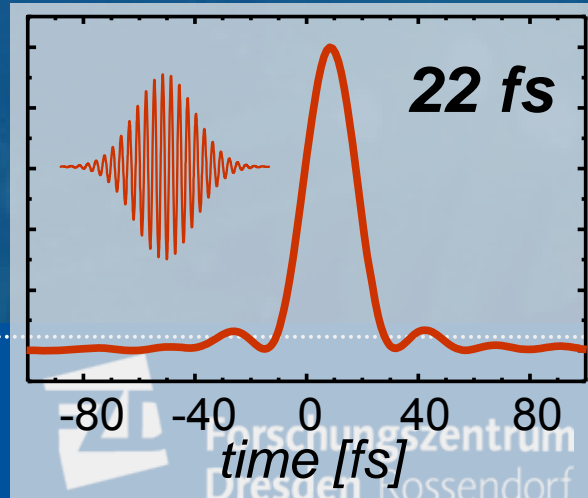
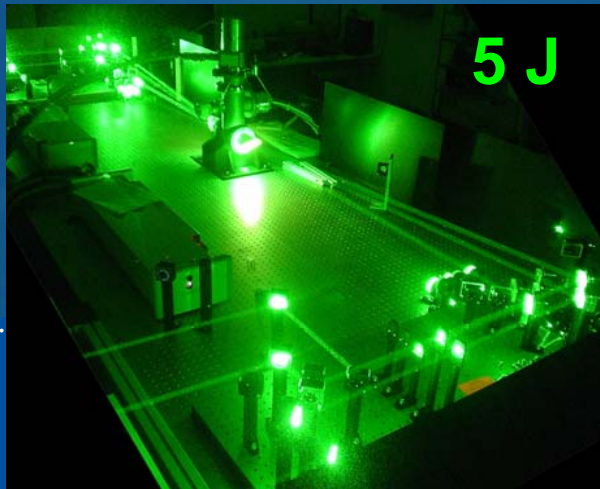
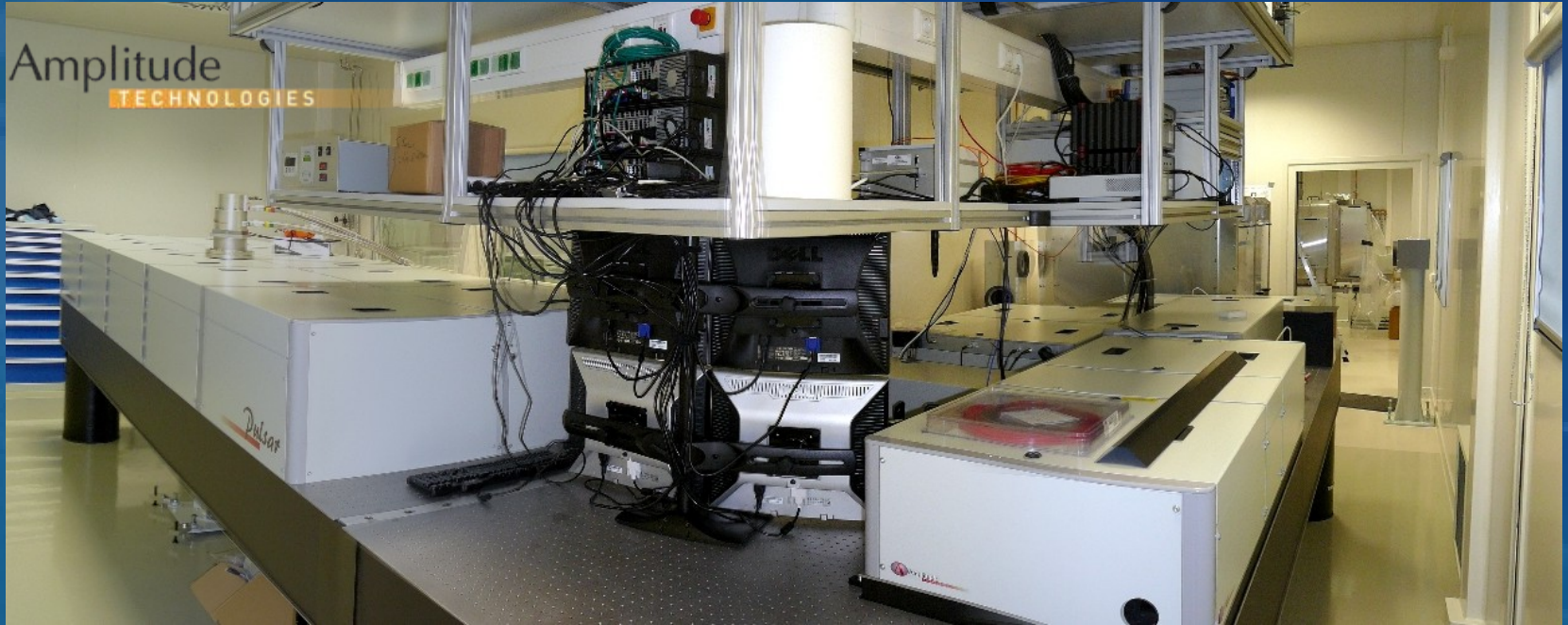






Amplitude

TECHNOLOGIES



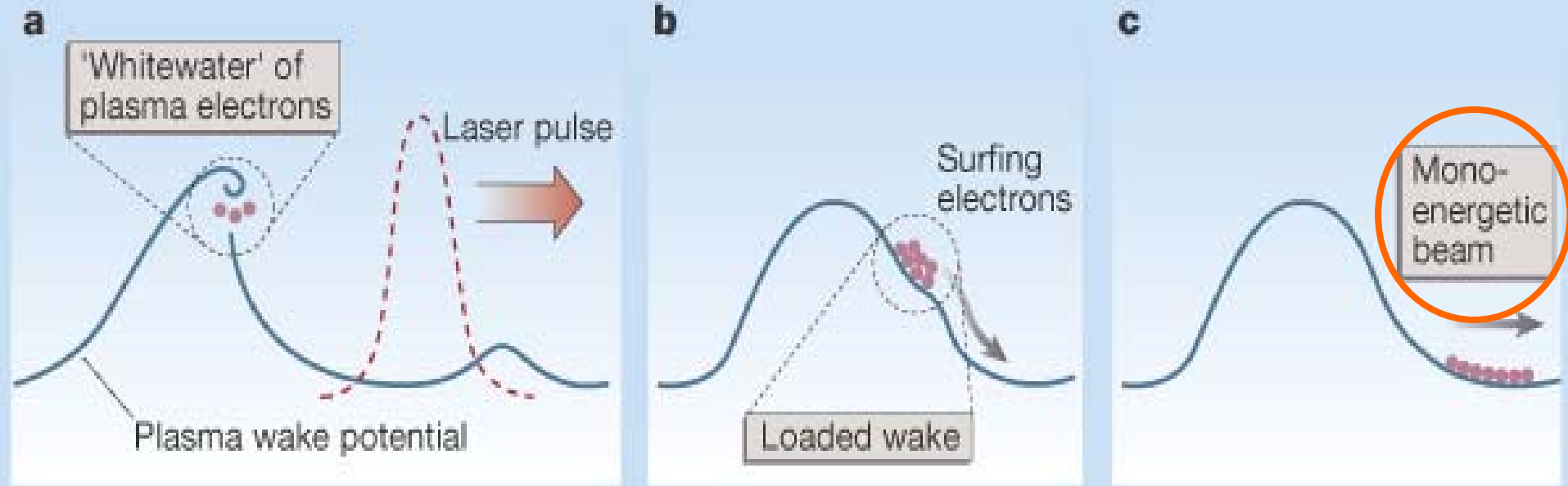
the FZD group (FWT)



R. Sauerbrey, U. Schramm

T. Kluge, S. Kraft, K. Zeil, S. Bock, (M. Bussmann, M. Siebold, A. Debus)

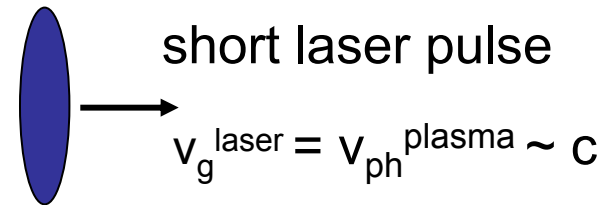
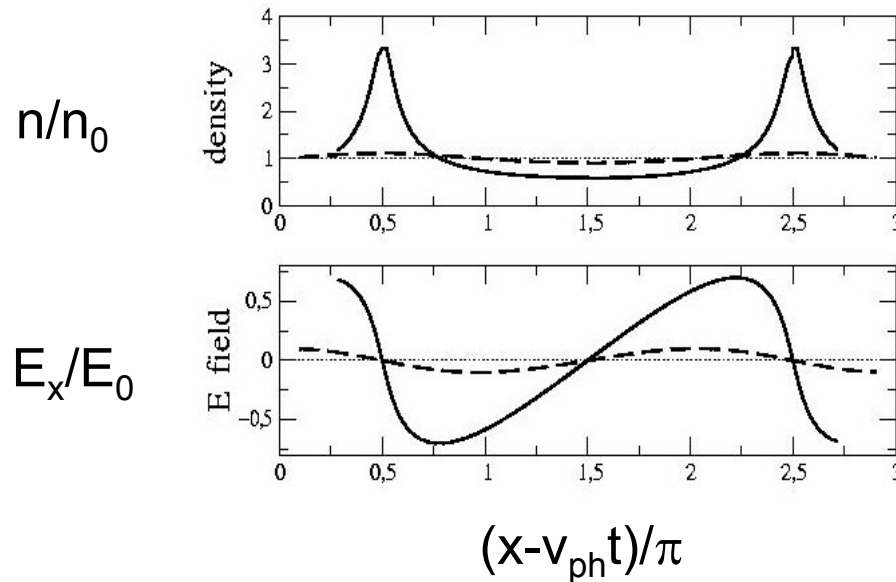




T. Katsouleas, Nature 431, 515 (2004)

... and yet surprisingly „real“ in the highly nonlinear broken wave – blow-out – bubble regime ...

In a transparent plasma a relativistic laser pulse with $L < \lambda_p = c/\omega_p$ drives a longitudinal plasma wave



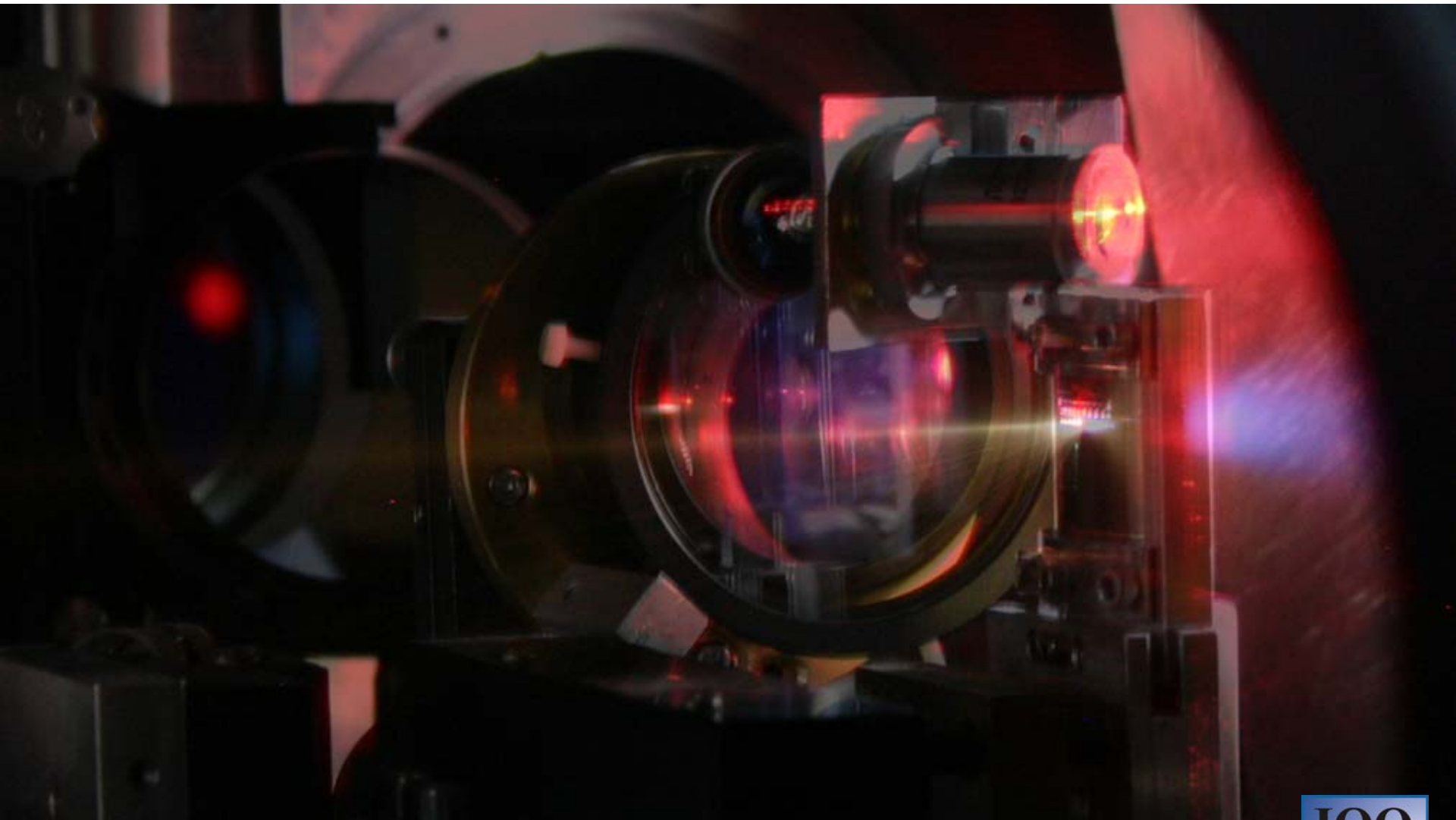
non-linear plasma wave for relativistic intensities

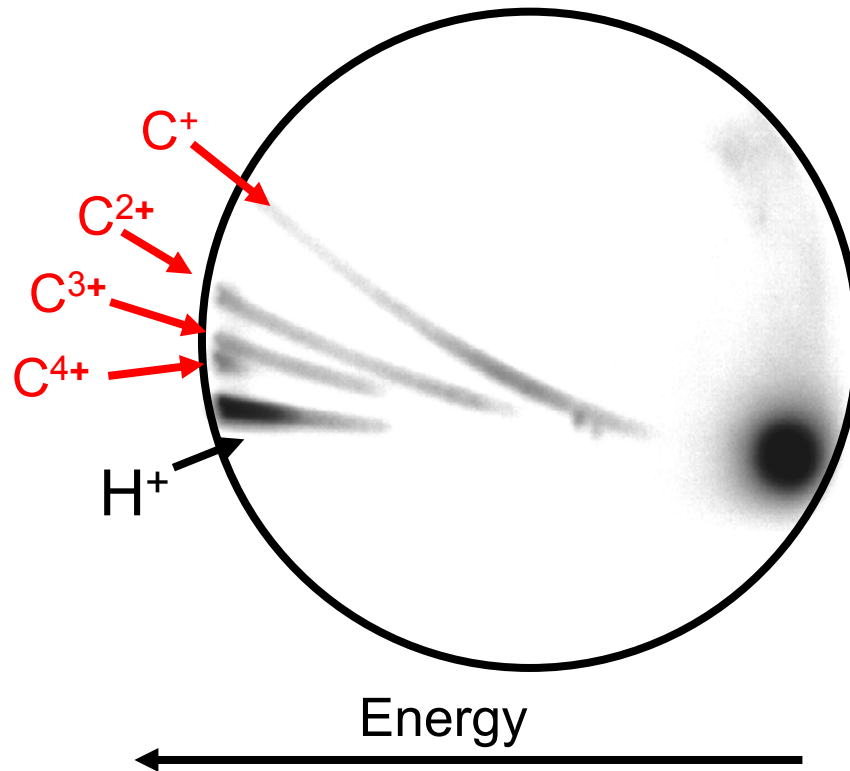
acceleration of particles in (traveling) waves

-> matched external injection or self-injection

Typical (10TW) laser parameters are 1J in 40-80fs

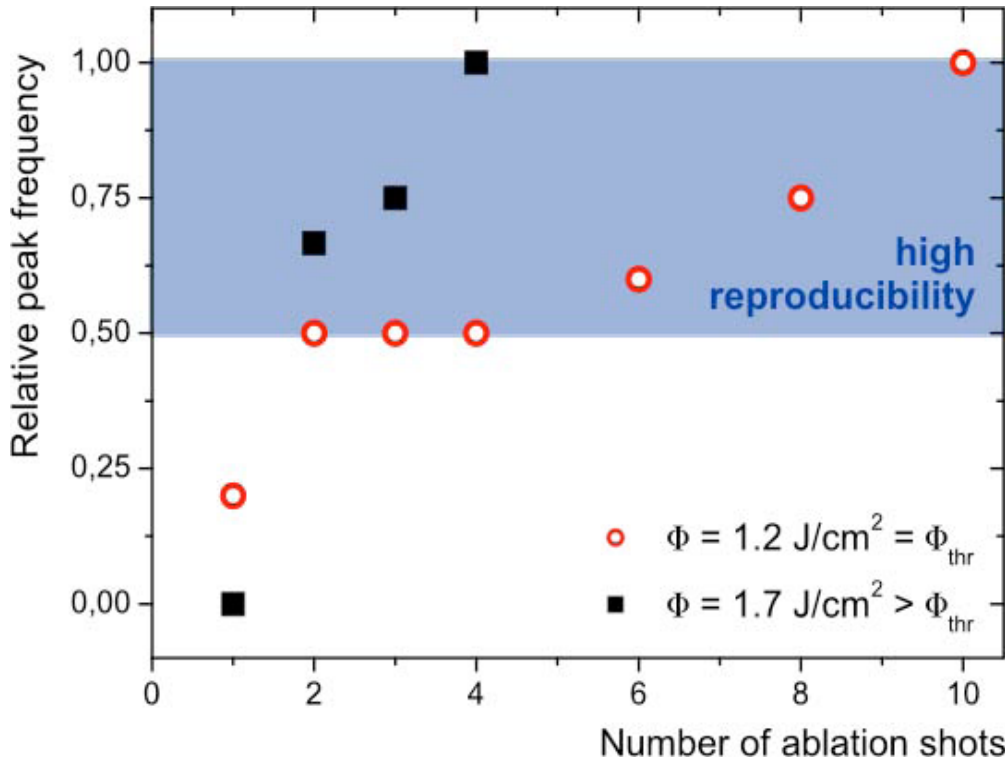
- too long for high density (underdense) plasma
- too weak to use much less density
(as threshold scales as $\omega_p^{-4/3}$ or $n_e^{-2/3}$ and including the larger spot size as $n_e^{-13/6}$)
- ▶ one might use capillary guiding for an increased acceleration distance, where lower densities are favourable and
- ▶ rely on self-compression and -focusing (with all its instability problems)





Ions are separated by a mass and energy selective spectrometer and hit a position sensitive detector.

All species of ions located at the backside of the foil can be accelerated



80% fidelity

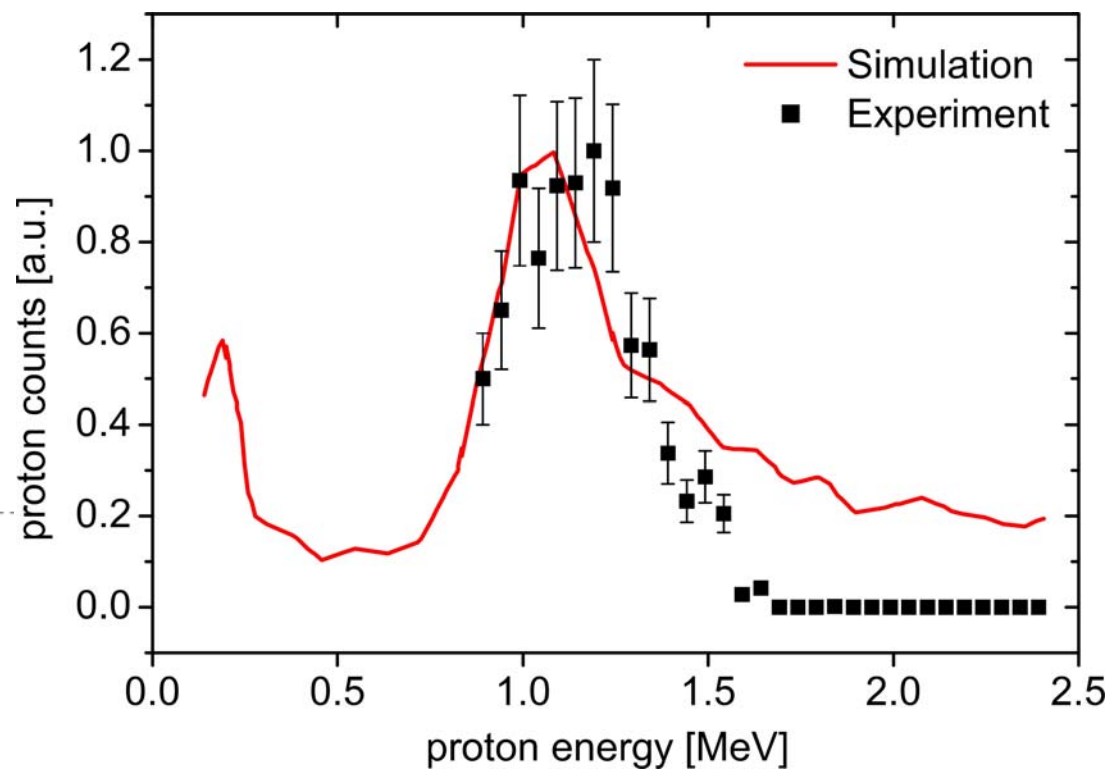
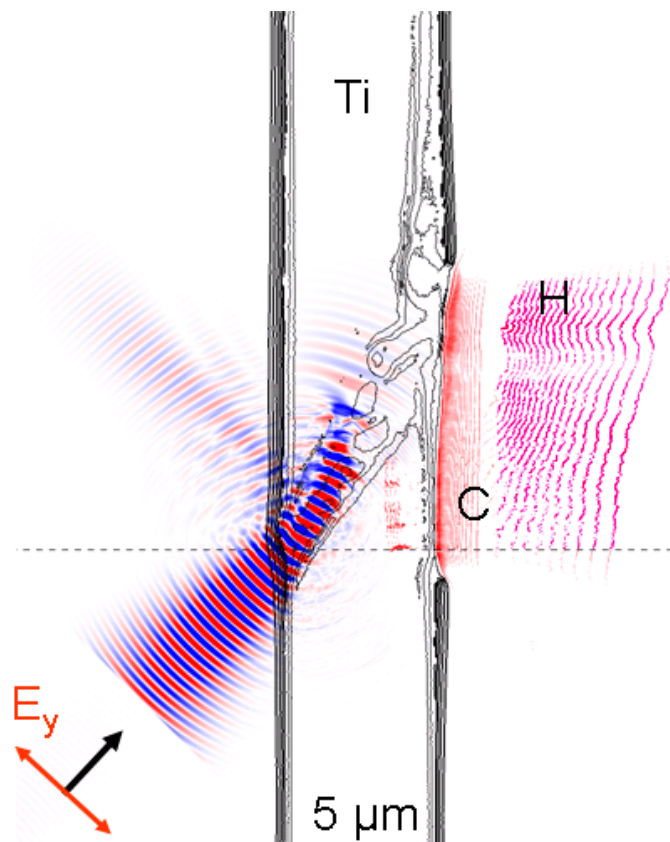
- Existence of threshold fluence » $1,2 \text{ J / cm}^2$ @ $t_{\text{pulse}} \gg 5 \text{ ns}$, $\lambda = 532 \text{ nm}$
- Observations of initial incubation effects
- Recombination time for adsorbants $> 5\text{s}$ at given chamber pressure of $p \gg 10^{-5} \text{ mbar}$

Simulation of Jena results (2006)

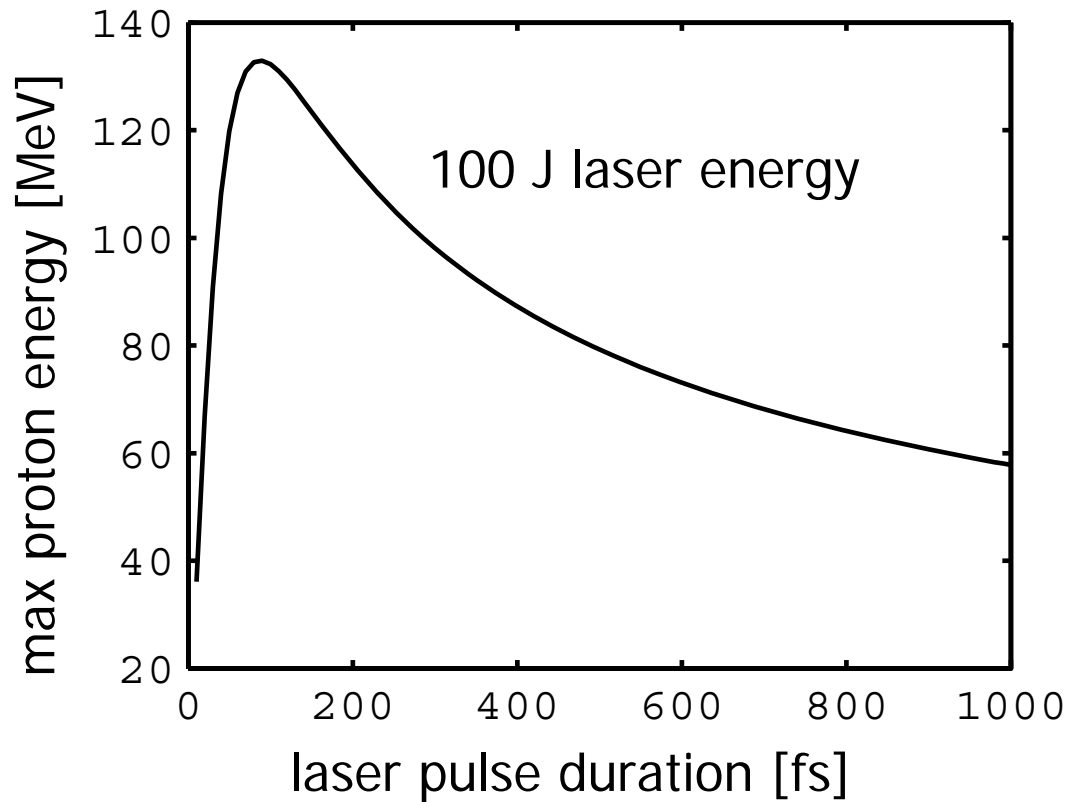


2D-PIC simulation by T. Esirkepov for following conditions :

$I_L = 3 \times 10^{19} \text{ W/cm}^2$, $5 \mu\text{m}$ Ti-foil + $0.5 \mu\text{m}$ PMMA dot ($20 \times 20 \mu\text{m}^2$)



For each laser system there is an optimal pulse duration for TNSA ion acceleration, which is not necessarily the shortest



[J. Schreiber, et al, PRL 97 (2006) 045005]

POLARIS laser system:

- Petawatt laser available in Jena by 2008 (diode pumped Yb³⁺:Glass)
- 4 out of 5 amplification stages realized including compressor (8 J, 150 fs)

$$I_{\text{POLARIS}} = 10^{21} \text{ W/cm}^2 \text{ @ } 0.1 \text{ Hz}$$

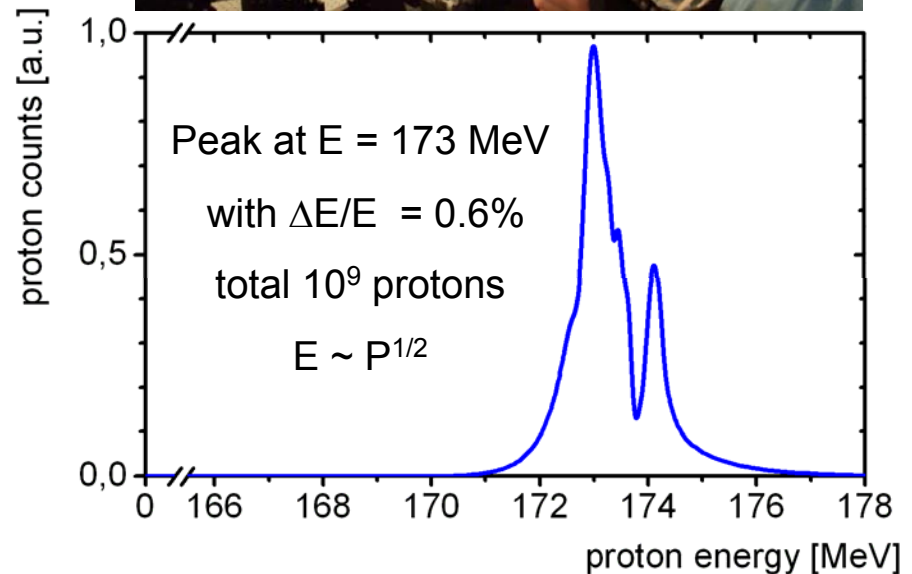
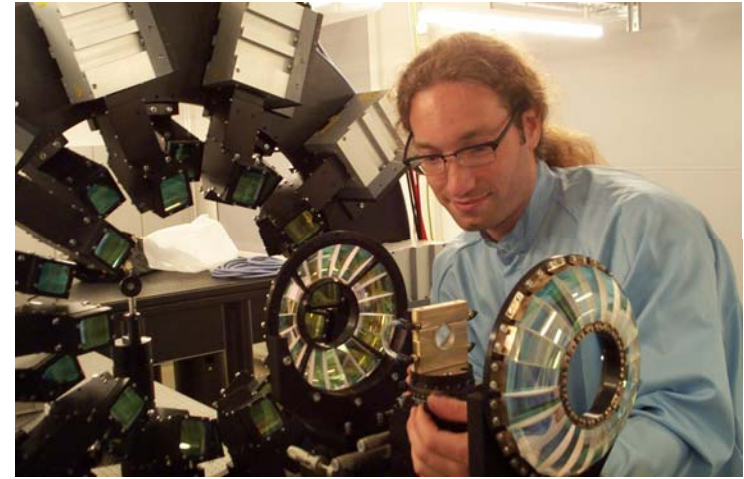
$$(E = 150 \text{ J}, \lambda = 1042 \text{ nm}, \tau = 150 \text{ fs})$$

POLARIS simulation: Diode pumped chirped pulse amplification to the joule level. *Applied Physics B - Lasers and Optics* 79, 1021 (2004)

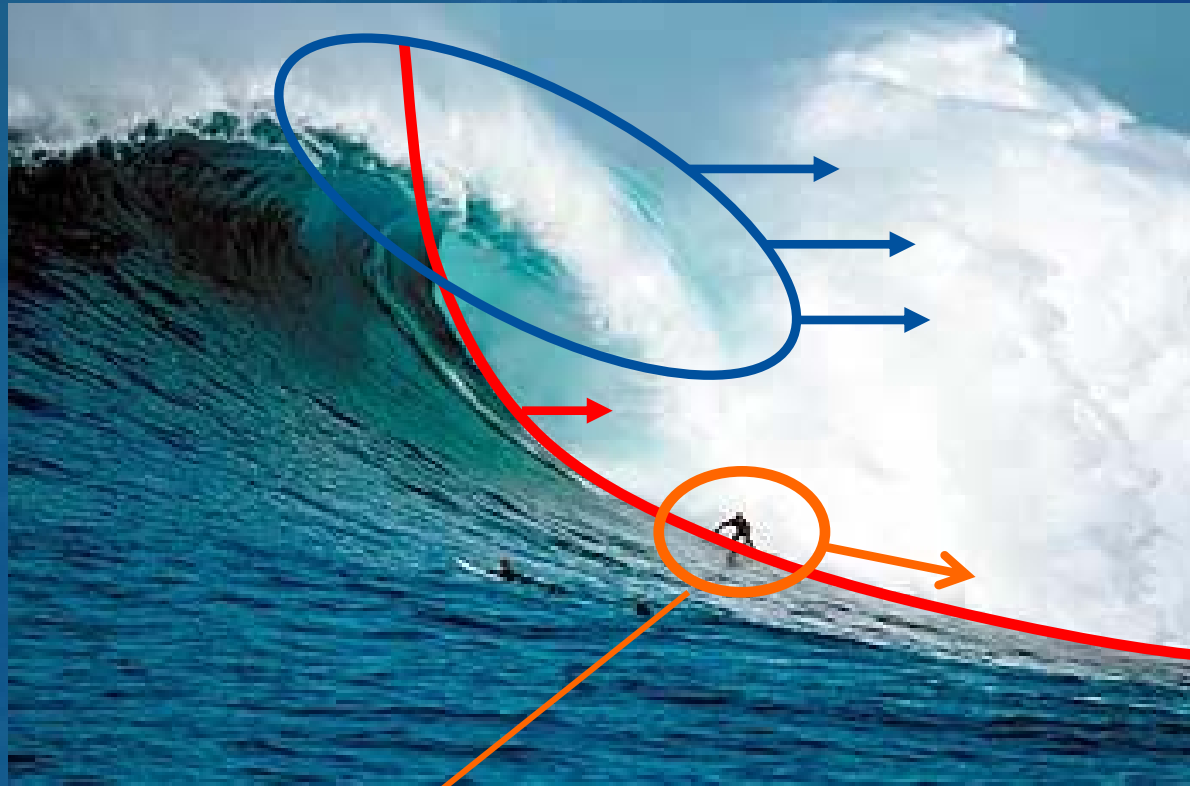
$I_{\text{POLARIS}} = 10^{21} \text{ W/cm}^2$

2,5 μm Ti-foil + 0.1 μm PMMA Dot (\varnothing 2.5 μm)

$\tau_{\text{ASE}} = 1 \text{ ns}$ @ $I_{\text{ASE}} / I_{\text{POLARIS}} = 10^{-7}$



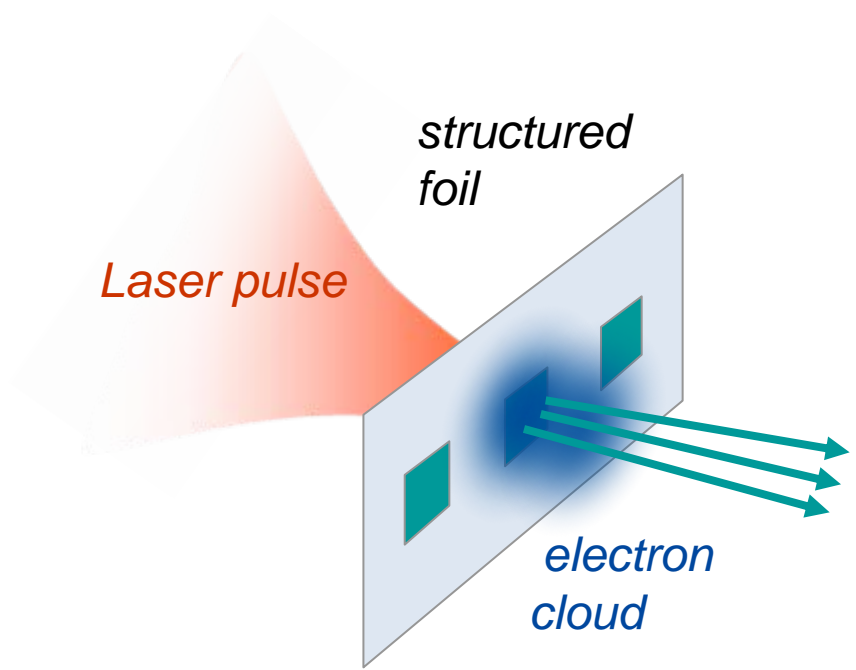
nonlinear wavebreaking (self injection) $v > v_{ph}$



dephasing

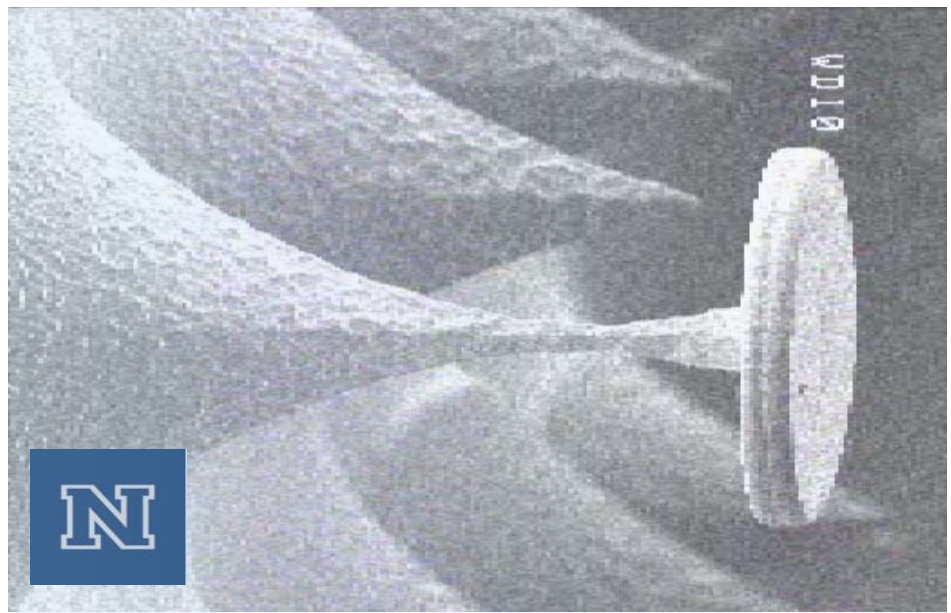
test particle $v_e > v_{ph}$
(external injection)

acceleration potential
(anharmonic, moving with v_{ph})



Mass-limited localized target (areas) lead to monoenergetic features

**More complex shapes might help
In guiding electrons ...**



Features of laser accelerated beams

high charge (up to nC)
short pulses (down to 10fs)



high peak current
(up to 100kA)



space charge



„compact“ accelerators - local machines

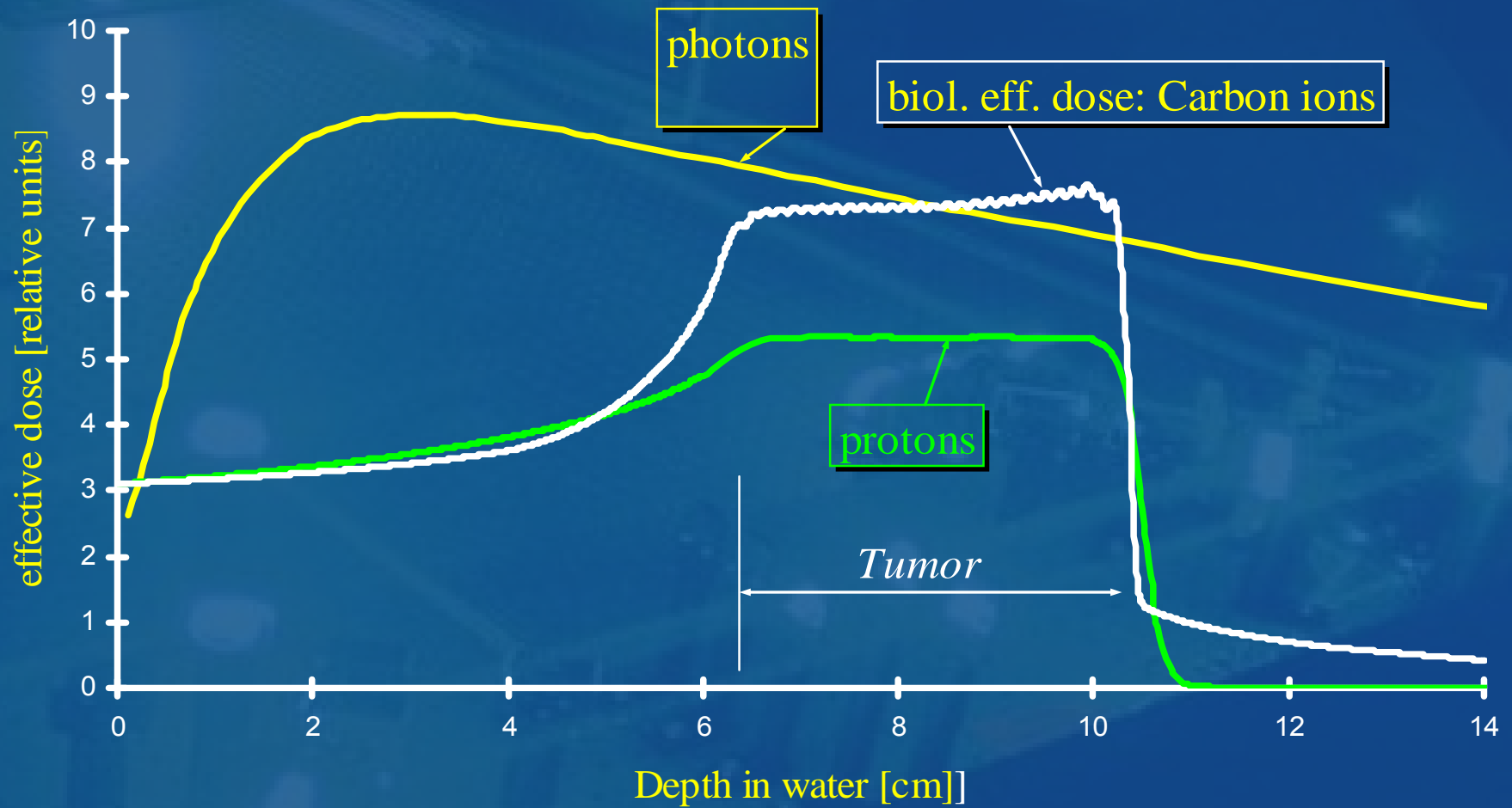


low rep-rate (<10Hz) - restrict to low average
yet high peak dose



excellent emittance, yet poor divergence

Ion beam therapy – the idea



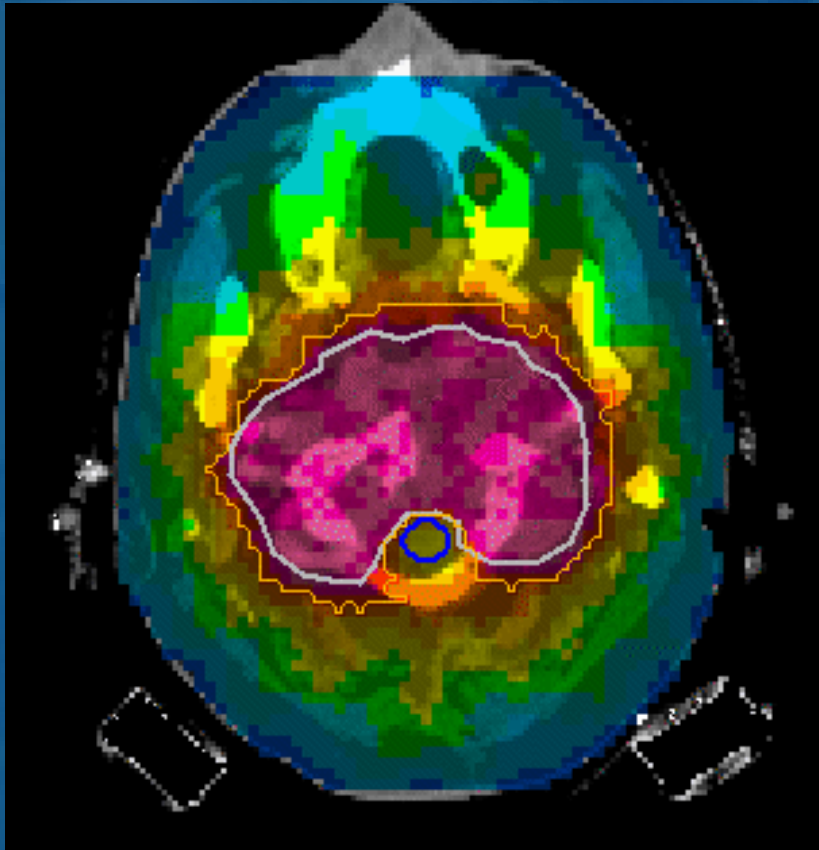
excitation of a longitudinal wave (wake)



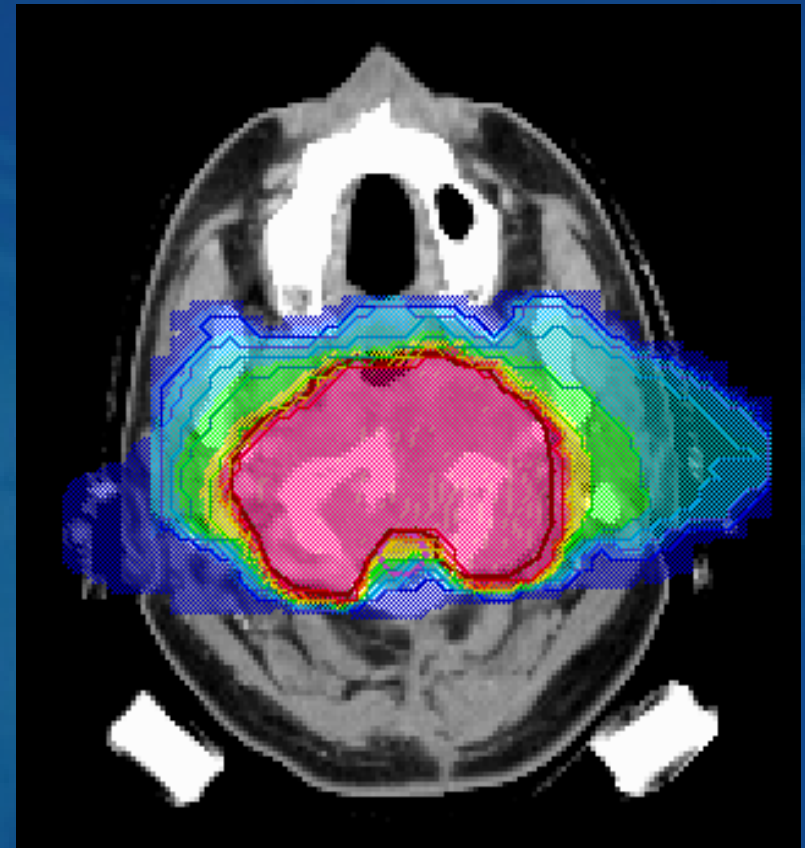
Ion beam therapy – treatment planning



Photons: 9 fields



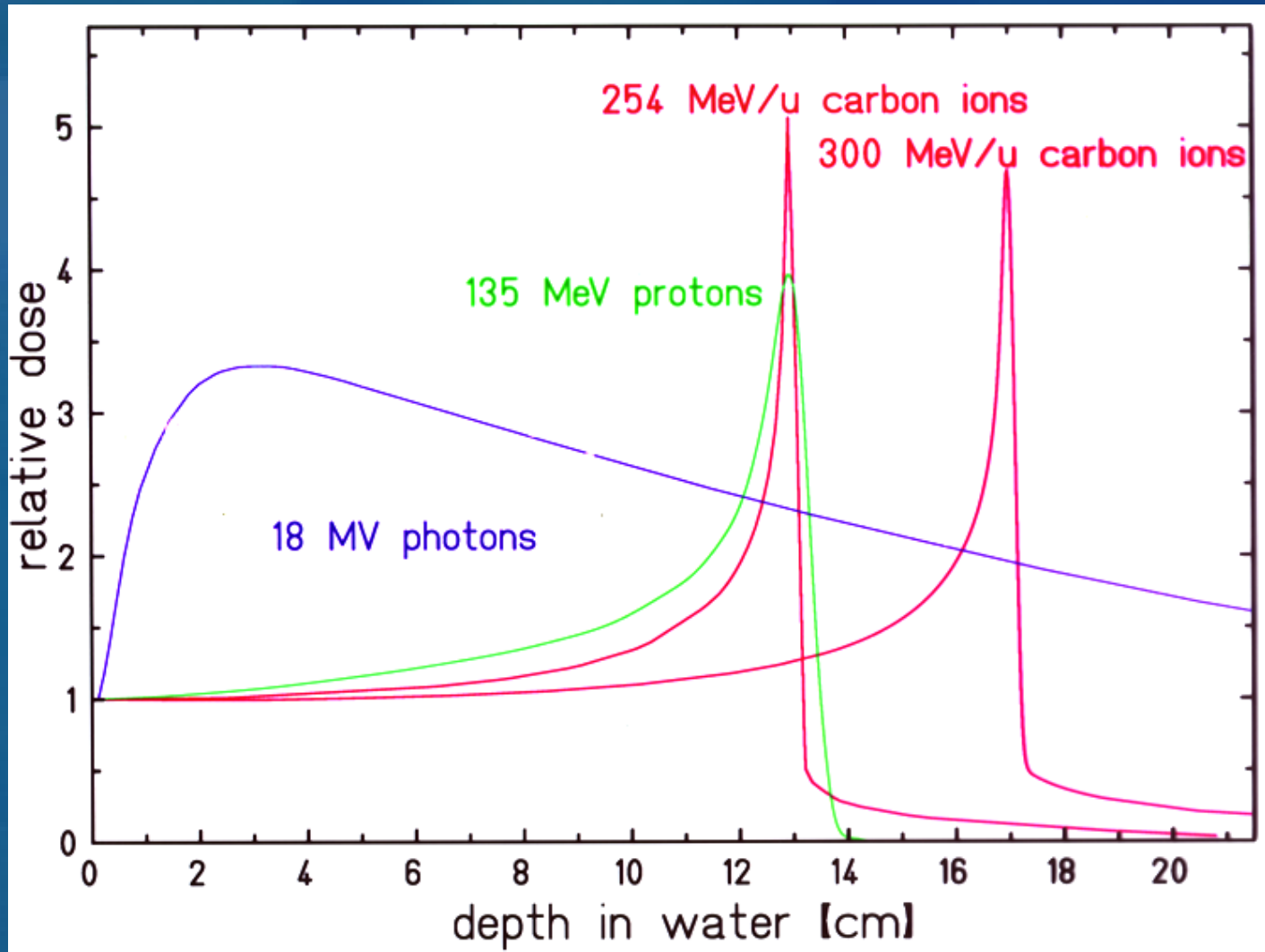
carbon-ions: 2 fields

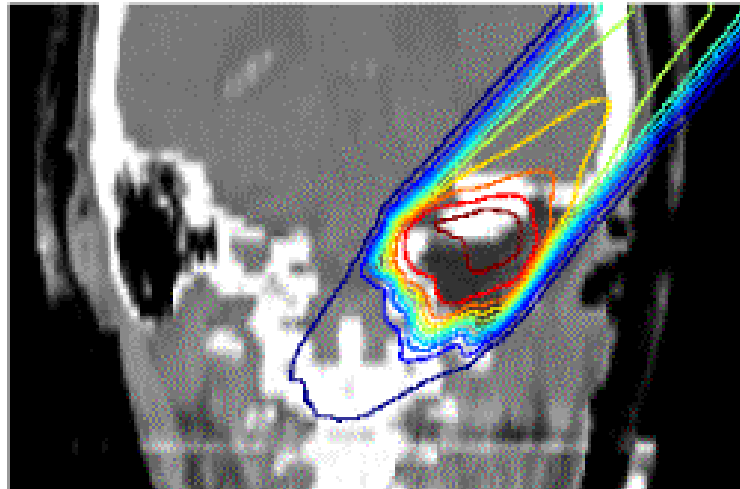


dose in % of the maximum dose

Courtesy O. Jäkel, DKFZ Heidelberg

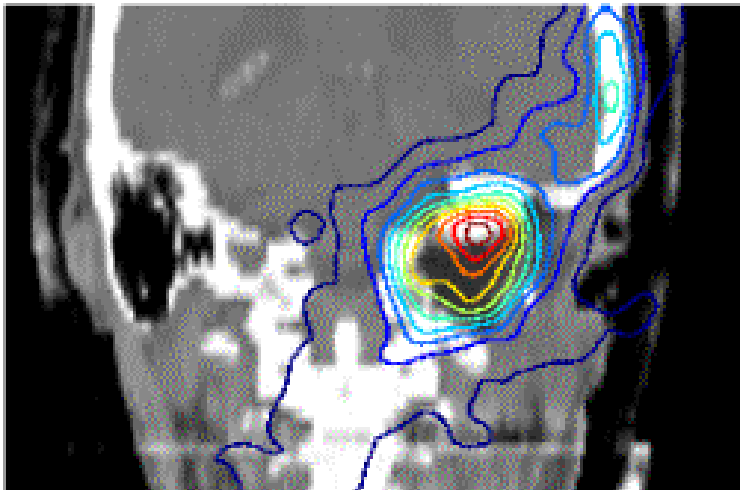
Ion beam therapy – required energies



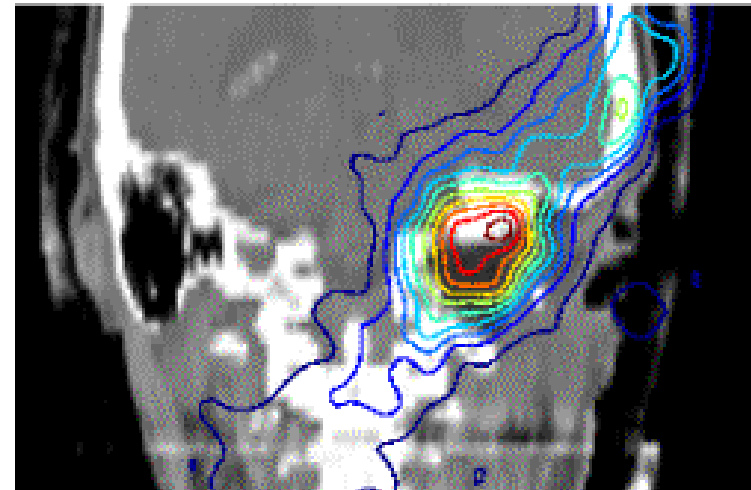


dose plan

W.Enghardt et al. ,
FZD Dresden



measured

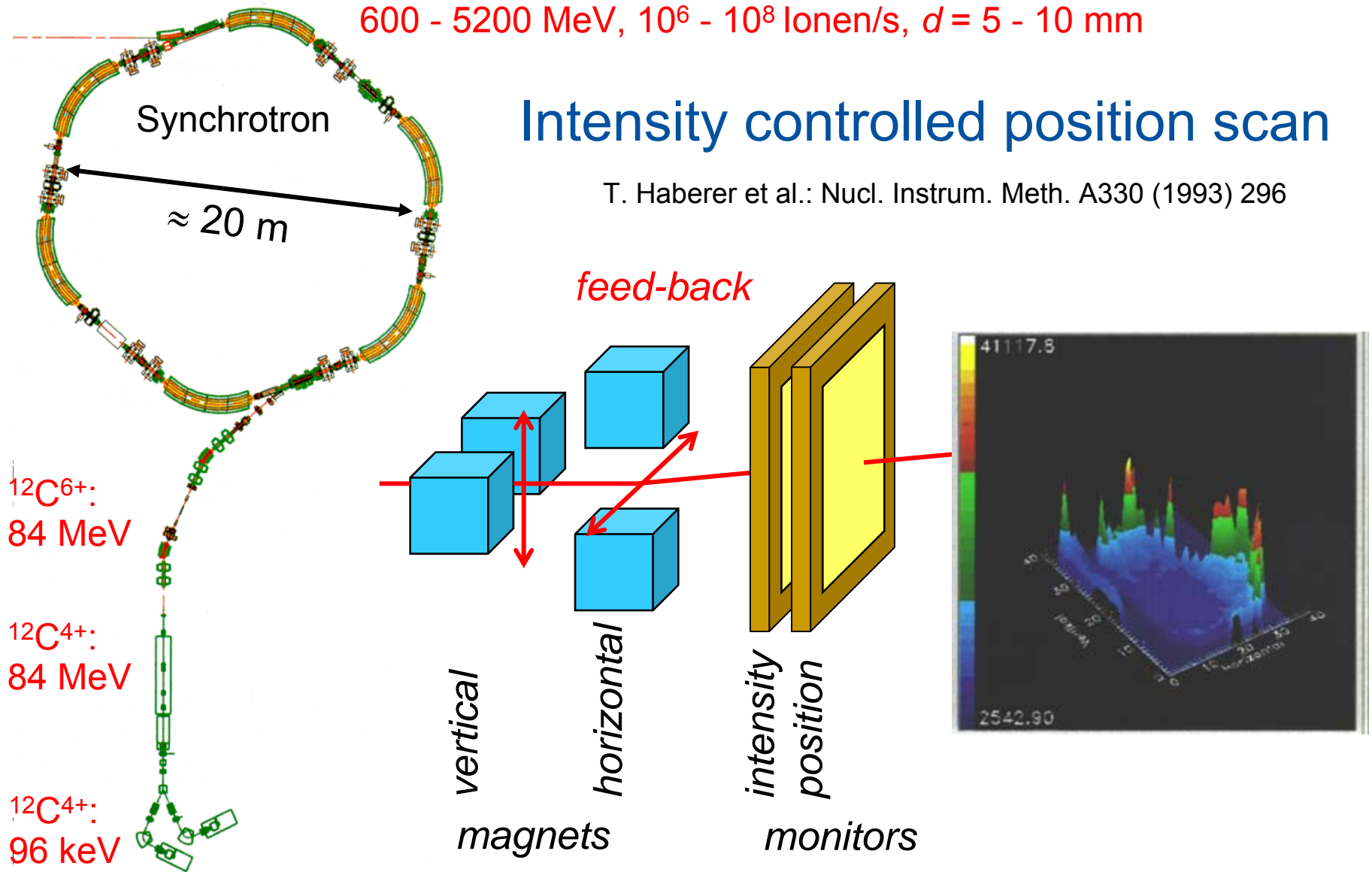


simulated

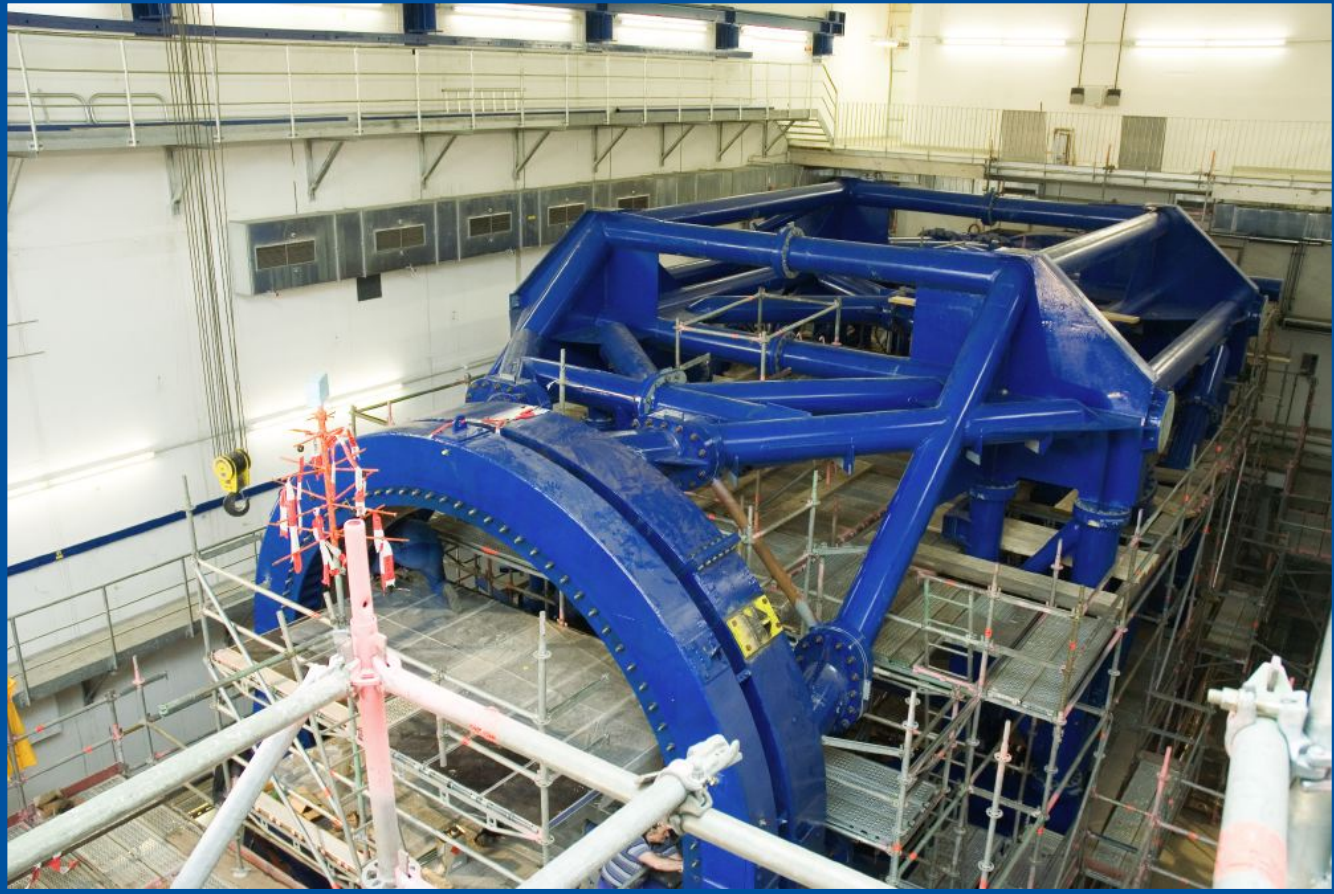
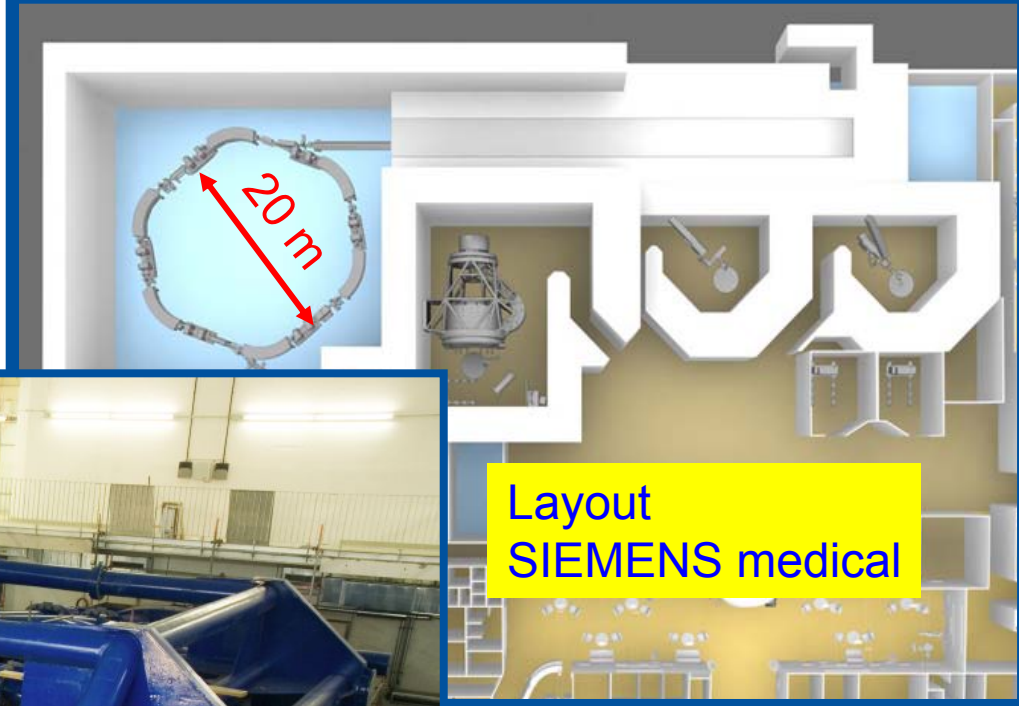
600 - 5200 MeV, $10^6 - 10^8$ Ionen/s, $d = 5 - 10$ mm

Intensity controlled position scan

T. Haberer et al.: Nucl. Instrum. Meth. A330 (1993) 296

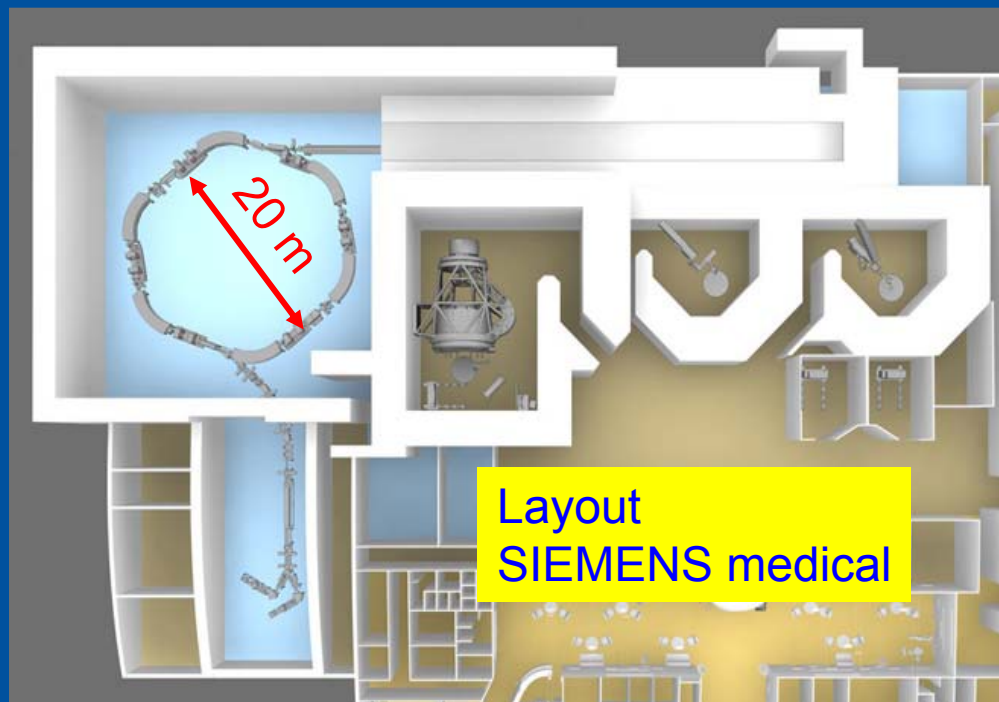


The GSI / HIT approach

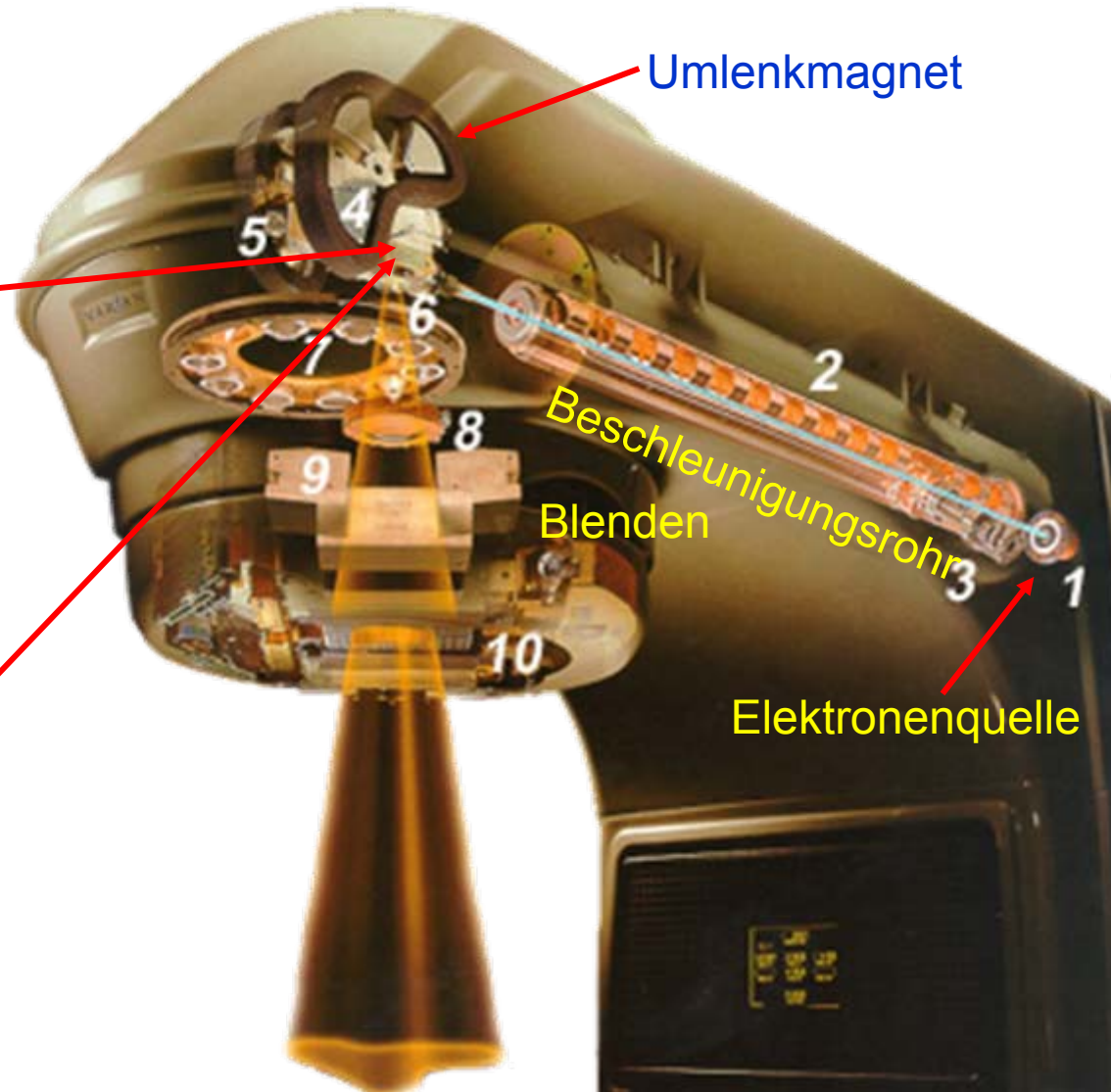
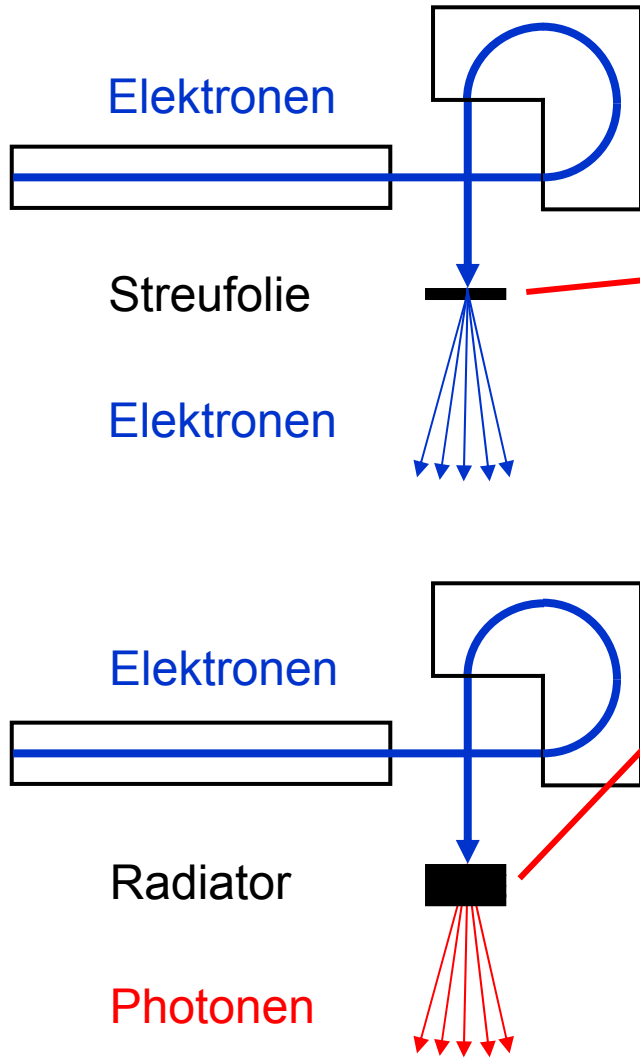


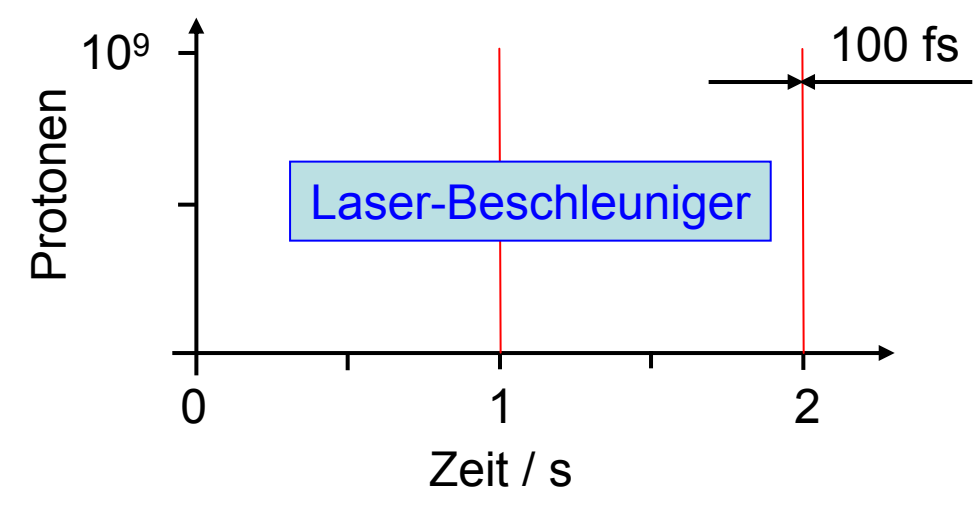
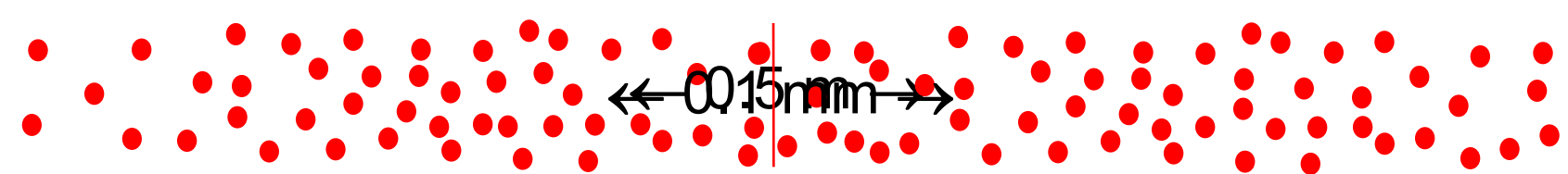
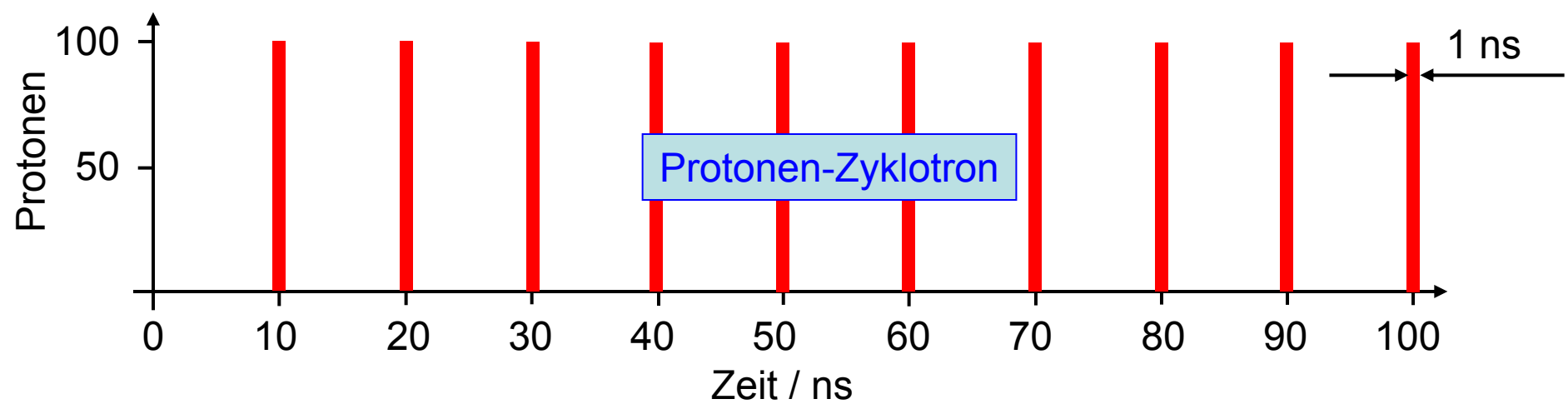
Pictures by courtesy of T. Haberer HIT Heidelberg

Conventional electron / photon therapy device



10x cheaper ...





JETI: $P \sim 1.2 \text{ J}$
 $f = 2.5 \text{ Hz}$

17.10.2007

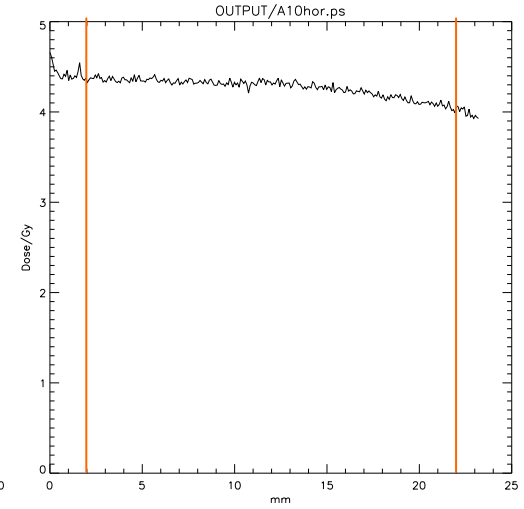
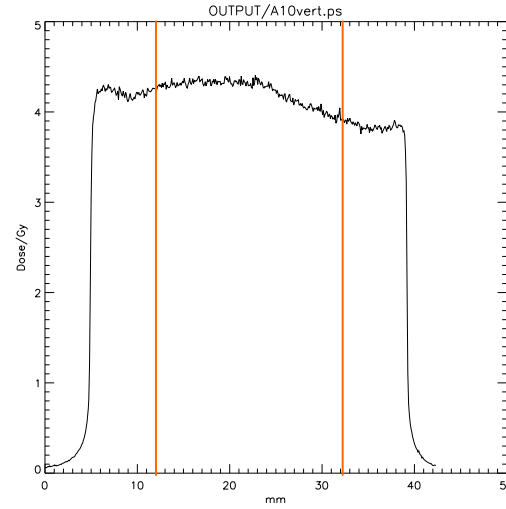
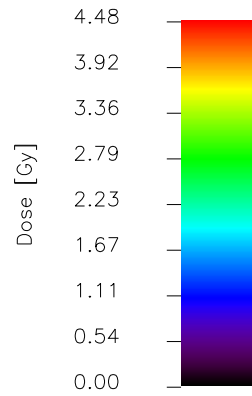
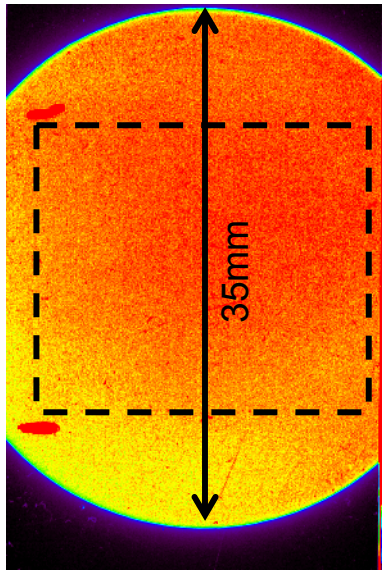
Zelllinie: FaDu_{DD} (Plattenepithelkarzinom, Kopf/Hals)
 N_{Pulse} : 150 ... 3000
Dosis: 0.263 Gy ... 4.17 Gy
 $t_{\text{Bestrahlung}}$: 60 s ... 21 min
Probenanzahl: 31 (Doppelbestimmung)

18.10.2007

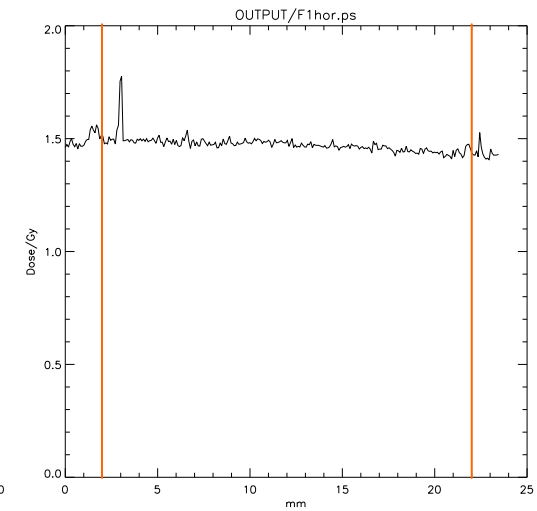
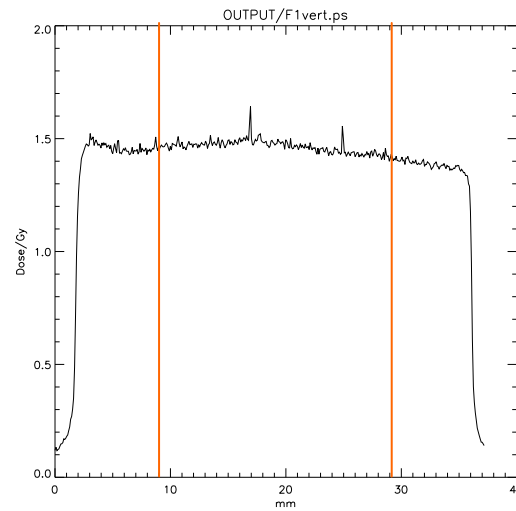
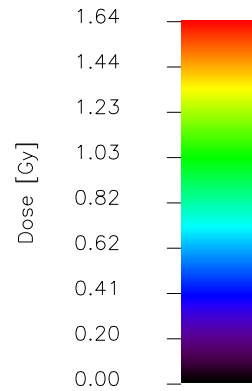
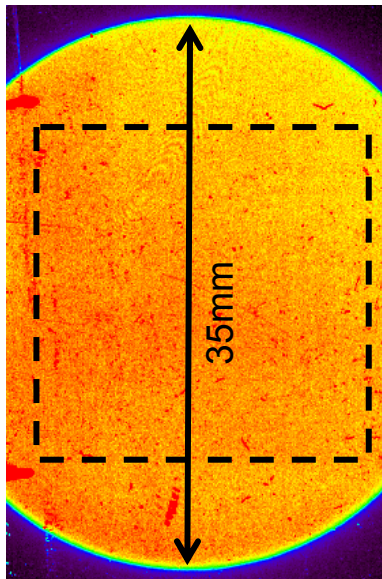
Zelllinie: 184A1 (humane Brustdrüsenepithelzellen, normal)
 N_{Pulse} : 450 ... 2400
Dosis: 0.978 Gy ... 5.210 Gy
 $t_{\text{Bestrahlung}}$: ~3 min ... ~ 16 min
Probenanzahl: 16 (Einfach- und Doppelbestimmung)

Biologischer Endpunkt: DNA Doppelstrangbrüche
(γ H2AX + 53BP1, 2h + 24h nach Bestrahlung)

Dose distribution

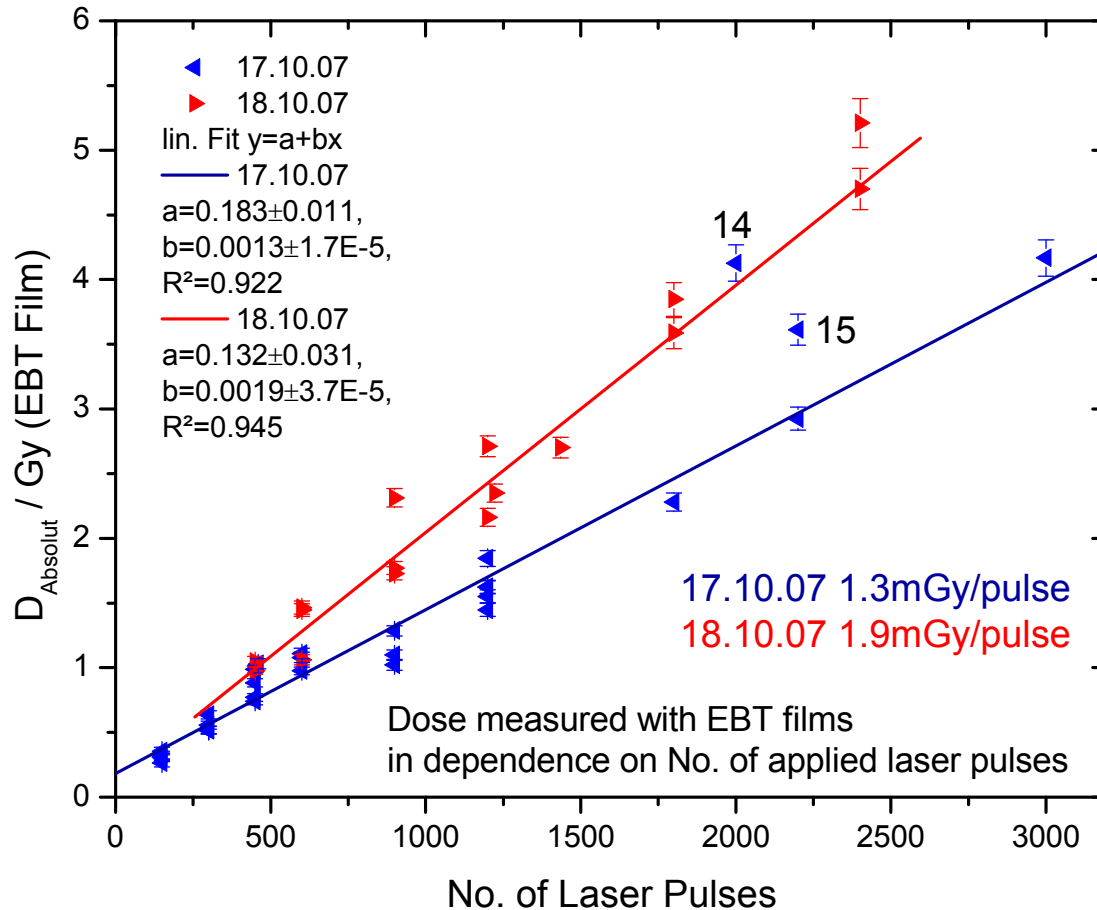


Film A10: 2000 Pulse, mittl. Dosis 4.128 Gy, 17.10.07



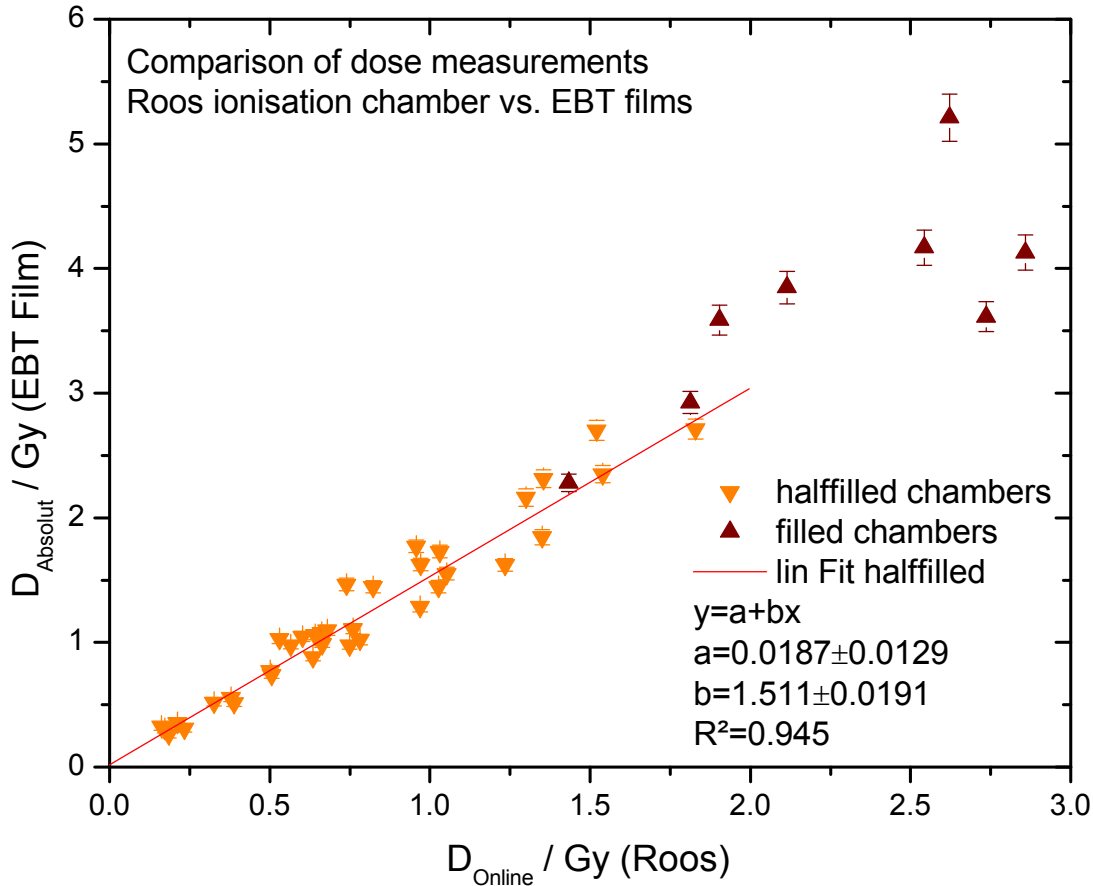
Film F1: 600 Pulse, mittl. Dosis 1.465 Gy, 18.10.07

Vergleich der Anzahl applizierter Laserpulse mit der erreichten Absolutdosis in der Zellprobe



- ⇒ Stabile Pulsdosisleistung innerhalb eines Bestrahlungstages
- ⇒ Probe 14/15: Bestrahlung nach Mittagspause, Abkühlung der Vakuumpumpen

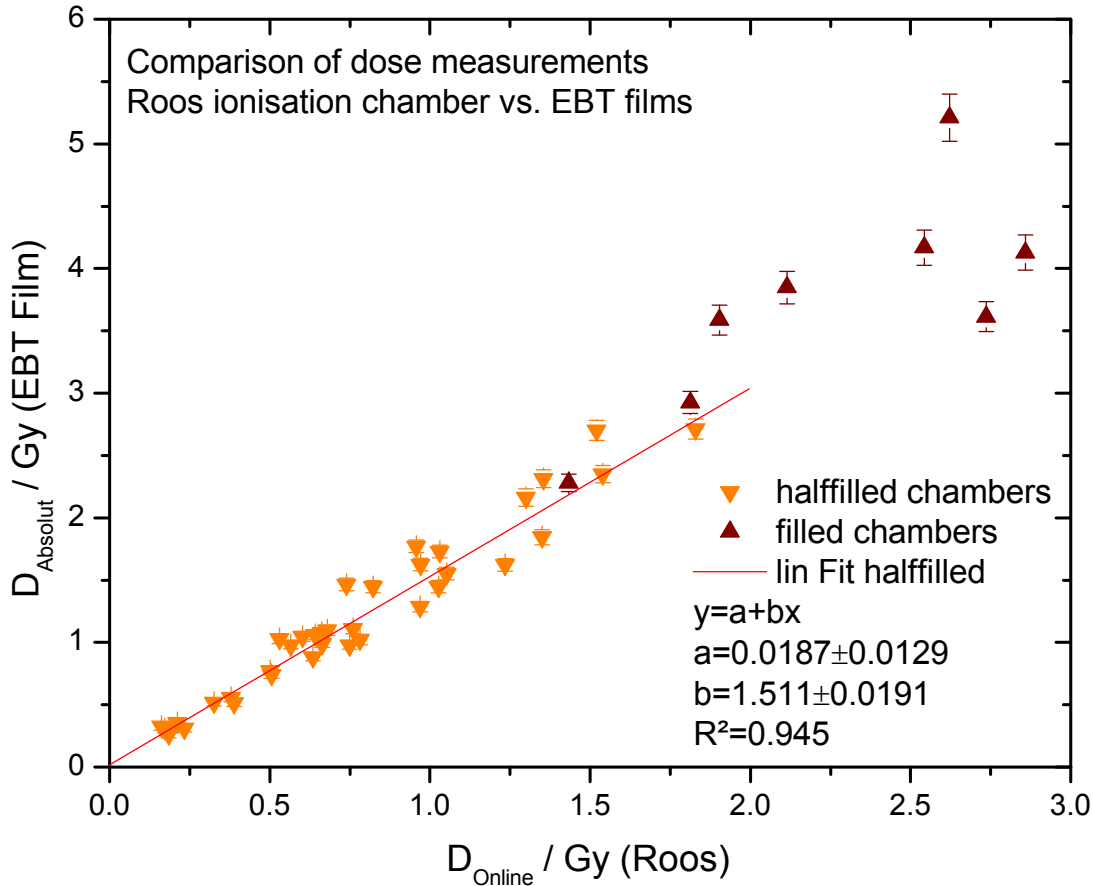
Vergleich der gemessenen Dosen mit EBT Filmen und Roos-Kammer



$D_{\text{Absolut}} / D_{\text{Online}} \approx 1.5$ (mit Zellen)

- ⇒ Gute Korrelation für „halbvolle“ Probenkammern → definiertes Probenvolumen
- ⇒ Starke Streuung für „volle“ Probenkammern → Schwankung des Volumens, Luftblasen, ...
- ⇒ Füllung der Probenkammern wegen Gefahr der Austrocknung bei $t_{\text{Bestrahlung}} > 10$ min

Vergleich der gemessenen Dosen mit EBT Filmen und Roos-Kammer

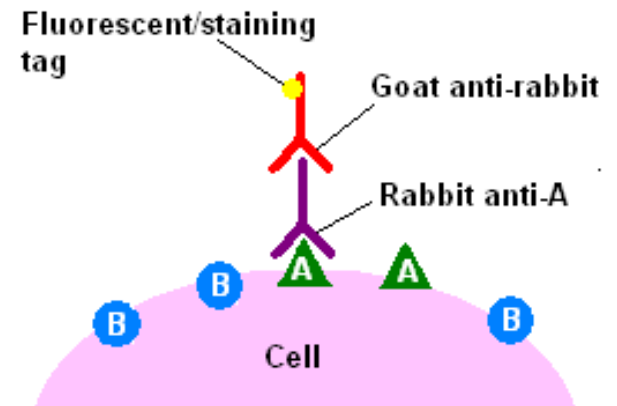
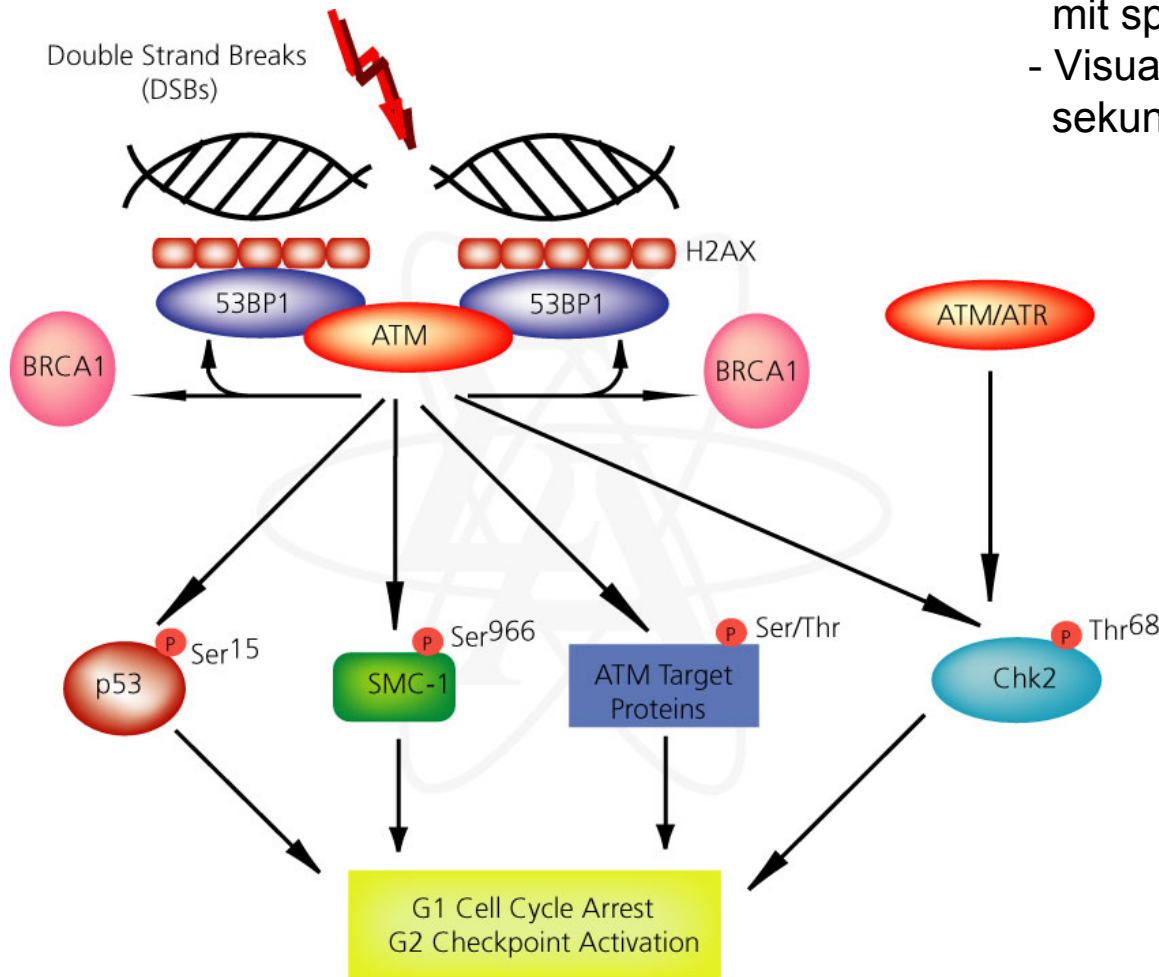


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Nachweis der Doppelstrangbrüche:

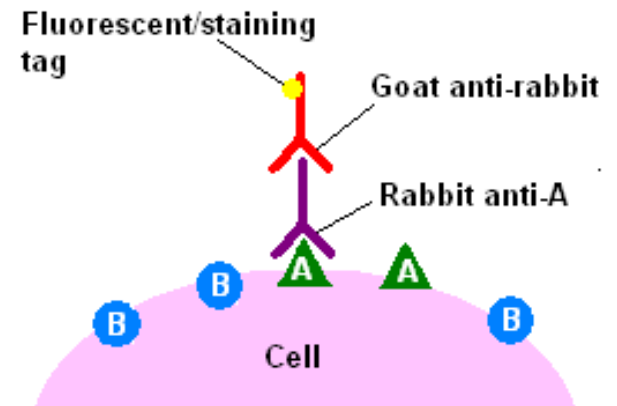
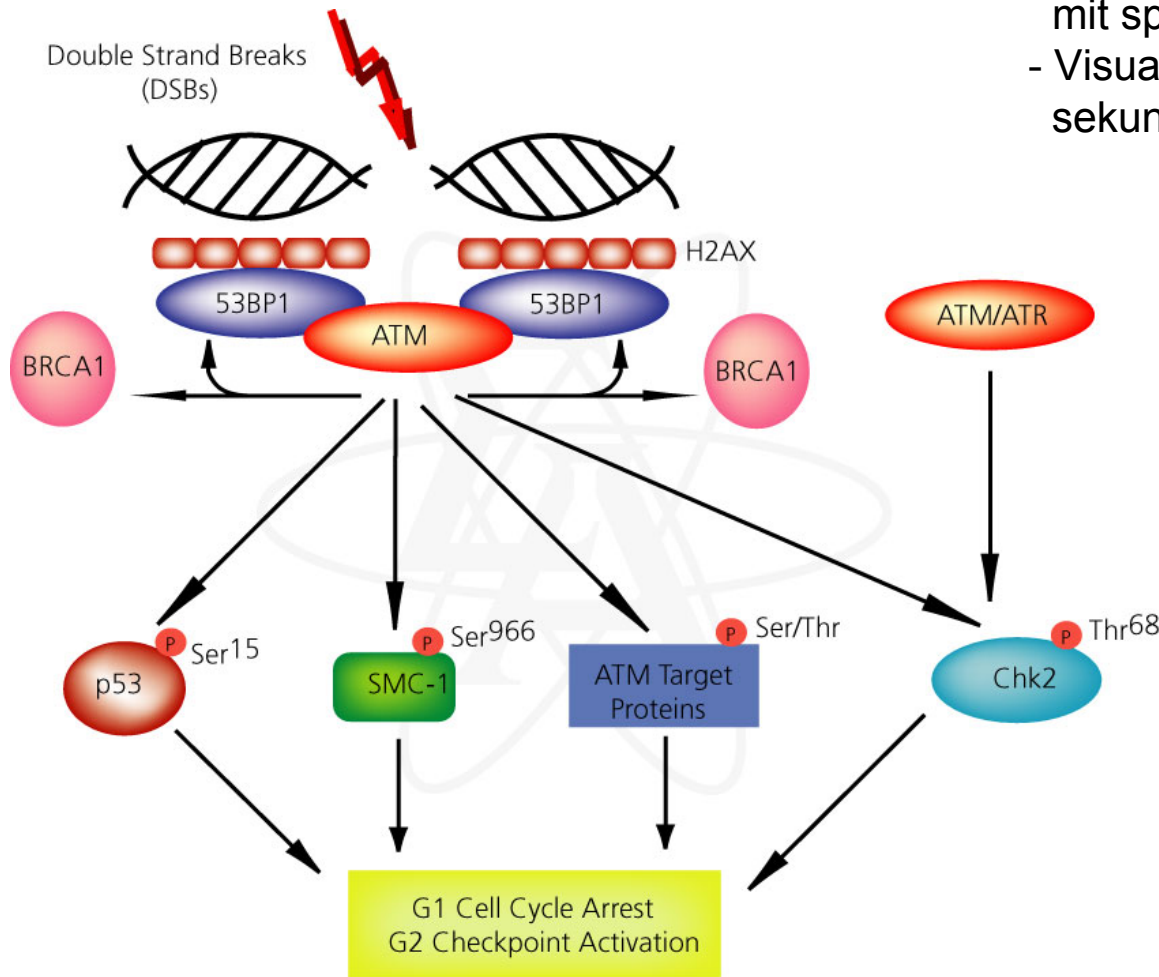
- Detektion von γ -H2AX und 53BP1 mit spezifischen Antikörpern
- Visualisierung mit fluoreszenzmarkierten sekundären AK



Nature Cell Biology 4, E277, 2002

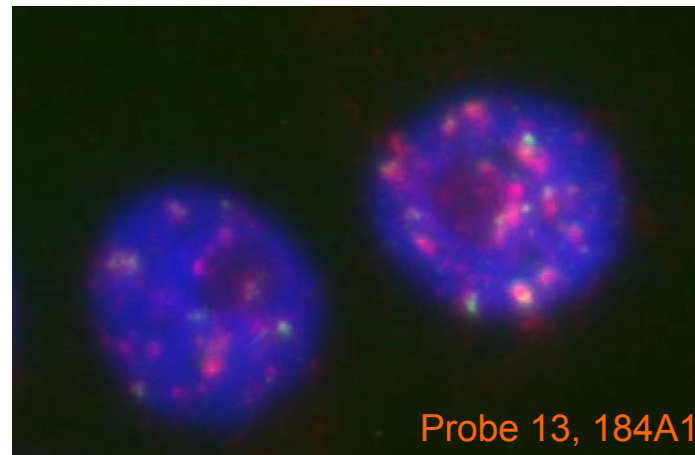
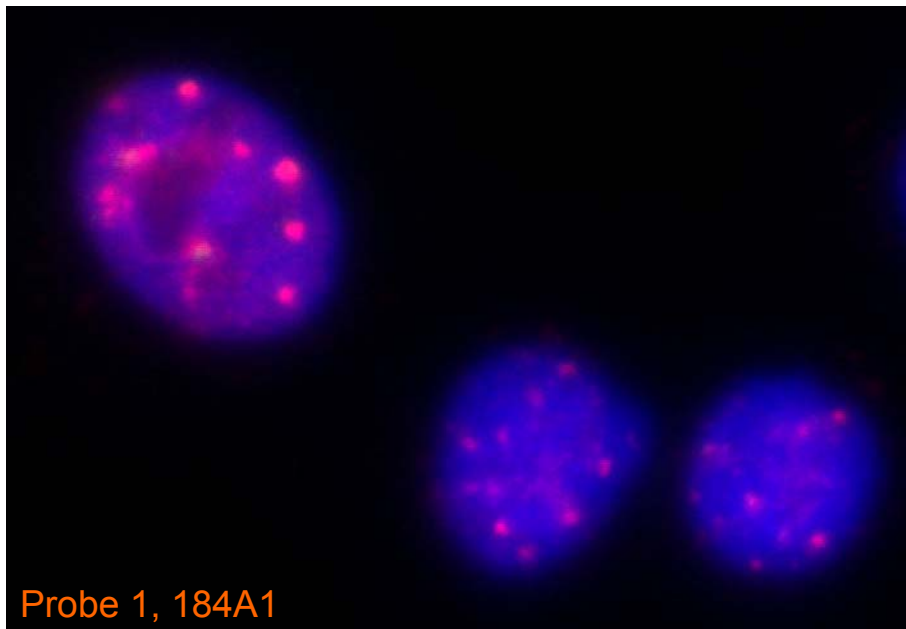
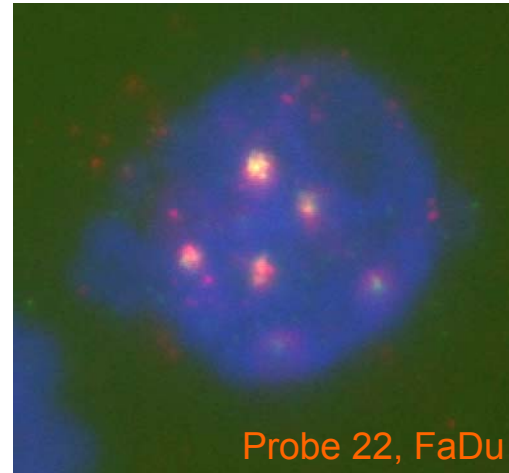
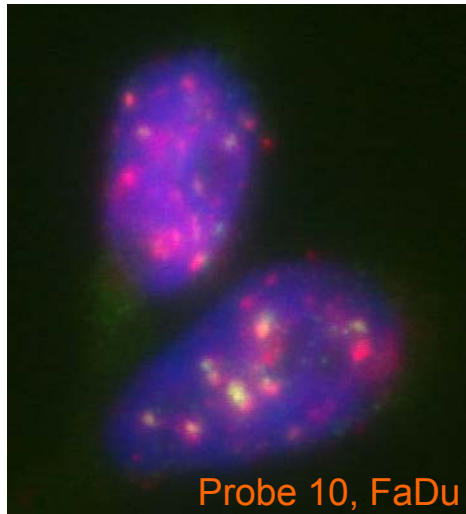
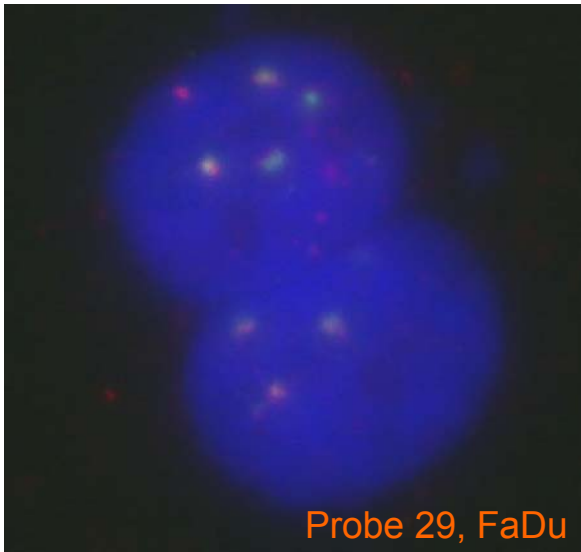
Nachweis der Doppelstrangbrüche:

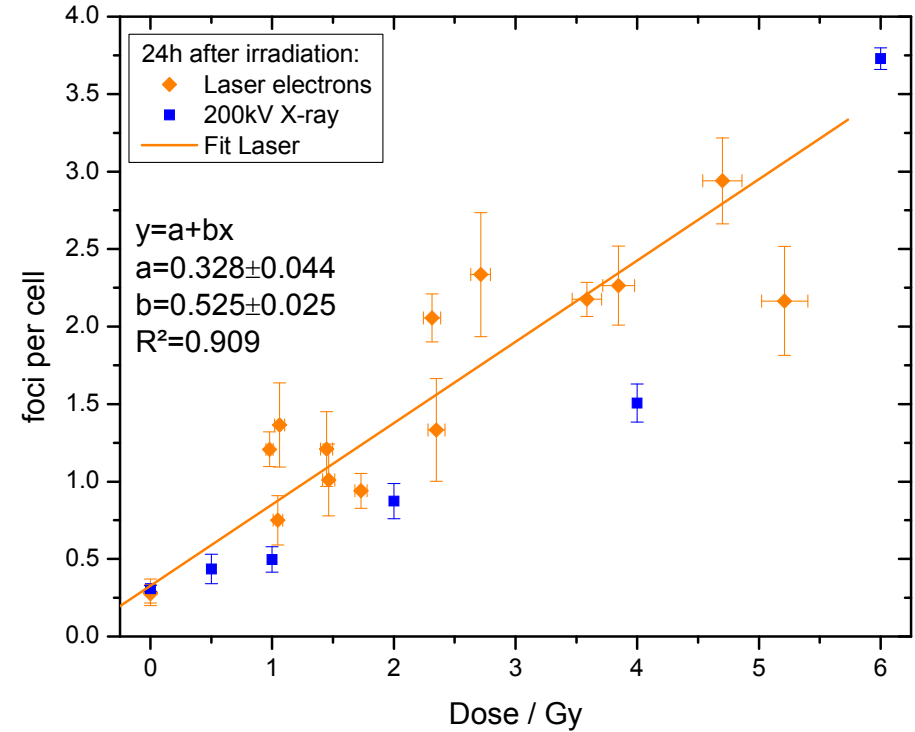
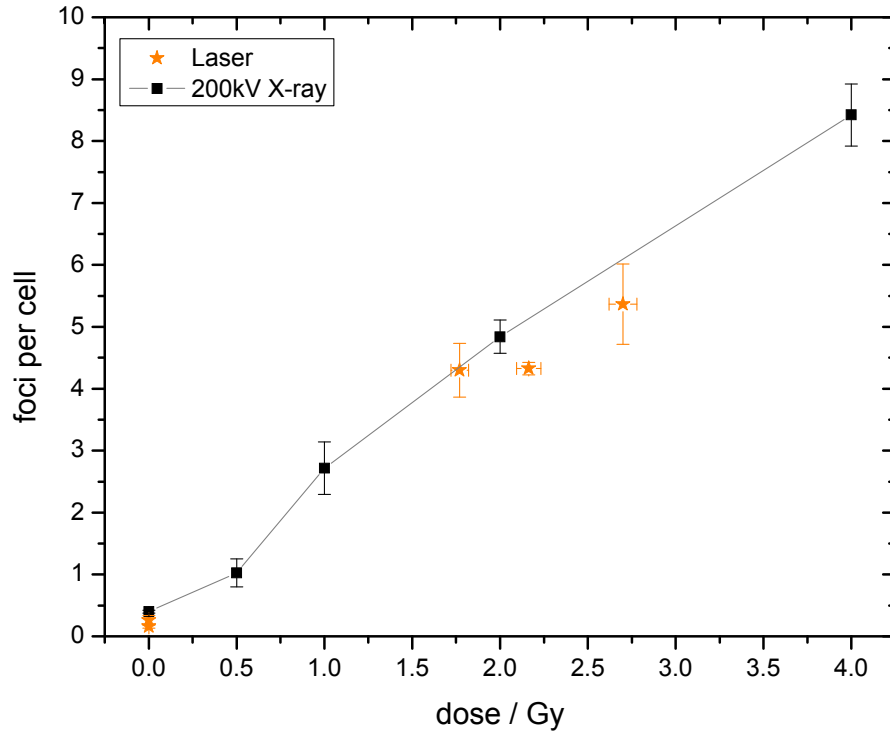
- Detektion von γ -H2AX und 53BP1 mit spezifischen Antikörpern
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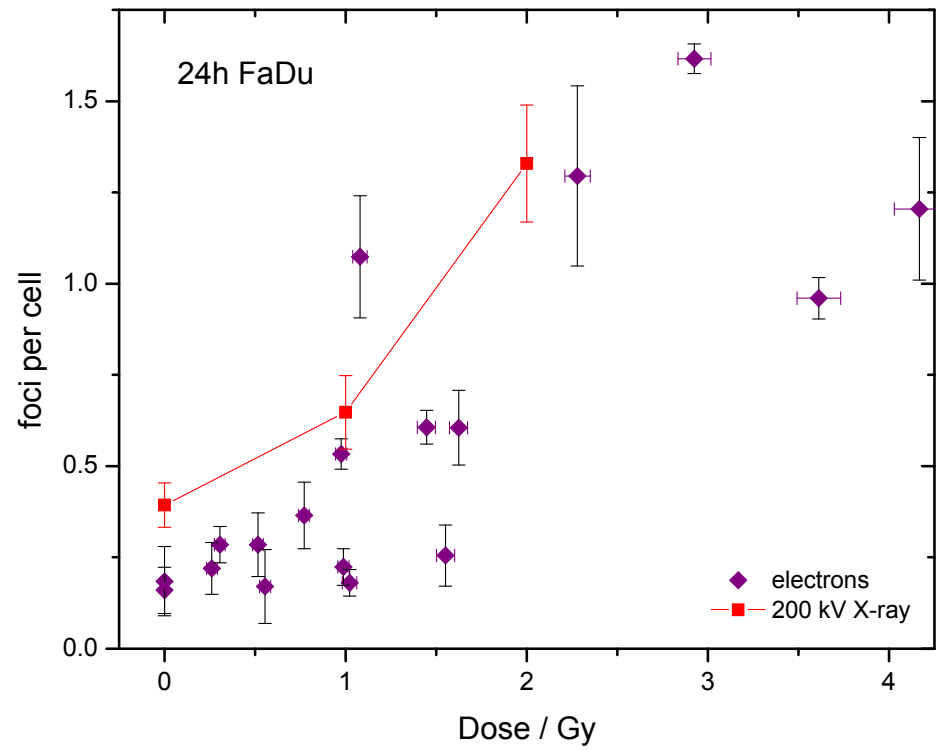
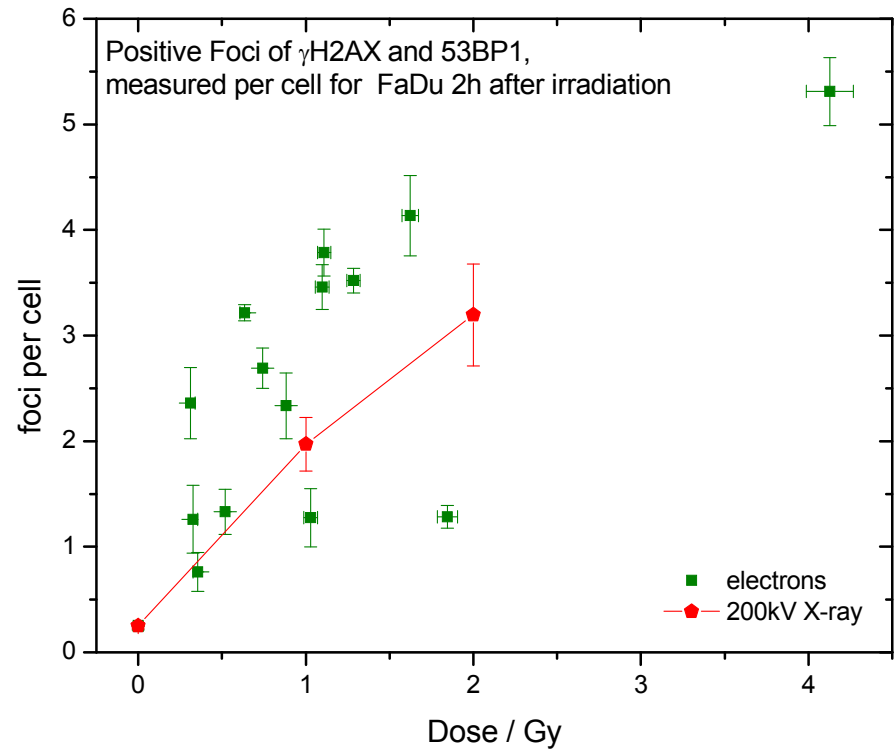
Nature Cell Biology 4, E277, 2002

Detection of DNA damage





- 2h Daten: ~ Übereinstimmung bzw. leicht verringert für MeV Elektronen
- 24h Daten: - höhere Anzahl verbleibende DSB für MeV Elektronen
 - LET? Energiespektrum?
 - Verändertes Reparaturverhalten?



- Vergleich der Werte: 2h MeV e⁻ < 200kV X ☹ 24h MeV e⁻ > 200kV X... Erklärung?
- Erwartung: Schädigung MeV Elektronen < 200 kV Photonen (LET)

Aber: Statistik 200 kV Daten unzureichend; 2 Experimente mit je 3 Dosispunkten
 → Wiederholung notwendig!

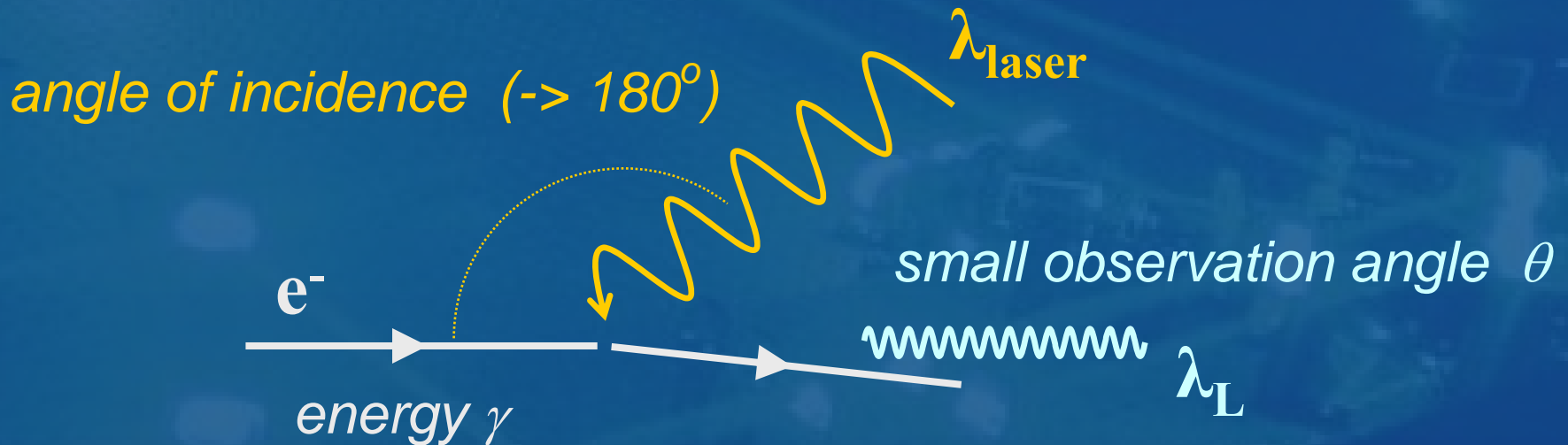
Bestrahlungen in Jena erfolgreicher als gedacht:

- Laser im „Dauereinsatz“ eingesetzt (Ausfall am 18.10.07)
- Stabile Pulsdosisleistung innerhalb eines Bestrahlungstages
- 47 bestrahlte Zellproben

Unklarheiten/nächste Schritte:

- Energiespektrum der Elektronen nicht ausreichend genau bekannt
⇒ Einfluss?, Messung?, Dosimetrie?
- Interpretation der gemessenen Daten schwierig, fehlende/zu wenige Referenzdaten
⇒ „Zeitnahe“ Bestimmung von Referenzwerten für FaDu notwendig
(an 200kV Röhre, parallel zu JeTi-Bestrahlungen)
- Referenzbestrahlungen:
 - Therapiebeschleuniger mit versch. Absorptionsmedien
 - Monoenergetische Elektronenstrahlung (ELBE)
 - 200 kV Röntgenröhre → bisher nur für 184A1
- Etablierung einer zweiten Tumorzelllinie + Referenzbestrahlungen

Thomson scattering – basic idea



photon energy, distribution, efficiency, coherence ...



Forschungszentrum
Dresden Rossendorf

Thomson scattering – head on



$$\lambda_L = \frac{\lambda_{laser}}{4\gamma^2} \left(1 + \frac{a^2}{2} + \gamma^2 \theta^2 \right)$$

Resonance conditions (undulator and `optical undulator`):

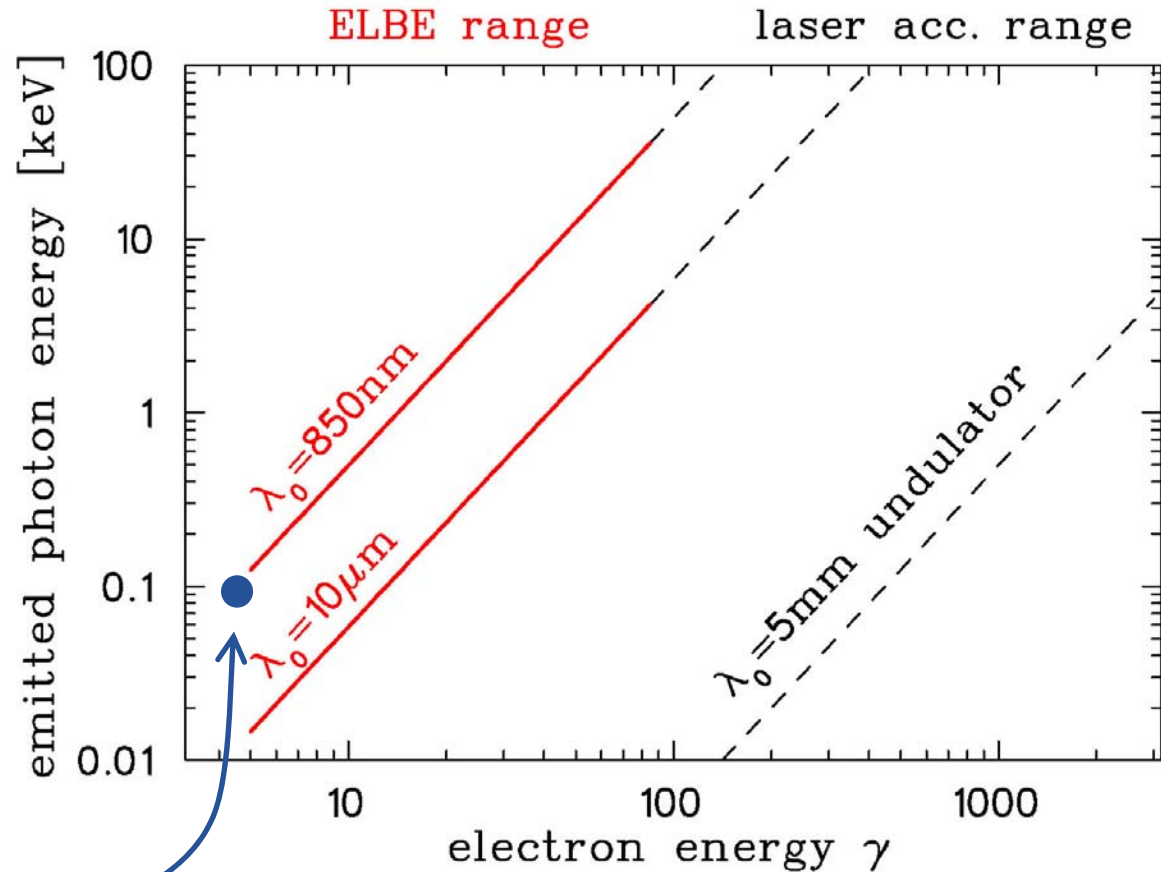
- *Thomson scattering:*

$$\lambda_L = \frac{\lambda_{laser}}{4\gamma^2} \left(1 + \frac{a^2}{2} + \gamma^2 \theta^2 \right)$$

- *Undulator radiation:*

$$\lambda_L = \frac{\lambda_u}{2\gamma^2} \left(1 + \frac{K^2}{2} + \gamma^2 \theta^2 \right)$$

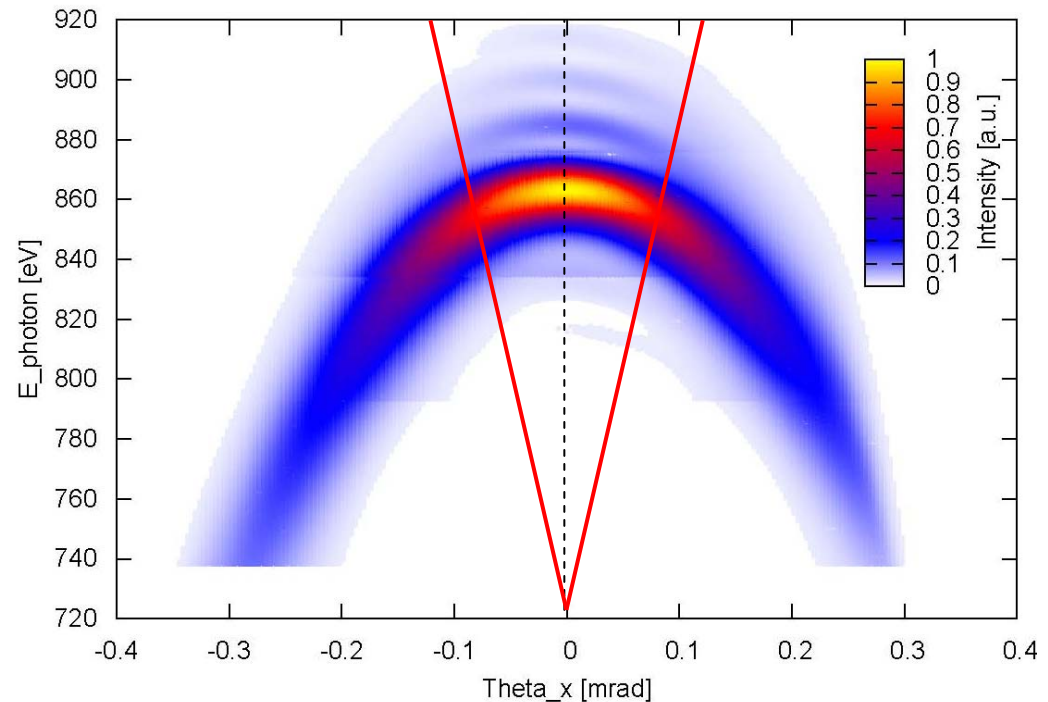
for 13.5 nm (92eV)
 $\gamma=4.4$ (or 1.7 MeV)



(treated as a counterpropagating undulator)

$$\lambda_L = \frac{\lambda_{laser}}{4\gamma^2} \left(1 + \frac{a^2}{2} + \gamma^2 \theta_{obs}^2 \right)$$

$$\theta_{emission,rms} \sim \frac{1}{\gamma \sqrt{N_{laser}}}$$



Wavelength @ rms emission angle:

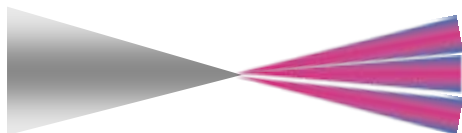
$$\lambda_L = \frac{\lambda_{laser}}{4\gamma^2} \left(1 + \frac{1}{N_l} \right)$$

Relative width @ zero angle:

$$\frac{\Delta\lambda_L}{\lambda_L} \sim \frac{1}{N_l}$$

For a reasonable number of oscillation periods of
 $N_{laser} \sim 100$ (i.e. 300fs for Ti:Sapphire laser)

► e-beam divergence defines radiation properties



Thomson scattering rate



yield limited by small cross-section ($\sim r_e^2$)

$$dN_{emission} / d\Omega \sim \alpha \cdot \gamma^2 \cdot a_0^2 \cdot N_e \cdot N_{laser}^2$$

$$N_{emission} \sim \alpha \cdot a_0^2 \cdot N_e \cdot N_{laser} \cdot f_{rep}$$

1 nC, few ps
 $\epsilon_n = 1$ mm mrad

~ 100 - 300 fs (bandwidth,
divergence)

1 (J-class) \rightarrow 10^8 photons / pulse @ 10Hz = 10^9 /s
few mrad cone and % bandwidth (~10nW)