

Bogoliubov Theory for Dipolar Bose Gas at Low Temperatures

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Bogoliubov theory

- **Hamilton operator:**

$$\hat{H} = \sum_{\mathbf{k}} \left(\frac{\hbar^2 \mathbf{k}^2}{2m} + \hbar \mathbf{w} \mathbf{k} - \mu \right) \hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}} + \frac{1}{2V} \sum_{\mathbf{k}, \mathbf{k}', \mathbf{q}} V_{\mathbf{q}} \hat{b}_{\mathbf{k}-\mathbf{q}}^\dagger \hat{b}_{\mathbf{k}'+\mathbf{q}}^\dagger \hat{b}_{\mathbf{k}'} \hat{b}_{\mathbf{k}}$$

- **C-number approximation:** $N_0 \gg 1$ $\hat{b}_0 \approx \hat{b}_0^\dagger \approx \sqrt{N_0}$

$$\hat{H} = V n_0 \left(\frac{n_0}{2} V_0 - \mu \right) + \frac{1}{2} \sum_{\mathbf{k} \neq 0} \left[\frac{\hbar^2 \mathbf{k}^2}{2m} - \mu + n_0 (V_0 + V_{\mathbf{k}}) \right] \left(\hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}} + \hat{b}_{-\mathbf{k}}^\dagger \hat{b}_{-\mathbf{k}} \right) + \frac{1}{2} \sum_{\mathbf{k} \neq 0} \hbar \mathbf{w} \mathbf{k} \left(\hat{b}_{\mathbf{k}}^\dagger \hat{b}_{\mathbf{k}} - \hat{b}_{-\mathbf{k}}^\dagger \hat{b}_{-\mathbf{k}} \right) + \frac{n_0}{2} \sum_{\mathbf{k} \neq 0} V_{\mathbf{k}} \left(\hat{b}_{\mathbf{k}} \hat{b}_{-\mathbf{k}} + \hat{b}_{\mathbf{k}}^\dagger \hat{b}_{-\mathbf{k}}^\dagger \right) + \dots$$

- **Bogoliubov transformation:** $\hat{B}_{\mathbf{k}} = u_{\mathbf{k}} \hat{b}_{\mathbf{k}} + v_{\mathbf{k}} \hat{b}_{-\mathbf{k}}^\dagger$, $\hat{B}_{\mathbf{k}}^\dagger = u_{\mathbf{k}}^* \hat{b}_{\mathbf{k}}^\dagger + v_{\mathbf{k}}^* \hat{b}_{-\mathbf{k}}$

$$u_{\mathbf{k}} = \frac{E_{\mathbf{k}} + \epsilon_{\mathbf{k}}}{2E_{\mathbf{k}}\epsilon_{\mathbf{k}}}, \quad v_{\mathbf{k}} = \frac{E_{\mathbf{k}} - \epsilon_{\mathbf{k}}}{2E_{\mathbf{k}}\epsilon_{\mathbf{k}}}, \quad E_{\mathbf{k}} = \sqrt{\epsilon_{\mathbf{k}}^2 + 2n_0 V_{\mathbf{k}}\epsilon_{\mathbf{k}}}, \quad \epsilon_{\mathbf{k}} = \frac{\hbar^2 \mathbf{k}^2}{2m} - \mu + n_0 V_0$$

- **Diagonalized Hamiltonian:**

$$\hat{H} = V n_0 \left(\frac{n_0}{2} V_0 - \mu \right) + \frac{1}{2} \sum_{\mathbf{k} \neq 0} (E_{\mathbf{k}} - \epsilon_{\mathbf{k}} - n_0 V_{\mathbf{k}}) + \sum_{\mathbf{k} \neq 0} (E_{\mathbf{k}} + \hbar \mathbf{w} \mathbf{k}) \hat{B}_{\mathbf{k}}^\dagger \hat{B}_{\mathbf{k}} + \dots$$

- **Effective potential:**

$$V_{\text{eff}} = V n_0 \left(\frac{n_0}{2} V_0 - \mu \right) + \frac{1}{2} \sum_{\mathbf{k} \neq 0} (E_{\mathbf{k}} - \epsilon_{\mathbf{k}} - n_0 V_{\mathbf{k}}) + \frac{1}{\beta} \sum_{\mathbf{k} \neq 0} \ln \left(1 - e^{-\beta E_{\mathbf{k}}} \right) + \dots$$

Landau theory

- **Quantum and thermal fluctuations:** smallness parameter η

$$V_{\text{eff}}(n_0) = V \left(-\mu n_0 + \frac{n_0^2}{2} V_0 \right) + \frac{\eta}{2} \sum_{\mathbf{k} \neq 0} \left[E_{\mathbf{k}} - \left(\frac{\hbar^2 \mathbf{k}^2}{2m} - \mu + n_0 (V_0 + V_{\mathbf{k}}) \right) \right] + \frac{\eta}{\beta} \sum_{\mathbf{k} \neq 0} \ln \left(1 - e^{-\beta E_{\mathbf{k}}} \right) + \dots$$

- **Condensate density:** $\frac{\partial V_{\text{eff}}(n_0)}{\partial n_0} = 0 \implies n_0 = n_0^{(0)} + \eta n_0^{(1)} + \dots$

- **Free energy:** $F(T, V, \mu) = V_{\text{eff}} \left(n_0^{(0)} + \eta n_0^{(1)} + \dots \right)$

$$F(T, V, \mu) = -\frac{V \mu^2}{2 V_0} + \frac{\eta}{2} \sum_{\mathbf{k} \neq 0} \left[\sqrt{\left(\frac{\hbar^2 \mathbf{k}^2}{2m} \right)^2 + 2\mu \frac{V_{\mathbf{k}} \hbar^2 \mathbf{k}^2}{2m}} - \frac{\hbar^2 \mathbf{k}^2}{2m} - \mu \frac{V_{\mathbf{k}}}{V_0} \right] + \frac{\eta}{\beta} \sum_{\mathbf{k} \neq 0} \ln \left(1 - e^{-\beta \sqrt{\left(\frac{\hbar^2 \mathbf{k}^2}{2m} \right)^2 + 2\mu \frac{V_{\mathbf{k}} \hbar^2 \mathbf{k}^2}{2m}}} \right) + \dots$$

- **Superfluid density:** $(n_s)_{ij} = n \delta_{ij} - (n_n)_{ij}$

$$p_i = -\frac{\partial F}{\partial w_i} = mV (n_n)_{ij} w_j + \dots, \quad (n_n)_{ij} = \eta \frac{\hbar^2 \beta}{mV} \sum_{\mathbf{k} \neq 0} \frac{e^{\beta E_{\mathbf{k}}}}{(e^{\beta E_{\mathbf{k}}} - 1)^2} k_i k_j + \dots$$

Dipolar Bose gas

- **Interaction potential:** $C_{dd} = \mu_0 m^2$ (atomic origin), $C_{dd} = 4\pi d^2$ (molecular origin)

$$V(\mathbf{x}) = g \delta(\mathbf{x}) + \frac{C_{dd}}{4\pi |\mathbf{x}|^3} (1 - 3 \cos^2 \theta)$$

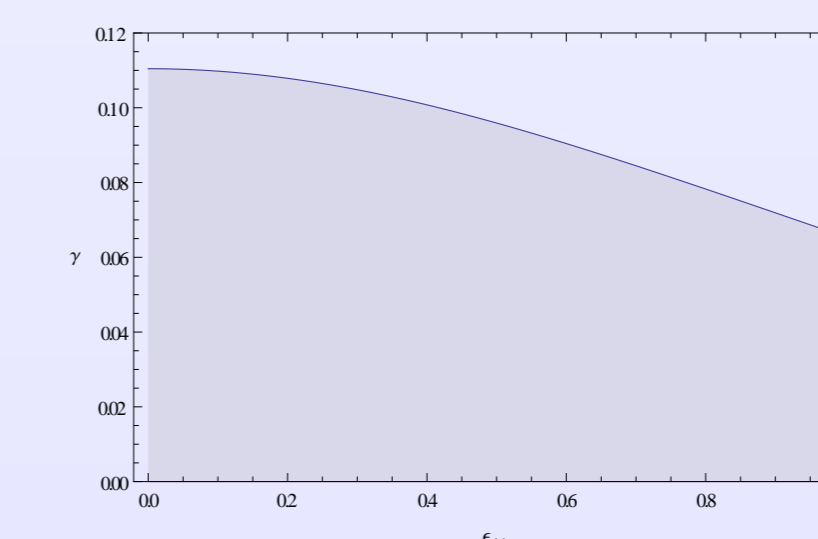
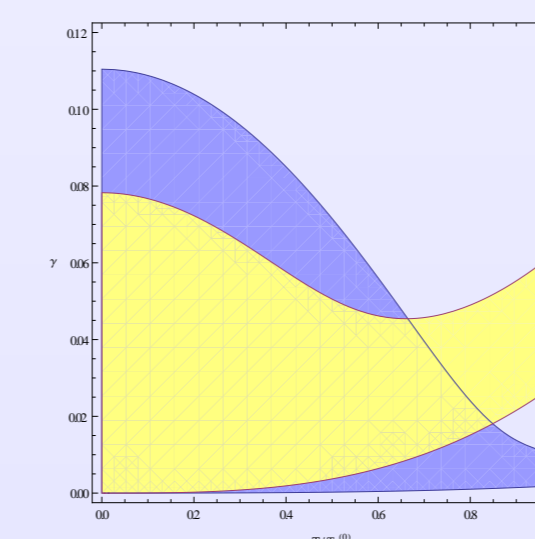
$$V_{\mathbf{k}} = g [1 + \epsilon_{dd} (3 \cos^2 \theta - 1)], \quad \epsilon_{dd} = \frac{C_{dd}}{3g}$$

- **Condensate depletion:**

$$\Delta n_0 = \eta \frac{8}{3\sqrt{\pi}} (na)^{\frac{3}{2}} I \left(\epsilon_{dd}, \frac{3}{2} \right) + \eta \frac{\pi^{\frac{3}{2}}}{6\lambda_{\text{dB}}^4} n^{-\frac{1}{2}} a^{-\frac{1}{2}} I \left(\epsilon_{dd}, -\frac{1}{2} \right) - \eta \frac{\pi^{\frac{7}{2}}}{960\lambda_{\text{dB}}^8} n^{-\frac{5}{2}} a^{-\frac{5}{2}} I \left(\epsilon_{dd}, -\frac{5}{2} \right) + \dots, \quad \lambda_{\text{dB}} = \sqrt{\frac{2\pi \hbar^2}{mk_B T}}$$

$$I(\epsilon_{dd}, \alpha) = \int_0^1 du [1 + \epsilon_{dd} (3u^2 - 1)]^\alpha$$

- **Validity range:** blue: $\epsilon_{dd} = 0$, yellow: $\epsilon_{dd} = 0.8$



$$0 \leq \frac{\Delta n_0}{n} \leq \frac{1}{2} \quad \begin{matrix} T=0 \\ \implies \\ \gamma = na^3 \end{matrix} \quad 0 \leq \gamma(\epsilon_{dd}) \leq \left[\frac{3\sqrt{\pi}}{16I(\epsilon_{dd}, \frac{3}{2})} \right]^2$$

- **Superfluid density:** $(n_n)_{\parallel} \leq (n_n)_{\perp}$

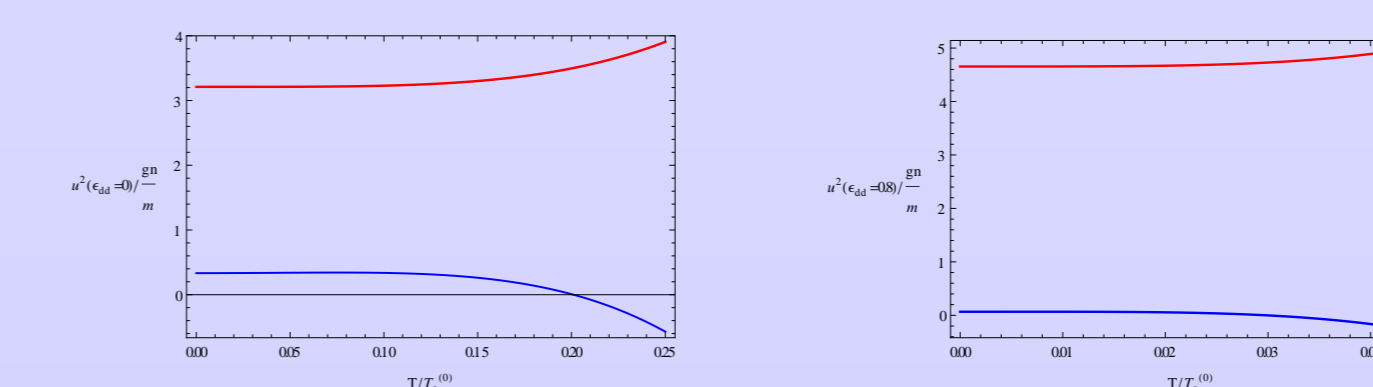
Sound velocities

- **Anisotropic Landau-Khalatnikov two-fluid model [1]**

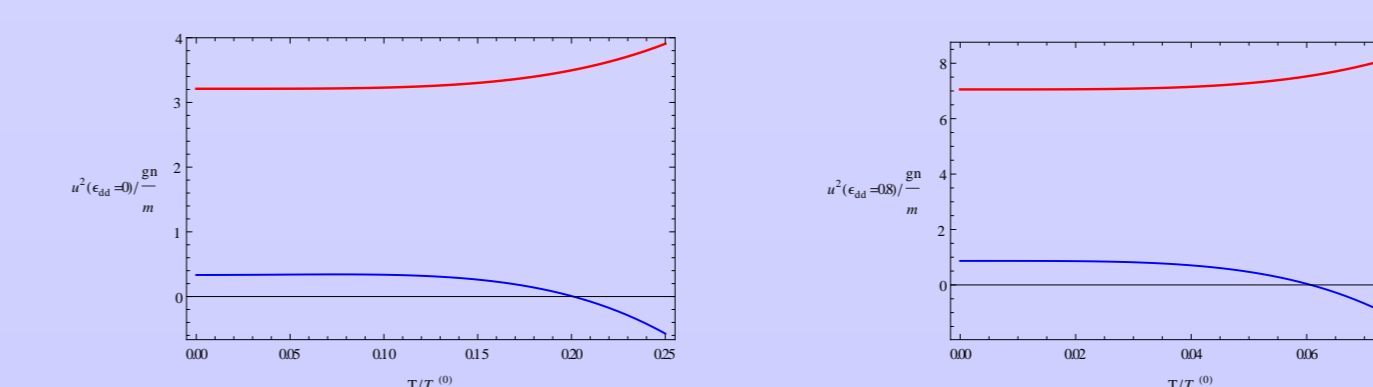
- **Long-wave length dispersion relations:**

$$\omega_{1/2}^2 = \begin{pmatrix} k \sin \vartheta & k \cos \vartheta \\ 0 & u_{\parallel}^2 \end{pmatrix} \begin{pmatrix} u_{\perp}^2 & 0 \\ 0 & u_{\parallel}^2 \end{pmatrix} \begin{pmatrix} k \sin \vartheta \\ k \cos \vartheta \end{pmatrix}$$

- **Sound velocities perpendicular to dipoles ($\gamma = 0.06$):**



- **Sound velocities parallel to dipoles ($\gamma = 0.06$):**



first sound: pressure and density wave (red), second sound: entropy and temperature wave (blue)