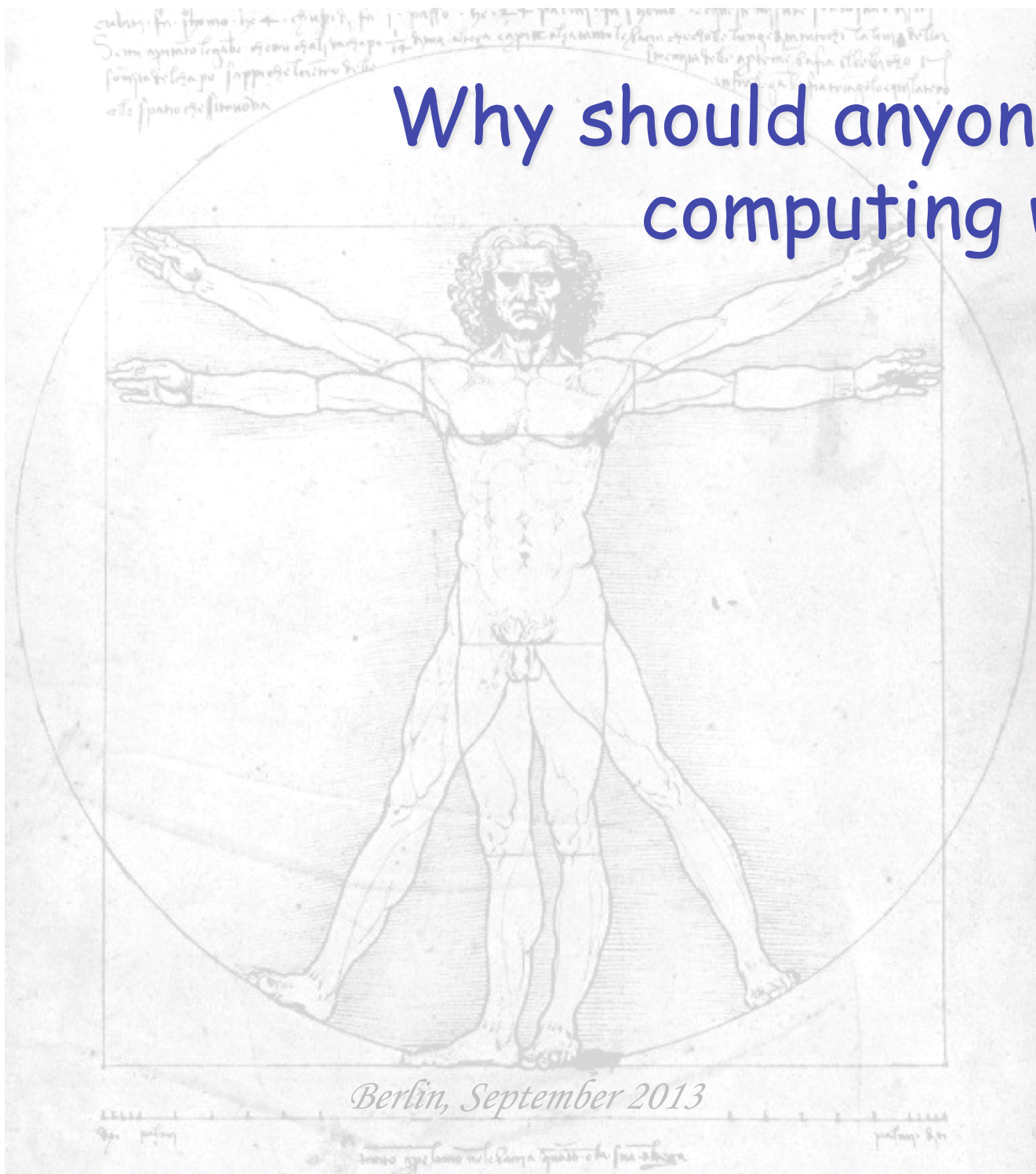


Why should anyone care about computing with anyons?

Jiannis K. Pachos

Toric Code and
Quantum Doubles



EPSRC

Engineering and Physical Sciences
Research Council



UNIVERSITY OF LEEDS

Anyons and Quantum Computation

- Error correction needs a huge overhead.
- Instead of performing active error correction let physics do the job.
- Perform QC in a physical medium that is **gapped** and **highly correlated**:
 - **Energy penalty for errors (gapped).**
 - **Make logical errors non-local (very unlikely).**
- Similar to **quantum error correction**, but without active control.

Toric Code: ECC

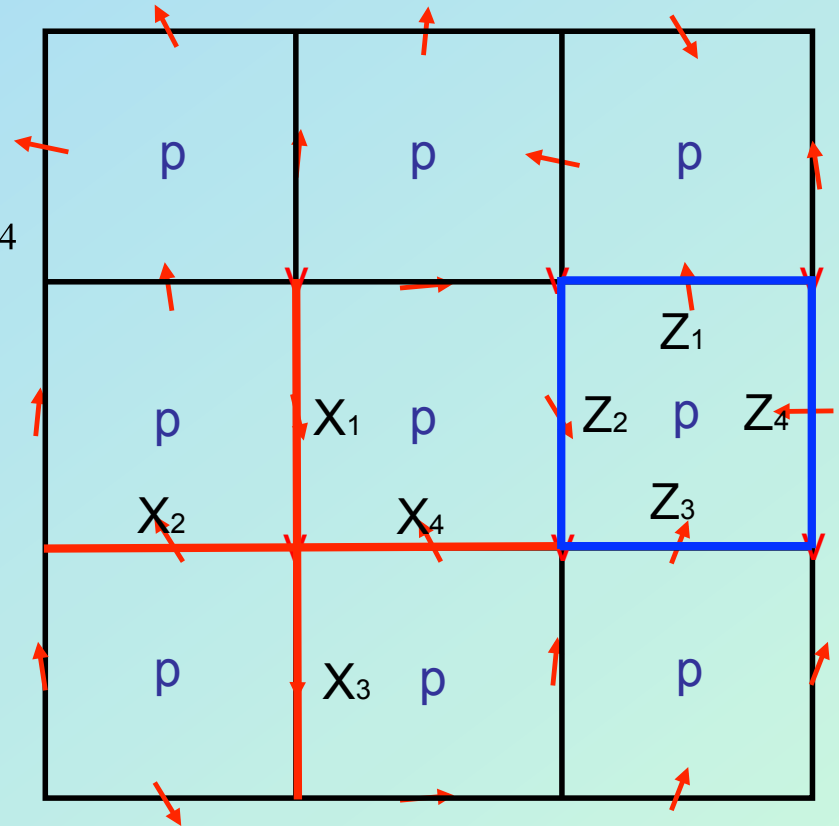
Consider the lattice Hamiltonian

$$H = - \sum_p Z_{p1} Z_{p2} Z_{p3} Z_{p4} - \sum_v X_{v1} X_{v2} X_{v3} X_{v4}$$

Spins on the edges.

Two different types of interactions: ZZZZ or XXXX acting on plaquettes and vertices respectively.

The four spin interactions involve spins of the same vertex/plaquette.



Toric Code: ECC

Consider the lattice Hamiltonian

$$H = - \sum_p Z_{p1} Z_{p2} Z_{p3} Z_{p4} - \sum_v X_{v1} X_{v2} X_{v3} X_{v4}$$

Good quantum numbers:

$$[H, Z_{p1} Z_{p2} Z_{p3} Z_{p4}] = 0$$

$$[H, X_{v1} X_{v2} X_{v3} X_{v4}] = 0$$

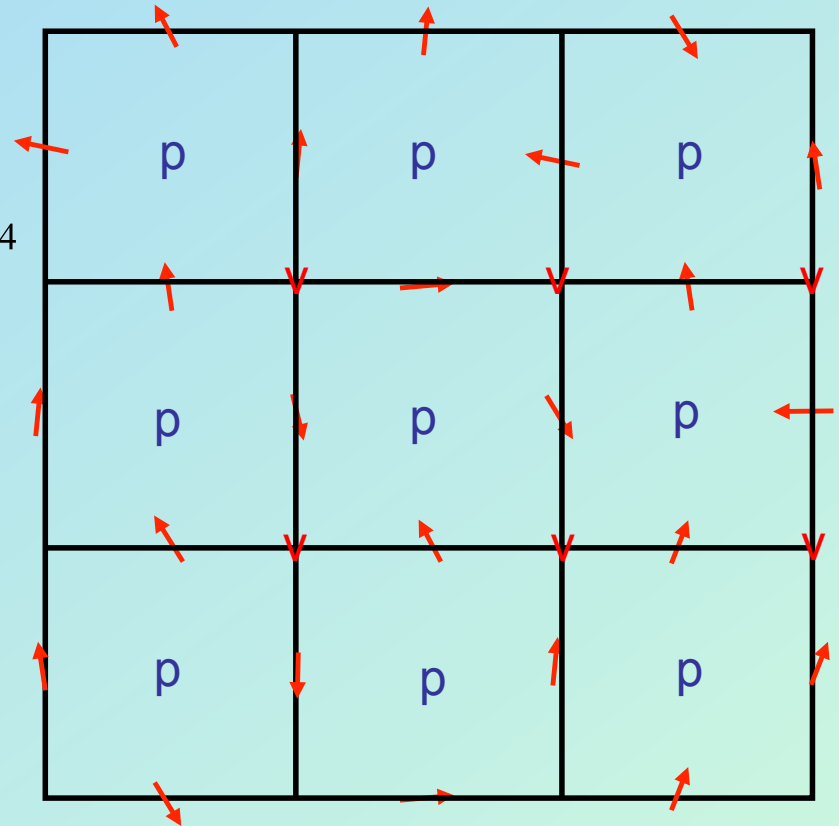
$$(X_{v1} X_{v2} X_{v3} X_{v4})^2 = 1$$

$$(Z_{p1} Z_{p2} Z_{p3} Z_{p4})^2 = 1$$

⇒ eigenvalues of XXXX and ZZZZ: ± 1

Also Hamiltonian **exactly solvable**:

$$[X_{v1} X_{v2} X_{v3} X_{v4}, Z_{p1} Z_{p2} Z_{p3} Z_{p4}] = 0$$



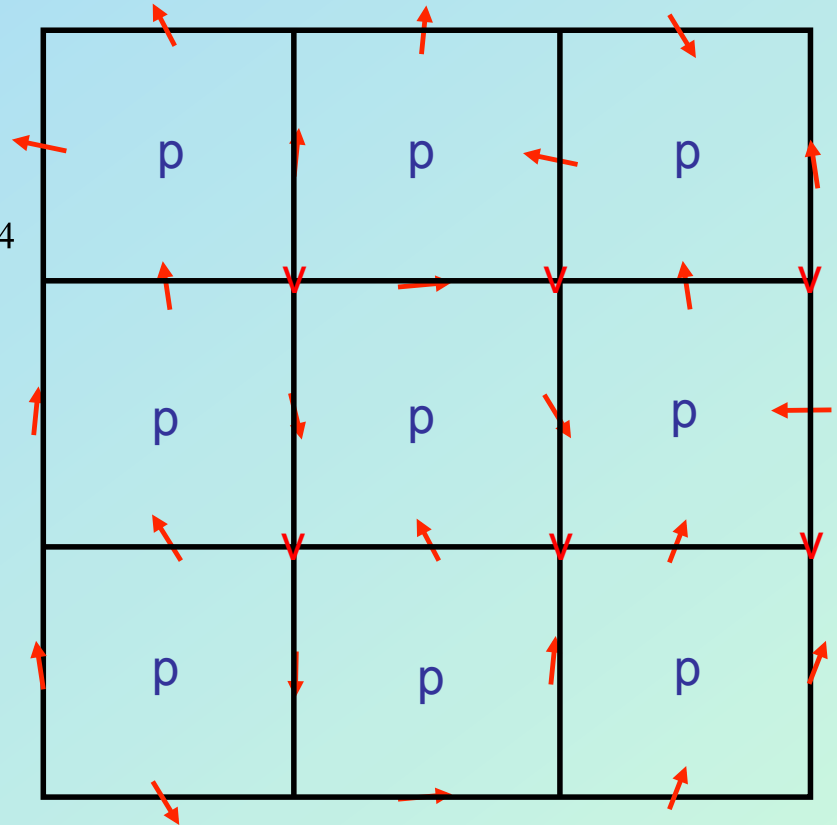
Toric Code: ECC

Consider the lattice Hamiltonian

$$H = - \sum_p Z_{p1} Z_{p2} Z_{p3} Z_{p4} - \sum_v X_{v1} X_{v2} X_{v3} X_{v4}$$

Indeed, the ground state is:

$$|\xi\rangle = \prod_v \frac{1}{\sqrt{2}} (\mathbb{I} + X_{v1} X_{v2} X_{v3} X_{v4}) |00\dots 0\rangle$$



The $|00\dots 0\rangle$ state is a ground state of the $ZZZZ$ term.

The $(\mathbb{I} + XXXX)$ term projects that state to the ground state of the $XXXX$ term.

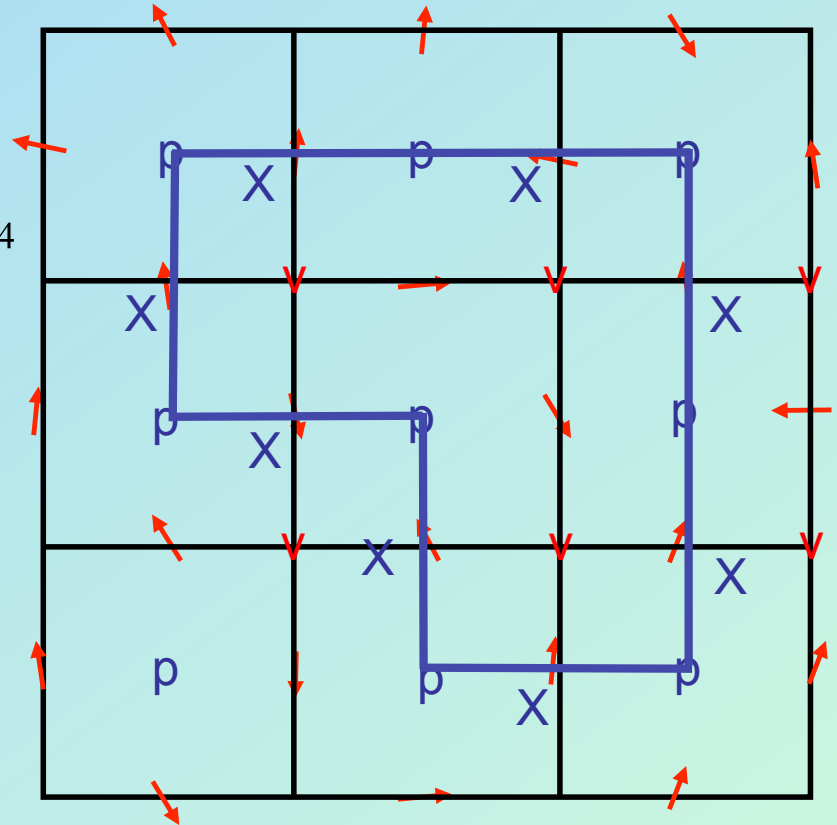
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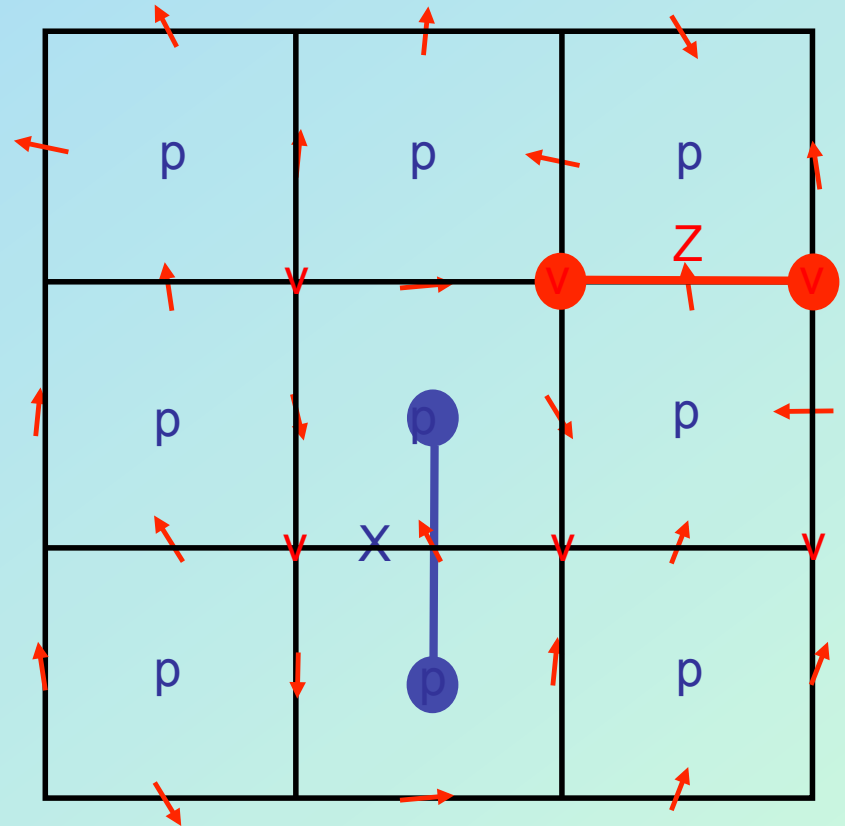
The ground state is a **superposition** of all X loops.

It is **stabilized** by the application of all X loop operators.

Equivalently for Z loops.

Toric Code: ECC

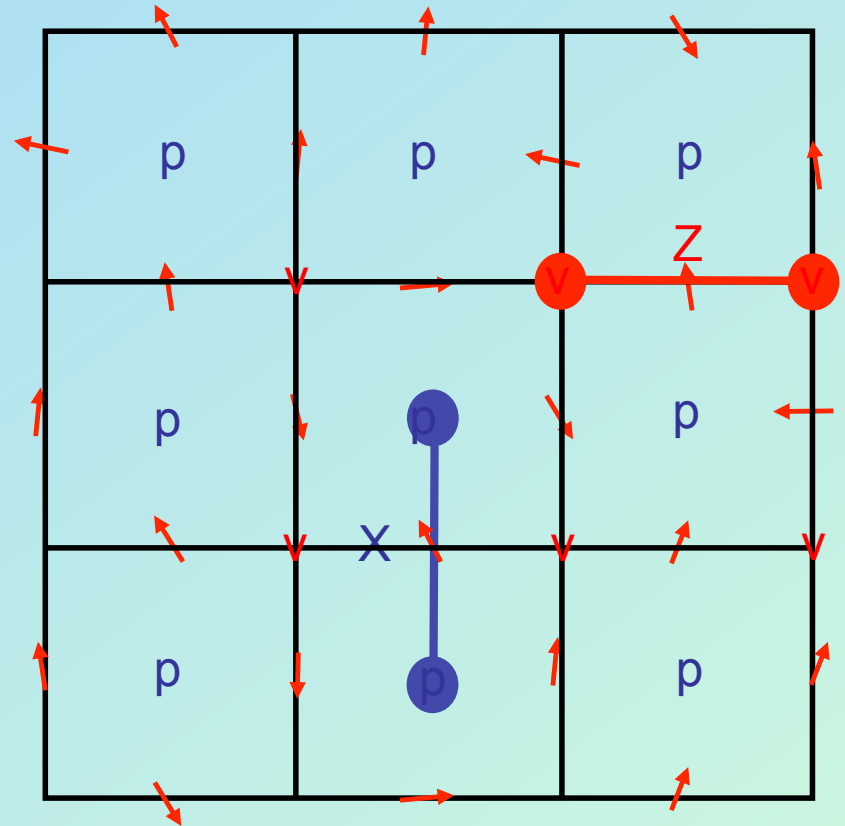
- **Excitations** are produced by Z or X rotations of one spin.
- These rotations **anticommute** with the X - or Z -part of the Hamiltonian, respectively.
- Z excitations on v vertices.
- X excitations on p plaquettes.



X and Z excitations behave as anyons with respect to each other.

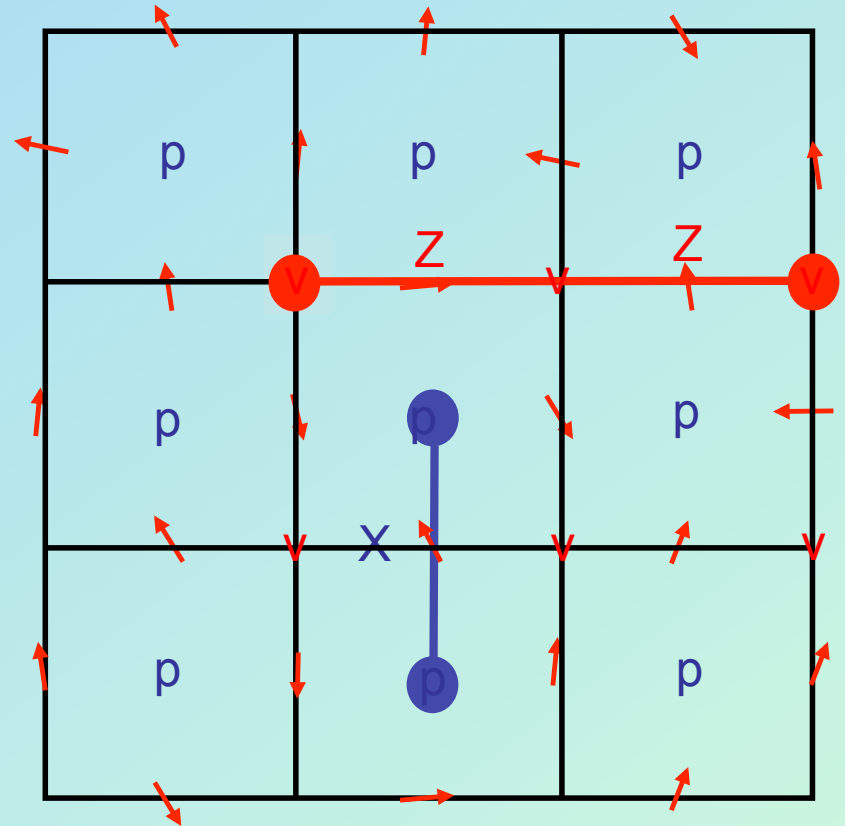
Toric Code: ECC

One can demonstrate the anyonic statistics between X and Z .
First create excitations with Z and X rotations.



Toric Code: ECC

One can demonstrate the anyonic statistics between X and Z .
First create excitations with Z and X rotations.
Then rotate Z excitation around the X one.

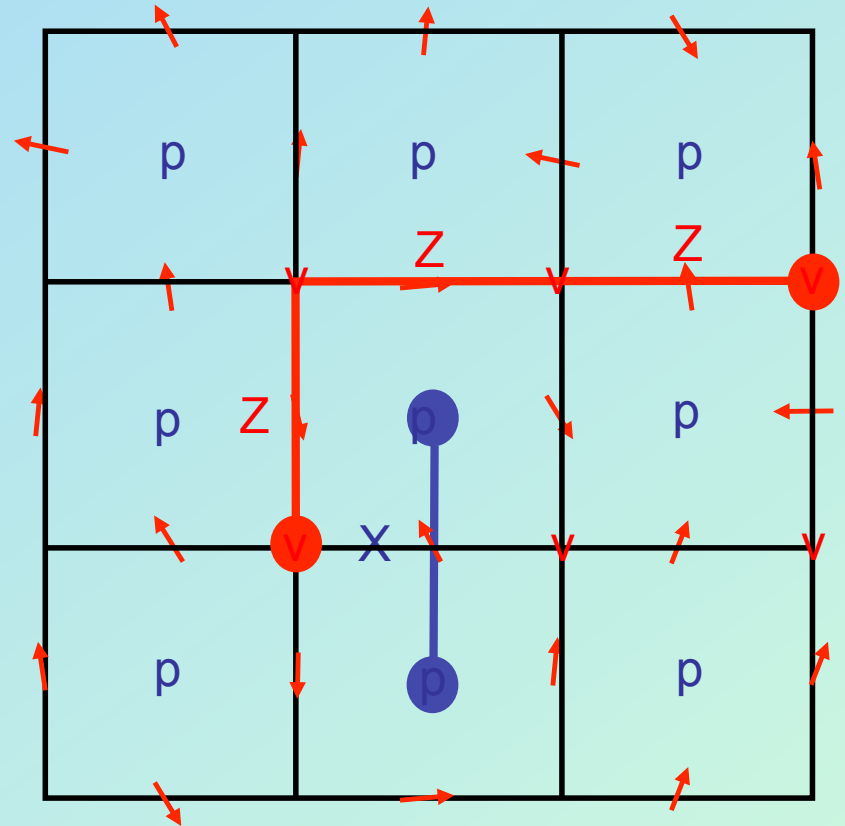


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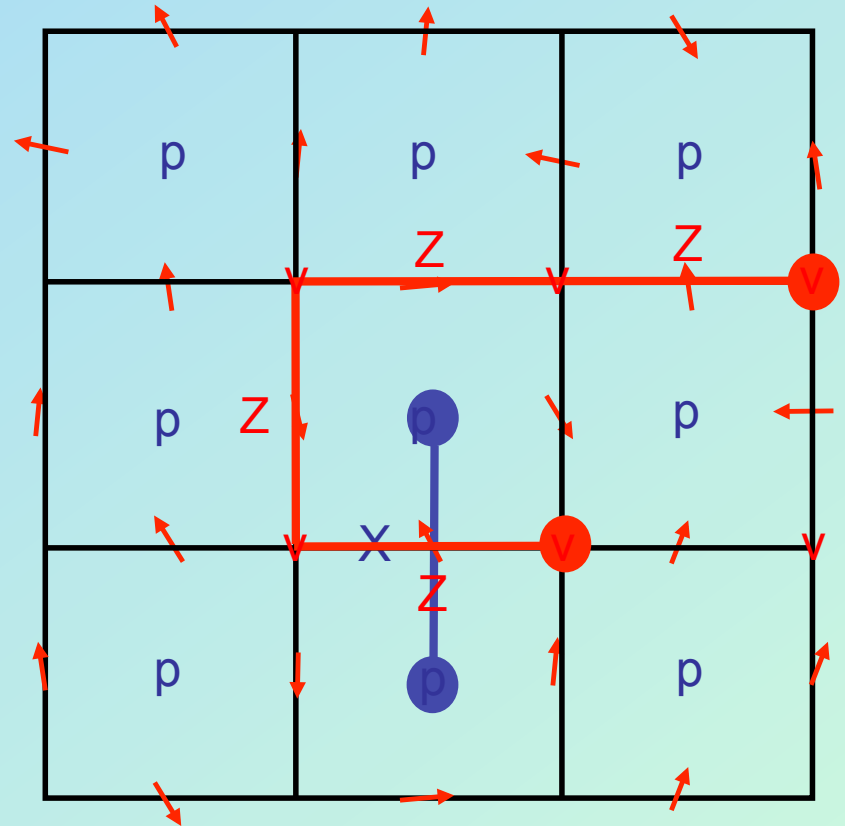


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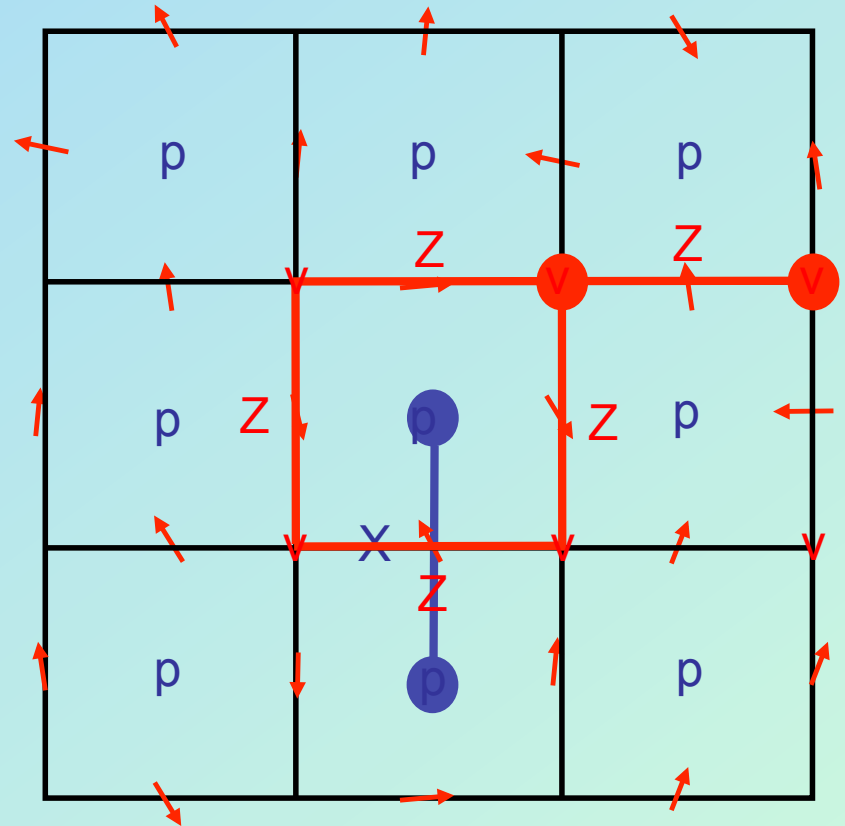
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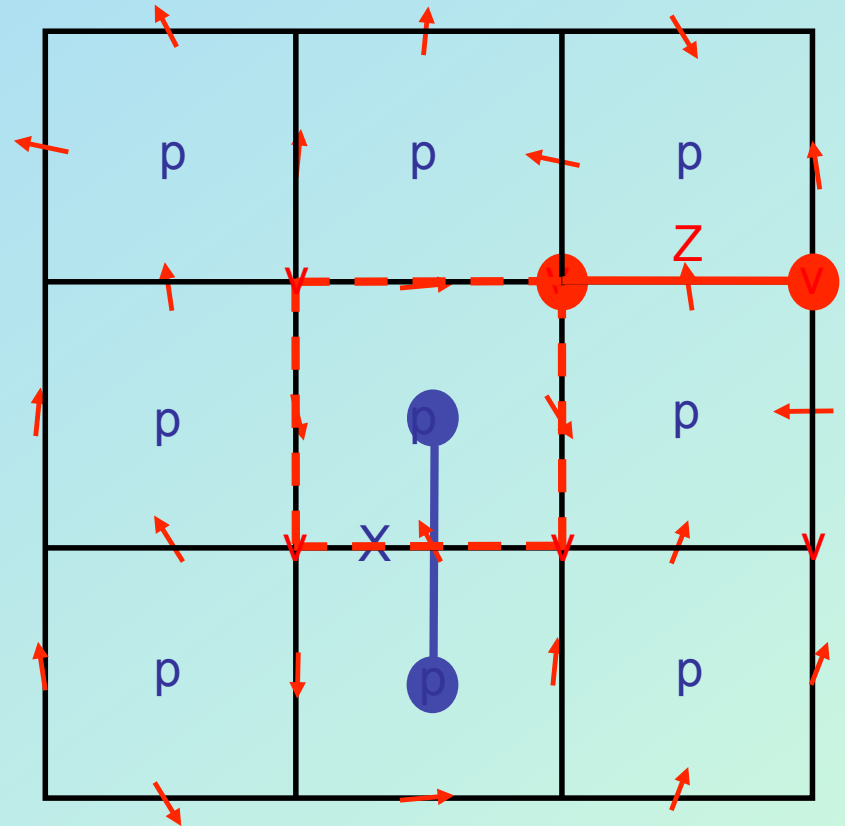
Toric Code: ECC

One can demonstrate the anyonic statistics between X and Z.

First create excitations with Z and X rotations.

Then rotate Z excitation around the X one.

This results in plaquette operator detecting the X excitation. Gives -1



$$\begin{aligned}
 |Final\rangle &= X_4 X_3 X_2 X_1 |Z\rangle = (X_4 X_3 X_2 X_1) Z_3 |\xi\rangle \\
 &= -Z_3 (X_4 X_3 X_2 X_1) |\xi\rangle = -|Initial\rangle
 \end{aligned}$$

Toric Code: ECC

$$\begin{aligned} |Final\rangle &= X_4 X_3 X_2 X_1 |Z\rangle = (X_4 X_3 X_2 X_1) Z_1 |\xi\rangle \\ &= - Z_1 (X_4 X_3 X_2 X_1) |\xi\rangle = - |Initial\rangle \end{aligned}$$

Anyonic statistics

After a complete rotation of an X anyon around a Z anyon (two successive exchanges) the resulting state gets a phase π (a minus sign): hence **ANYONS** with statistical angle $\pi / 2$

A property we used is that $X_4 X_3 X_2 X_1 |\xi\rangle = |\xi\rangle$

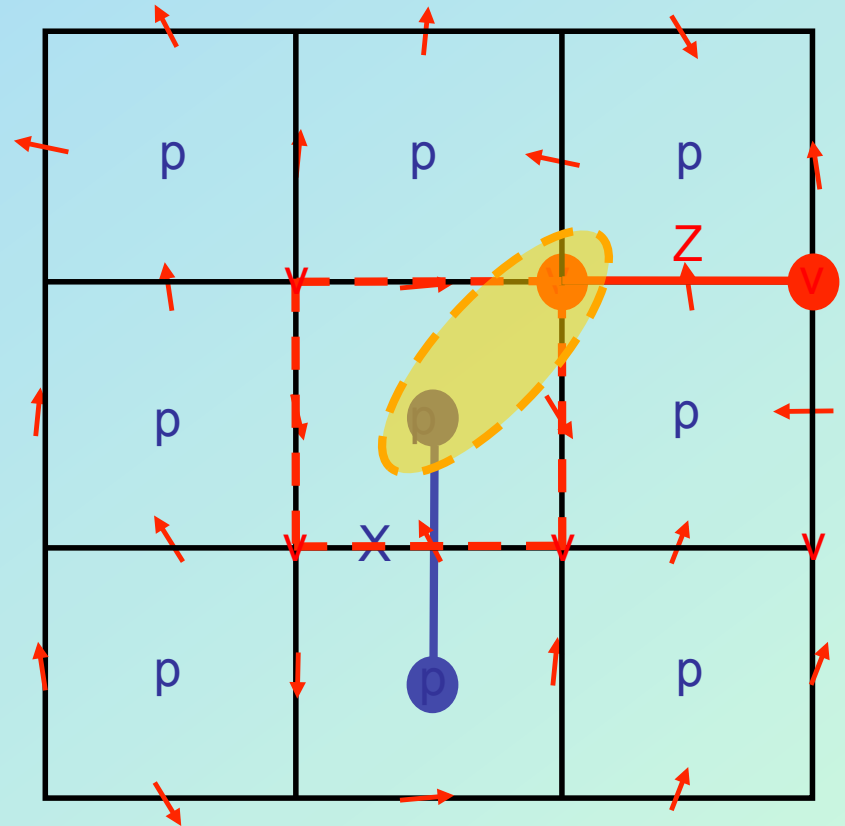
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Toric Code: Anyons

Hence Toric Code has particles:

$1, e (X), m (Z), \varepsilon$ (fermion)

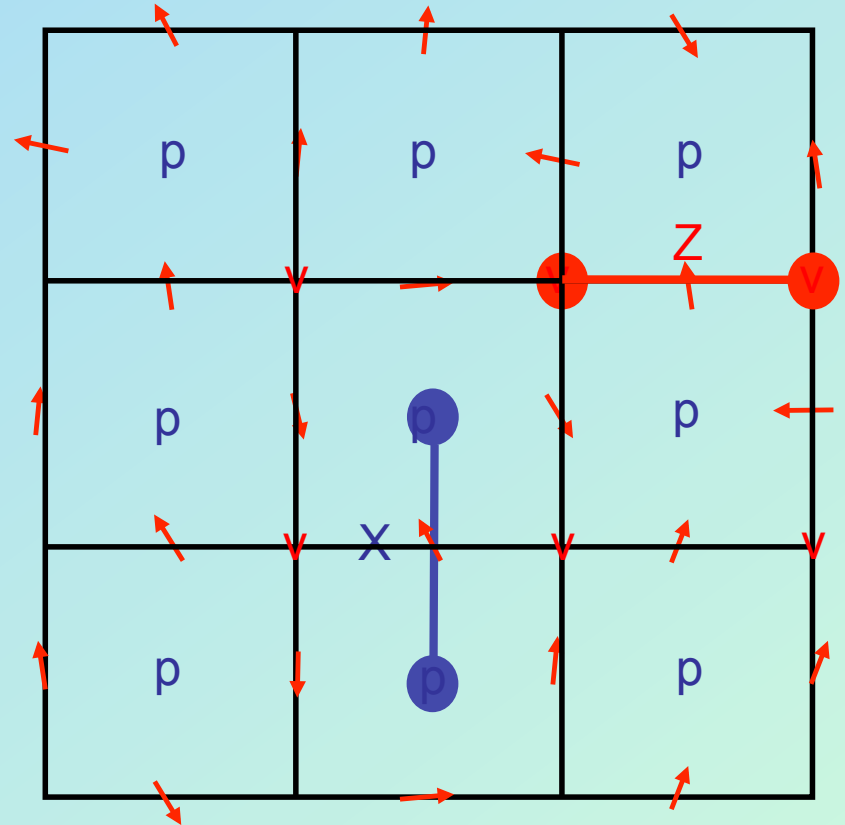
Fusion rules:

$$e \times e = 1, m \times m = 1, \varepsilon \times \varepsilon = 1$$

$$e \times m = \varepsilon, e \times \varepsilon = m, m \times \varepsilon = e$$

Fusion moves: F are trivial

Braiding moves R: $R_{em}^{\varepsilon} = i, R_{\varepsilon\varepsilon}^1 = -1$



Toric Code: Encoding

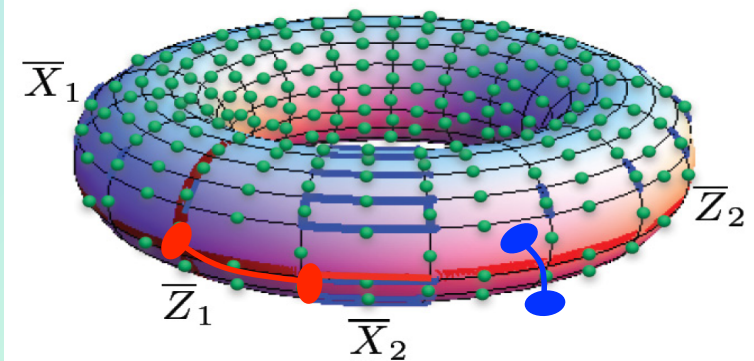
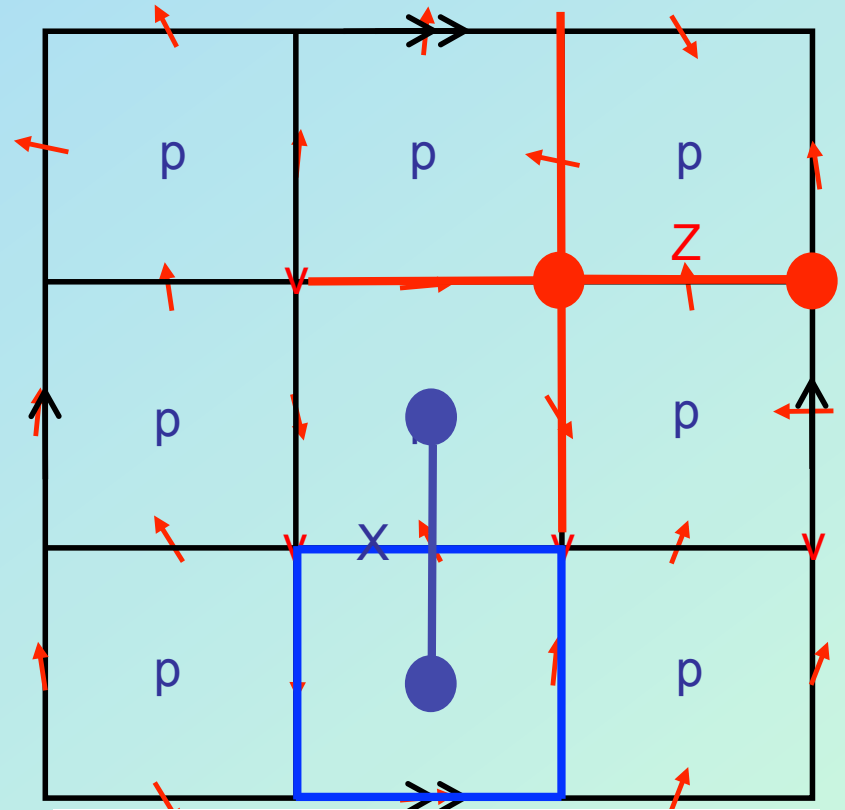
Toric code as a quantum error correcting code.

Consider periodic boundary conditions: TORUS of size L

Syndrome: Anyons

Error correction: detect anyons/errors and connect shortest distance between the same type of anyons.

$$\prod_v X_{v1} X_{v2} X_{v3} X_{v4} = \prod_p Z_{p1} Z_{p2} Z_{p3} Z_{p4} = 1$$



Toric Code: Encoding

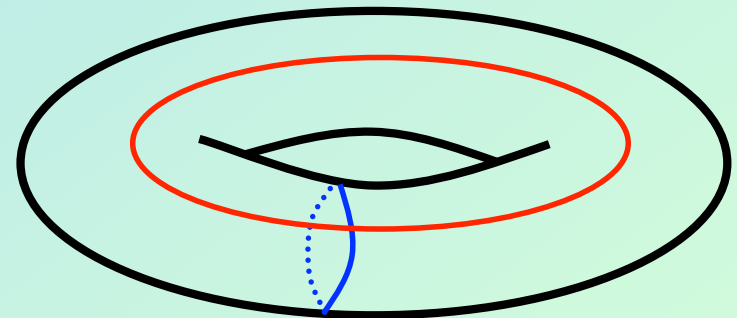
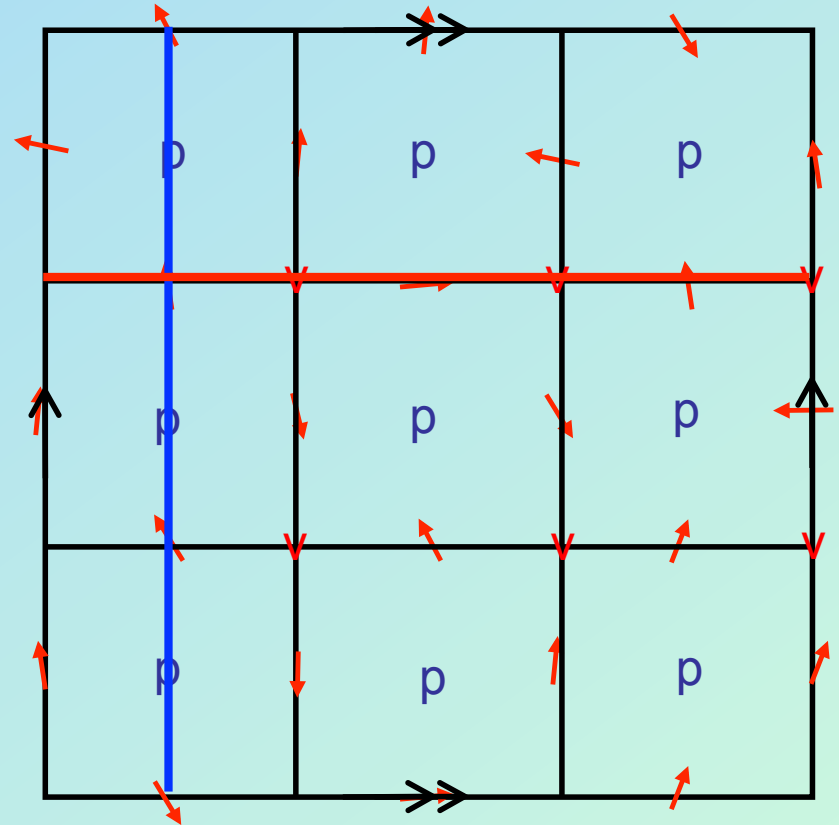
Toric code as a quantum error correcting code.

Consider periodic boundary conditions: TORUS of size L

Syndrome: Anyons

Error correction: detect anyons/errors and connect shortest distance between the same type of anyons.

Logical Gates: non-trivial loops



Toric Code: Encoding

Logical Space and Gates

$$|\Psi_1\rangle$$

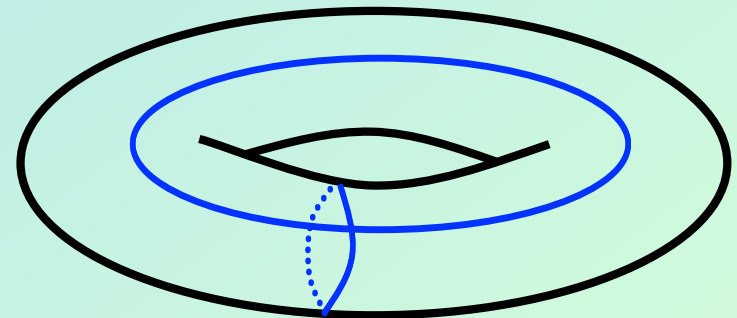
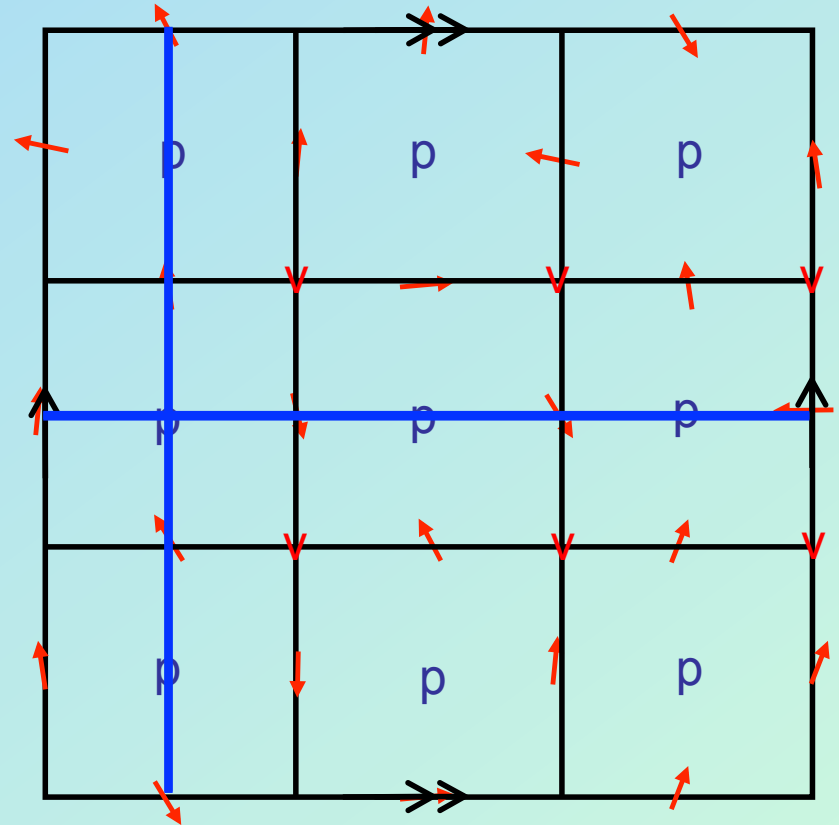
$$|\Psi_2\rangle = C_X^1 |\Psi_1\rangle$$

$$|\Psi_3\rangle = C_X^2 |\Psi_1\rangle$$

$$|\Psi_4\rangle = C_X^2 C_X^1 |\Psi_1\rangle$$

Can store two qubits and perform Clifford group operations!

Higher genus, g , stores $2g$ qubits.



Quantum Double Models

Toric Code is an example of quantum double models.

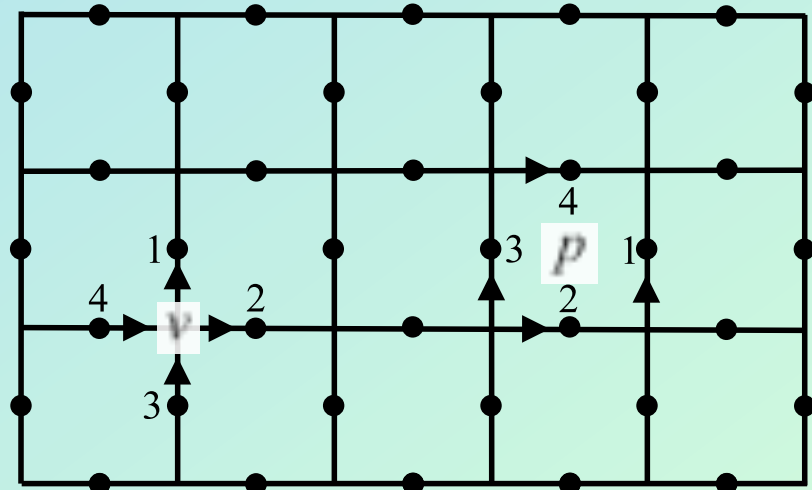
Corresponding group $Z_2 = \{1, e; e^2 = 1\}$ that gives rise to qubit states $|1\rangle, |e\rangle$.

Imagine a general finite group $G = \{g_1, g_2, \dots, g_d\}$ and the corresponding qudit with states $|g_i\rangle, i=1, \dots, d$.

Consider a qudit positioned at each edge of a square lattice.

Define orientation on the lattice:

Upwards and Rightwards



Quantum Double Models

Define operators:

$$L_+^g |z\rangle = |gz\rangle, \quad L_-^g |z\rangle = |zg^{-1}\rangle, \quad T_+^h |z\rangle = \delta_{h,z} |z\rangle, \quad T_-^h |z\rangle = \delta_{h^{-1},z} |z\rangle$$

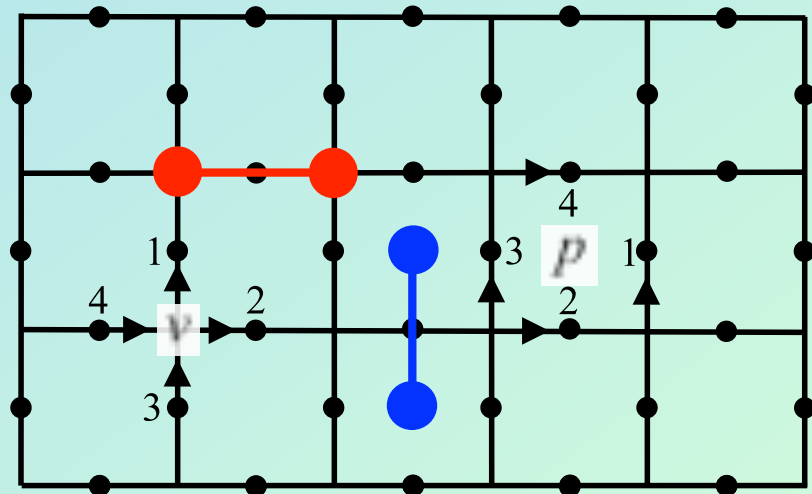
$$A(v) = \frac{1}{|G|} \sum_{g \in G} L_+^g(e_1) L_+^g(e_2) L_-^g(e_3) L_-^g(e_4), \quad B(p) = \sum_{h_1 \dots h_4 = 1} T_-^{h_1}(e_1) T_-^{h_2}(e_2) T_+^{h_3}(e_3) T_+^{h_4}(e_4)$$

Hamiltonian and ground state:

$$H = - \sum_v A(v) - \sum_p B(p)$$

$$A(v) |\xi\rangle = |\xi\rangle$$

$$B(p) |\xi\rangle = |\xi\rangle$$



Quantum Double Models

This is also an error correcting code defined from the stabilizer formalism.

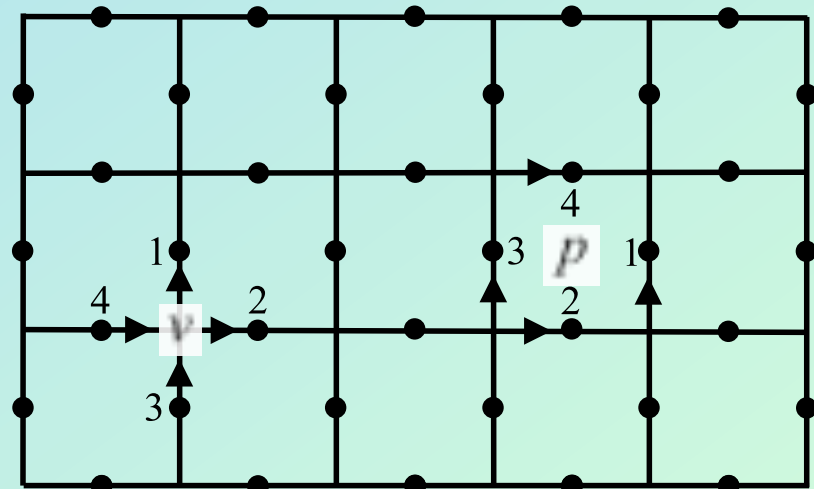
The syndromes are **anyons**, Abelian or non-Abelian, with the corresponding fusion rules, B and F matrices.

These properties can be explicitly determined.

Examples: $D(\mathbb{Z}_2)$, $D(\mathbb{Z}_2 \times \mathbb{Z}_2)$,

$D(S_3)$

$S_3 = \{1, x, y, y^2, xy, xy^2; x^2=1, y^3=1\}$



Quantum Double Models

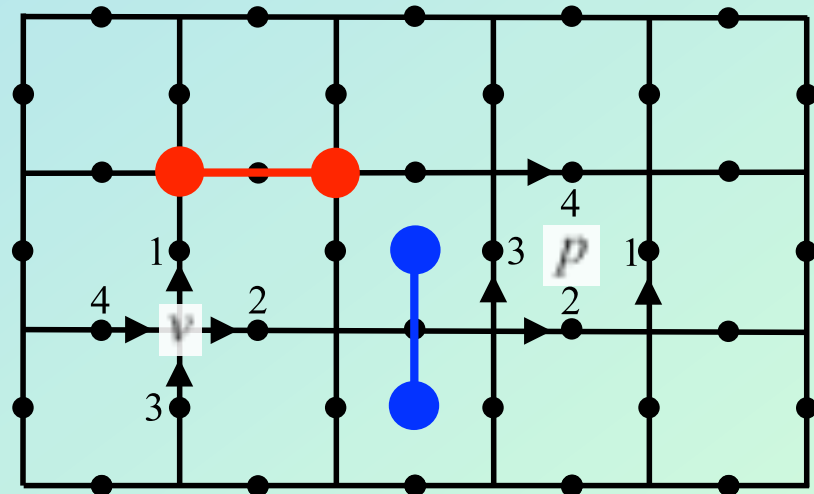
This is also an error correcting code defined from the stabilizer formalism.

The syndromes are **anyons**, Abelian or non-Abelian, with the corresponding fusion rules, B and F matrices.

Information can be encoded in the fusion space of non-Abelian anyons and manipulated by braiding them.

Realizations:

Josephson junctions, photons, optical lattices,...



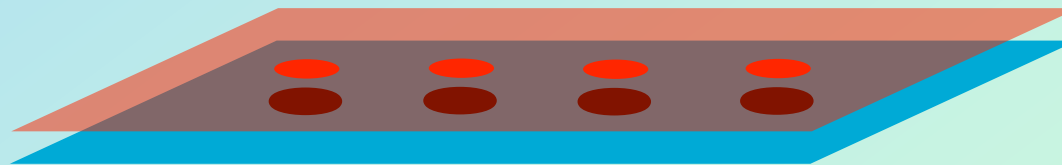
From Abelian to Nonabelian

Consider the toric code with higher **genus** surface or with **punctures**: encoding Hilbert space becomes larger.



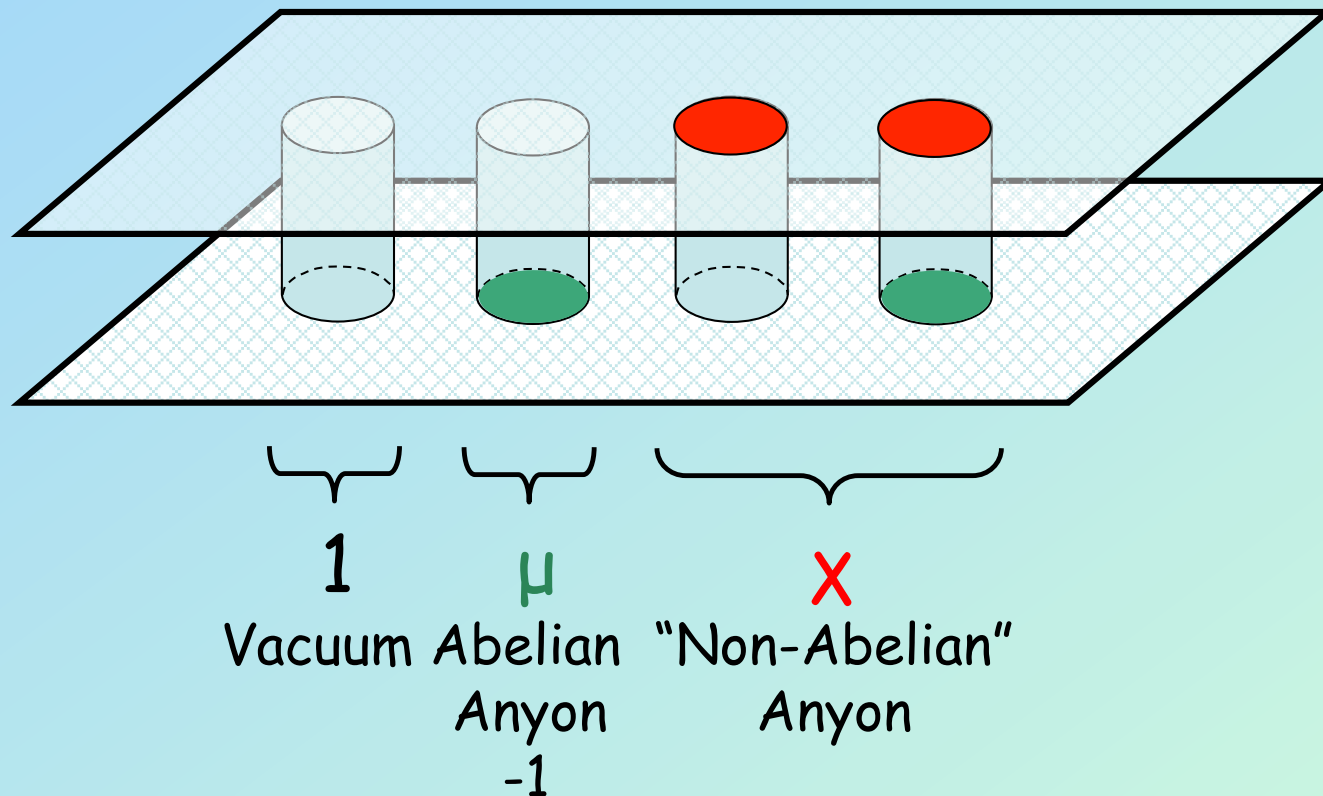
Punctures are *better* for storing and manipulating info.

If you do not like punctures:

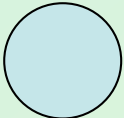



From Abelian to Nonabelian


- The scheme: $D(Z_2 \times Z_2)$ [or $D(S_3)$]
similar to two toric codes




Hamiltonian
bird's eye view:

Empty 

Lower only 

Upper only 

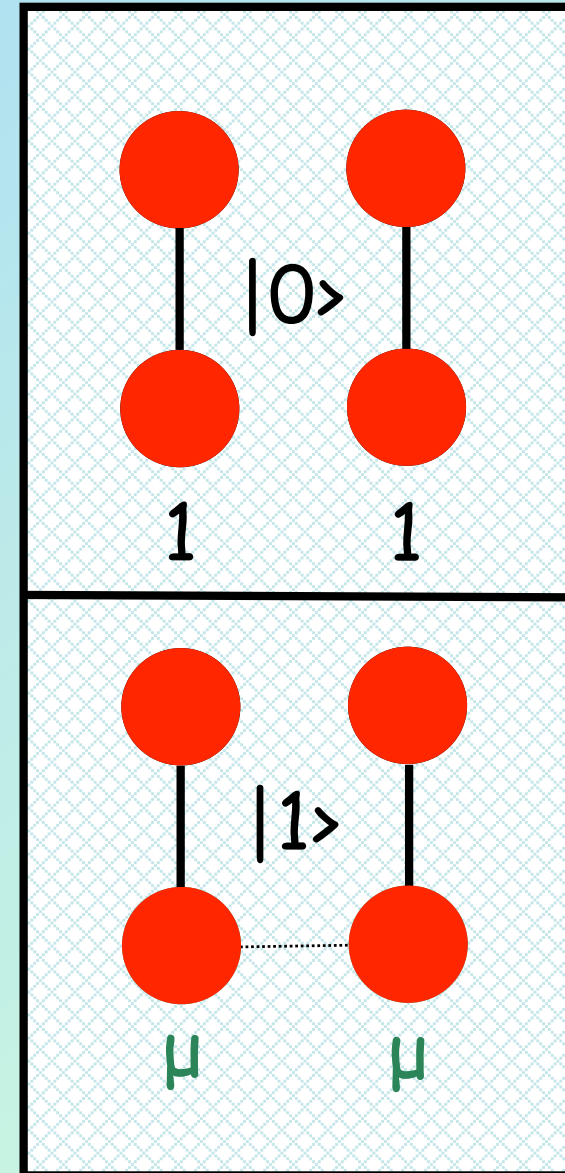
Upper and Lower 

From Abelian to Nonabelian

- Encode information in **fusion channels**:

$$\mu \times \mu = 1, \quad \chi \times \chi = 1 + \mu$$

- Qubit needs four anyons
- Logical $|0\rangle$ when each pair fuses to the vacuum 1
- Logical $|1\rangle$ when each pair fuses to μ
- $1, \mu$ indistinguishable to local operations when dressed with χ
- **Measurement by fusion**



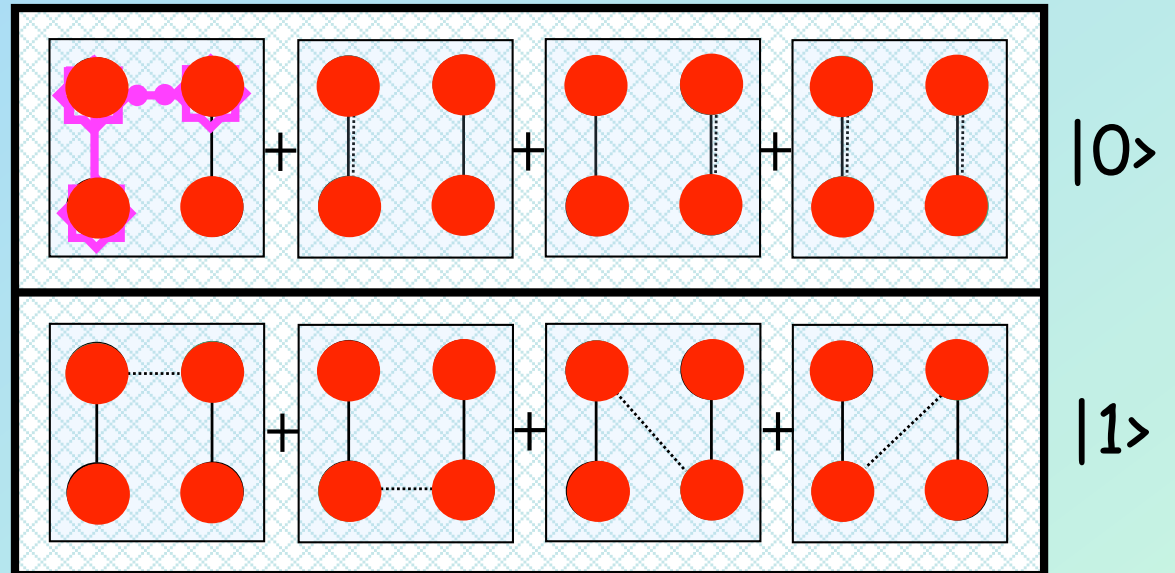
From Abelian to Nonabelian

- Fault-tolerance

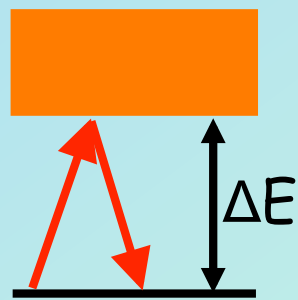
- Phase flips ●
- Bit flips ●

by non-local operators only

→ topo. protection



Energy gap present even during gate operations



- Redundancy and non-locality protects against **virtual transitions**
- Braiding is only Abelian.

Summary

- **Quantum Double models:**
 - Toric Code
 - Abelian encoding and quantum computation
 - Non-Abelian models
- **Degenerate encoding states**
- **Energy gap above encoding space**
- **Manipulations of code space:**

