



# Critical Properties of Homogeneous Bose-Einstein Condensates



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## 1. Variational Perturbation Theory

- Bose gas with weak repulsive two-particle interaction:

$$\mathcal{A}[\psi^*, \psi] = \int_0^{\hbar\beta} d\tau \int d^3x \left\{ \psi^*(\mathbf{x}, \tau) \left( \hbar \partial_\tau - \frac{\hbar^2 \nabla^2}{2M} \right) \psi(\mathbf{x}, \tau) - \mu |\psi(\mathbf{x}, \tau)|^2 + \frac{2\pi\hbar^2 a_s}{M} |\psi(\mathbf{x}, \tau)|^4 \right\}$$

### 1.1 Dimension $D = 3 + 1$

- Goal:** Extension of three-loop calculation for grand-canonical potential [1] to four-loops [2,3]

$$\Omega^{(4)} = 4 \text{ (chain)} + \frac{8}{3} \text{ (triangle)} + 4 \text{ (square)} + \frac{4}{3} \text{ (pentagon)} + \frac{1}{3} \text{ (hexagon)}$$

- Goal:** Variational calculation of critical temperature [4]

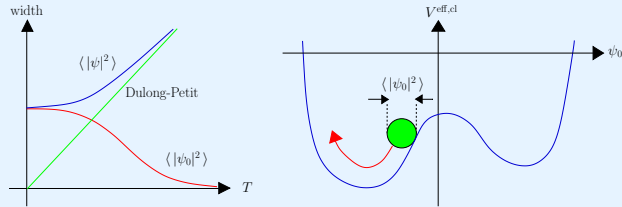
$$\frac{\Delta T_c}{T_c^{(0)}} = c_1 a_s n^{1/3} + [c_2 \ln(a_s n^{1/3}) + c_2] a_s^2 n^{2/3} + \mathcal{O}(a_s^3 n)$$

- Goal:** Arbitrary interaction potential [3]

$$\frac{\Delta T_c}{T_c^{(0)}} = \frac{M n^{1/3}}{3\pi\hbar^2 \zeta^{4/3}(3/2)} \sum_{k=1}^{\infty} \left[ \frac{-\pi n^{2/3}}{\zeta^{2/3}(3/2)} \right]^k \times \int d^3x V^{(\text{int})}(\mathbf{x})(\mathbf{x}^2)^k \sum_{l=1}^k \frac{\zeta(l+1/2)\zeta(3/2-k-l)}{l!(k-l)!}$$

- Goal:** Develop effective classical field theory [5,6]

$$Z = \int \mathcal{D}\psi_0^*(\mathbf{x}) \int \mathcal{D}\psi_0(\mathbf{x}) e^{-\mathcal{A}^{\text{eff}, \text{cl}}[\psi_0^*, \psi_0]/\hbar}$$



- Goal:** Justification of intuitive matching procedure [7]
- Goal:** Improvement of effective action calculation [8,9]
- Goal:** Condensate density, superfluid density, specific heat

### 1.2 Dimension $D = 3$

- Shift of critical temperature in dilute limit:

$$\frac{\Delta T_c}{T_c^{(0)}} = -\frac{2}{3} \frac{\Delta n}{n} = c_1 a_s n^{1/3}$$

- Calculation of  $\Delta n$  from classical limit of field theory:

$$\mathcal{A}[\Psi] = \int d^3x \left[ \frac{1}{2} |\nabla\Psi|^2 + \frac{1}{2} m^2 |\Psi|^2 + \frac{u}{4!} |\Psi|^4 \right]$$

$$\Psi_1 + i\Psi_2 = \sqrt{\frac{\hbar^3}{M k_B T}} \psi_0, \quad m^2 = -\frac{2M\mu}{\hbar^2}, \quad u = \frac{48\pi a_s M k_B T}{\hbar^3}$$

- Five-loop perturbation expansion [10]:  $c_1 = \sum_{l=-1}^3 c_1^{(l)} \left( \frac{u}{m_r} \right)^l$

critical limit  $m_r \rightarrow 0 \iff$  strong-coupling problem

- Applying variational perturbation theory [6,11] yields resummation of  $c_1$ -series [12]:  $c_1 \approx 0.91$

$1/N$ -expansion [13]:  $c_1 \approx 1.71$ , Monte-Carlo result [14]:  $c_1 \approx 1.29$

- Goal:** Six-loop variational calculation

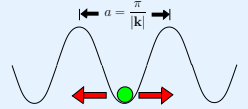
## 2. Mott Insulator Transition in Optical Lattices

- Periodic potential:  $V(\mathbf{x}) = V_0 \sum_{i=1}^3 \sin^2(k_i x_i)$ ,  $E_r = \frac{\hbar^2 \mathbf{k}^2}{2M}$   
Experimentally [15],  $V_0$  can be raised to  $22E_r$

- Goal:** Effective  $|\psi|^4$ -field theory for lowest energy band in Bose-Hubbard model [16]

$$\frac{\hbar^2 \mathbf{q}^2}{2M_{\text{eff}}} \cong 2J \sum_{i=1}^3 (1 - \cos q_i a)$$

$$\frac{4\pi\hbar^2 a_s^H}{M_{\text{eff}}} \cong \sqrt{\frac{8}{\pi}} |\mathbf{k}| a_s E_r \left( \frac{V_0}{E_r} \right)^{3/4}$$



### 2.1 Goldstone Boson Model

- Decompose field  $\psi(\mathbf{x}, \tau) = [\sqrt{n_0} + \rho(\mathbf{x}, \tau)] e^{i\gamma(\mathbf{x}, \tau)}$  and integrate out density fluctuations  $\rho$

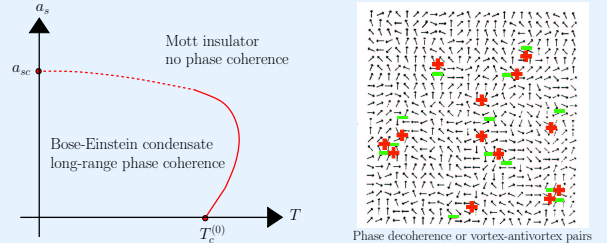
- Goal:** Study phase fluctuations in hydrodynamic limit

$$\mathcal{L}^{\text{hydr}} \approx \frac{\hbar^2 n_0}{2M_{\text{eff}}} \left[ (\nabla\gamma)^2 + \frac{1}{c^2} \partial_\tau \gamma \left( 1 - \frac{\hbar^2 \nabla^2}{4M_{\text{eff}}^2 c^2} \right)^{-1} \partial_\tau \gamma \right]$$

with Bogoliubov spectrum

$$\epsilon(\mathbf{p}) = \sqrt{c^2 \mathbf{p}^2 + \left( \frac{\mathbf{p}^2}{2M_{\text{eff}}} \right)^2}, \quad c = \sqrt{\frac{4\pi\hbar^2 a_s n}{M_{\text{eff}}^2}}$$

- Goal:** Determine whole phase diagram



### 2.2 Phase Decoherence

- Goal:** Study new  $O(N)$ -symmetric nonlinear  $\sigma$ -model

$$\mathcal{L}^\sigma = \frac{\hbar^2 n_0}{2M_{\text{eff}}} \left[ (\nabla\mathbf{m})^2 + \frac{1}{c^2} \partial_\tau \mathbf{m} \left( 1 - \frac{\hbar^2 \nabla^2}{4M_{\text{eff}}^2 c^2} \right)^{-1} \partial_\tau \mathbf{m} \right] + \lambda(\mathbf{m}^2 - 1)$$

- Goal:** Extend classical criterion for dephasing transition [17] to quantum regime
- Goal:** Characterize Mott insulator transition by energy gap [18]

$$\epsilon(\mathbf{p}) = \sqrt{c^2 \mathbf{p}^2 + \left( \frac{\mathbf{p}^2}{2M_{\text{eff}}} \right)^2 + \Delta^2}$$

### 2.3 Disorder Gauge Field Theory

- Goal:** Describe ensembles of vortex lines by disorder field  $\phi$

- Goal:** Introduce auxiliary vector potential  $\mathbf{a}$  by duality transformation

- Goal:** Develop Meissner-Higgs theory [19,20] whose Feynman diagrams represent pictures of vortex lines

$$\mathcal{L}^{\text{GL}} = \frac{1}{2} \left[ (\nabla \times \mathbf{a})_{\parallel} \left( 1 - \frac{\hbar^2 \nabla^2}{4M_{\text{eff}}^2 c^2} \right) (\nabla \times \mathbf{a})_{\parallel} + (\nabla \times \mathbf{a})_{\perp}^2 \right] + \left| \left( \nabla - \frac{2\pi i}{\hbar} \sqrt{\frac{n_0}{M_{\text{eff}}}} \mathbf{a} \right) \phi \right|^2 + m^2 |\phi|^2 + g|\phi|^4$$

- Goal:** Characterize quantum phase transition by proliferation of vortex lines whose cores are Mott insulators [21]



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