On the Dirty Boson Problem

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- **1. Experimental Realizations of Dirty Bosons**
- 2. Theoretical Description of Dirty Bosons
- 3. Huang-Meng Theory (T=0)
- 4. Shift of Condensation Temperature
- **5. Hartree-Fock Mean-Field Theory**
- 6. Summary and Outlook

SFB/TR 12: Symmetries and Universality in Mesoscopic Systems

1.1 Overview of Set-Ups

- Superfluid Helium in Porous Media: (persistence of superfluidity) Reppy *et al.*, PRL **51**, 666 (1983)
- Laser Speckles: (controlled randomness) Billy *et al.*, Nature **453**, 891 (2008)
- Wire Traps: (undesired randomness)
 Krüger *et al.*, PRA **76**, 063621 (2007)
 Fortàgh and Zimmermann, RMP **79**, 235 (2007)
- Localized Atomic Species: (theoretical suggestion) Gavish and Castin, PRL 95, 020401 (2005)
- Incommensurate Lattices: (quasi-randomness)

Roati *et al.*, Nature **453**, 895 (2008)

1.2 Laser Speckles



0 Hz 50 Hz 360 Hz 1200 Hz (iii) O Hz (I) O Hz (ii) 1 700 Hz 0 30 60 1000 H 400 Hz 60 Hz 90 O Hz 200 time (ms) 100 300 0 400

global condensate vanishes

Lye et al., PRL 95, 070401 (2005)

1.3 Wire Trap



Distance: $d = 10 \ \mu \text{m}$

Wire Width: $100 \ \mu m$

Magnetic Field: 10 G, 20 G, 30 G

Deviation: $\Delta B/B \approx 10^{-4}$

Krüger *et al.*, PRA **76**, 063621 (2007) Fortàgh and Zimmermann, RMP **79**, 235 (2007)

2.1 Model System

Action of a Bose Gas:

$$\mathcal{A} = \int_0^{\hbar\beta} d\tau \int d^3x \, \left\{ \psi^* \left[\hbar \frac{\partial}{\partial \tau} - \frac{\hbar^2}{2M} \mathbf{\Delta} + \mathbf{U}(\mathbf{x}) + \mathbf{V}(\mathbf{x}) - \mathbf{\mu} \right] \psi + \frac{\mathbf{g}}{2} \, \psi^{*2} \psi^2 \right\}$$

Properties:

- harmonic trap potential: $\mathbf{U}(\mathbf{x}) = \frac{M}{2}\omega^2 \mathbf{x}^2$
- disorder potential: $V(\mathbf{x})$; bounded from below, i.e. $V(\mathbf{x}) \ge V_0$

$$\overline{V(\mathbf{x}_1)} = 0, \quad \overline{V(\mathbf{x}_1)V(\mathbf{x}_2)} = R(\mathbf{x}_1 - \mathbf{x}_2), \quad \dots$$

- chemical potential: μ
- repulsive interaction:

$$\mathbf{g} = \frac{4\pi\hbar^2 a}{M}$$

• periodic Bose fields: $\psi(\mathbf{x}, \tau + \hbar\beta) = \psi(\mathbf{x}, \tau)$

2.2 Random Potential

Disorder Ensemble Average:

$$\overline{\bullet} = \int \mathcal{D}V \bullet P[V], \quad \int \mathcal{D}V P[V] = 1, \quad P[V < V_0] = 0$$

Assumption:

$$\overline{V(\mathbf{x}_1)} = 0, \quad \overline{V(\mathbf{x}_1)V(\mathbf{x}_2)} = R^{(2)}(\mathbf{x}_1 - \mathbf{x}_2) = \frac{R}{(2\pi\xi^2)^{3/2}} e^{-\frac{(\mathbf{x}_1 - \mathbf{x}_2)^2}{2\xi^2}}$$

Characteristic Functional:

$$\overline{\exp\left\{i\int d^{D}x\,j(\mathbf{x})V(\mathbf{x})\right\}}$$
$$=\exp\left\{\sum_{n=2}^{\infty}\frac{i^{n}}{n!}\int d^{D}x_{1}\cdots\int d^{D}x_{n}\,R^{(n)}(\mathbf{x}_{1},\ldots,\mathbf{x}_{n})\,j(\mathbf{x}_{1})\cdots j(\mathbf{x}_{n})\right\}$$

2.3 Grand-Canonical Potential

Aim:

$$\Omega = -\frac{1}{\beta} \overline{\ln \mathcal{Z}}$$
$$\mathcal{Z} = \oint D\psi D\psi^* e^{-\mathcal{A}[\psi^*,\psi]/\hbar}$$

Problem:

$$\overline{\ln \mathcal{Z}} \neq \ln \overline{\mathcal{Z}}$$

Solution: Replica Trick

$$\Omega = -\frac{1}{\beta} \lim_{N \to 0} \frac{\overline{\mathcal{Z}^N} - 1}{N}$$

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2.4 Replica Trick

Disorder Averaged Partition Function:

$$\overline{\mathcal{Z}^{N}} = \oint \left\{ \prod_{\alpha'=1}^{N} D^{2} \psi_{\alpha'} \right\} e^{-\sum_{\alpha=1}^{N} \mathcal{A}([\psi_{\alpha}^{*}, \psi_{\alpha}])/\hbar} = \oint \left\{ \prod_{\alpha=1}^{N} D^{2} \psi_{\alpha} \right\} e^{-\mathcal{A}^{(N)}/\hbar}$$

Replicated Action:

$$\mathcal{A}^{(N)} = \int_{0}^{\hbar\beta} d\tau \int d^{D}x \sum_{\alpha=1}^{N} \left\{ \psi_{\alpha}^{*}(\mathbf{x},\tau) \left[\hbar \frac{\partial}{\partial \tau} - \frac{\hbar^{2}}{2M} \Delta + U(\mathbf{x}) - \mu \right] \psi_{\alpha}(\mathbf{x},\tau) \right. \\ \left. + \frac{g}{2} \left| \psi_{\alpha}(\mathbf{x},\tau) \right|^{4} \right\} + \sum_{n=2}^{\infty} \frac{1}{n!} \left(\frac{-1}{\hbar} \right)^{n-1} \int_{0}^{\hbar\beta} d\tau_{1} \cdots \int_{0}^{\hbar\beta} d\tau_{n} \int d^{D}x_{1} \cdots \int d^{D}x_{n} \\ \left. \times \sum_{\alpha_{1}=1}^{N} \cdots \sum_{\alpha_{n}=1}^{N} R^{(n)}(\mathbf{x}_{1},\ldots,\mathbf{x}_{n}) \left| \psi_{\alpha_{1}}(\mathbf{x}_{1},\tau_{1}) \right|^{2} \cdots \left| \psi_{\alpha_{n}}(\mathbf{x}_{n},\tau_{n}) \right|^{2} \right]$$

 \implies Disorder amounts to attractive interaction for n = 2

 \implies Higher-order disorder cumulants negligible in replica limit $N \rightarrow 0$

3.1 Condensate Density

Assumptions:

homogeneous Bose gas:

$$U(\mathbf{x}) = 0$$

 δ -correlated disorder:

 $R(\mathbf{x}) = R\,\delta(\mathbf{x})$

Bogoliubov Theory:

background method:

$$\psi_{\alpha}(\mathbf{x},\tau) = \Psi_{\alpha} + \delta\psi_{\alpha}(\mathbf{x},\tau)$$

replica symmetry:

 $\Psi_{\alpha} = \sqrt{n_0}$

Result:

$$n_0 = n - \frac{8}{3\sqrt{\pi}}\sqrt{a n_0}^3 - \frac{M^2 R}{8\pi^{3/2}\hbar^4}\sqrt{\frac{n_0}{a}}$$

Huang and Meng, PRL **69**, 644 (1992) Falco, Pelster, and Graham, PRA **75**, 063619 (2007)

3.2 Superfluid Density

Galilei Boost:

$$\Delta \mathcal{A} = \int_{0}^{\hbar\beta} d\tau \int d^{3}x \,\psi^{*}(\mathbf{x},\tau) \,\mathbf{u} \,\frac{\hbar}{i} \,\nabla \,\psi(\mathbf{x},\tau)$$
$$d\Omega = -S \,dT - p \,dV - N \,d\mu - \mathbf{p} \,d\mathbf{u}$$
$$\mathbf{p} = -\frac{\partial \Omega(T,V,\mu,\mathbf{u})}{\partial \mathbf{u}}\Big|_{T,V,\mu} = MV \,n_{n} \,\mathbf{u} + \dots$$

Result:
$$n_s = n - n_n = n - \frac{4}{3} \frac{M^2 R}{8\pi^{3/2}\hbar^4} \sqrt{\frac{n_0}{a}}$$

Huang and Meng, PRL **69**, 644 (1992) Falco, Pelster, and Graham, PRA **75**, 063619 (2007) ٠

3.3 Collective Excitations

Hydrodynamic Equation in Trap With Disorder:

$$m \frac{\partial^2}{\partial t^2} \delta n(\mathbf{x}, t) - \boldsymbol{\nabla} \Big[g n_{\mathrm{TF}}(\mathbf{x}) \boldsymbol{\nabla} \delta n(\mathbf{x}, t) \Big]$$
$$= -\boldsymbol{\nabla}^2 \Big[3g n_R(\mathbf{x}) \delta n(\mathbf{x}, t) \Big] - \boldsymbol{\nabla} \left[\frac{4g}{3} n_R(\mathbf{x}) \boldsymbol{\nabla} \delta n(\mathbf{x}, t) \right]$$

 $n_R(\mathbf{x})$: Huang-Meng depletion in trap $n_{\mathrm{TF}}(\mathbf{x}) = \left[\mu - V(\mathbf{x})\right]/g$: Thomas-Fermi density

Violation of Kohn Theorem:

Surface dipole mode

$$(n = 0, l = 1)$$
:
$$\frac{\delta\omega_{dip}(\xi = 0)}{\omega_{dip}} = -\frac{5\pi}{16} \frac{M^2 R}{8\pi^{3/2} \hbar^4 \sqrt{n_{TF}(\mathbf{0})a}}$$

Falco, Pelster, and Graham, PRA 76, 013624 (2007)

3.4 Comparison With Experiment

Typical Values:



 \implies Disorder effect vanishes in laser speckle experiment

Improvement:

laser speckle setup with correlation length $\xi = 1 \ \mu m$

Aspect et al., NJP 8, 165 (2006)

\implies Disorder effect should be measurable

Falco, Pelster, and Graham, PRA 76, 013624 (2007)

3.5 Rederivation of Huang-Meng Depletion

Gross-Pitaevskii Equation:

$$\left\{-\frac{\hbar^2}{2M}\boldsymbol{\Delta} + V(\mathbf{x}) + g\left|\Psi(\mathbf{x})\right|^2\right\}\Psi(\mathbf{x}) = \mu\,\Psi(\mathbf{x})$$

Perturbative Expansion:

$$\Psi(\mathbf{x}) = \sqrt{\frac{\mu}{g}} + \Psi_1(\mathbf{x}) + \Psi_2(\mathbf{x}) + \dots$$

$$\implies \text{Condensate density:} \quad n_0(\mu) = |\overline{\Psi(\mathbf{x})}|^2$$
$$\implies \text{Particle density:} \quad n(\mu) = \overline{|\Psi(\mathbf{x})|^2}$$

Disorder-Induced Depletion:

$$n_0 = n - n \int \frac{d^3k}{(2\pi)^3} \frac{R(\mathbf{k})}{\left(\frac{\hbar^2 \mathbf{k}^2}{2M} + 2gn\right)^2} + \dots \qquad \begin{array}{c} R(\mathbf{k}) = 1 \\ \implies \\ \end{array} \qquad \begin{array}{c} \text{Huang-Meng} \\ \text{depletion} \end{array}$$

Krumnow, von Hase, and Pelster (to be published)

4.1 Earlier Results

trapped Bose gas	homogeneous Bose gas
$T_c^{(0)} = \frac{\hbar\omega_{\rm g}}{k_B} \left[\frac{N}{\zeta(3)}\right]^{1/3}$	$T_c^{(0)} = \frac{2\pi\hbar^2}{k_B M} \left[\frac{n}{\zeta(3/2)}\right]^{2/3}$
$\frac{\Delta T_c}{T_c^{(0)}} = -3.426 \frac{a}{\lambda_c^{(0)}}$ Giorgini et al., PRA 54 , R4633 (1996) Gerbier et al., PRL 92 , 030405 (2004)	$\frac{\Delta T_c}{T_c^{(0)}} = 1.3 an^{1/3}$ Kleinert, MPLB 17 , 1011 (2003) Kastening, PRA 69 , 043613 (2004)
$R(\mathbf{x}) = ?$ $\frac{\Delta T_c}{T_c^{(0)}} = ?$	$R(\mathbf{x}) = R \delta(\mathbf{x})$ $\frac{\Delta T_c}{T_c^{(0)}} = -\frac{M^2 R}{3\pi [\zeta(3/2)]^{2/3} \hbar^2 n^{1/3}}$ Lopatin and Vinokur, PRL 88 , 235503 (2002)

Procedure: $n = n(\mu), \quad \mu \nearrow \mu_c \quad \Rightarrow \quad T_c$

4.2 Our Results



Timmer, Pelster, and Graham, EPL 76, 760 (2006)

5.1 Order Parameters

Definition:

$$\lim_{\substack{|\mathbf{x}-\mathbf{x}'|\to\infty}} \overline{\langle \psi(\mathbf{x},\tau)\psi^*(\mathbf{x}',\tau)\rangle} = n_0$$
$$\lim_{\mathbf{x}-\mathbf{x}'\mid\to\infty} \overline{|\langle \psi(\mathbf{x},\tau)\psi^*(\mathbf{x}',\tau)\rangle|^2} = (n_0+q)^2$$

Note:

q is similar to Edwards-Anderson order parameter of spin-glass theory

Hartree-Fock Mean-Field Theory:

Self-consistent determination of n_0 and q for $R(\mathbf{x} - \mathbf{x}') = R \,\delta(\mathbf{x} - \mathbf{x}')$

Phase Classification:

gas	Bose glass	superfluid
$q = n_0 = 0$	$q > 0, n_0 = 0$	$q > 0, n_0 > 0$

5.2 Hartree-Fock Results

Isotherm: T = const.

Phase Diagram: $\mu = \text{const.}$

disorder strength R = const.



Graham and Pelster, IJBC 19, 2745 (2009)

5.3 Properties of Bose Glass

- Absence of ODLRO and phase coherence $(n_0 = 0)$, presence of long-range glassy order (q > 0)
- Localization length of excitations: $\xi = \frac{2\pi\hbar^4}{M^2R}$
- Life time (width) of excitations: $\frac{M\xi}{\hbar |\mathbf{k}|}$
- Superfluid density: Localized part of condensate *q* contributes to normal but not to superfluid density
- Note: Localized condensation (q > 0) without global condensate, but not global condensation $n_0 > 0$ without localized part

6.1 Summary and Outlook

• Frozen Disorder Potential:

arises both artificially (laser speckles) or naturally (wire trap)

• Bosons:

local condensates in minima + global condensate + thermally excited

• Localization Versus Transport:

disorder reduces superfluidity

• Phase Diagram:

yet unknown for strong disorder

Navez, Pelster, and Graham, APB 86, 395 (2007)

• Disordered Bosons in Lattice:

Bose Glass versus Mott phase

6.2 Quantum Phase Transition

Gross-Pitaevskii Equation:

$$\left\{-\frac{\hbar^2}{2M}\boldsymbol{\Delta} + V(\mathbf{x}) + g\Psi^2(\mathbf{x})\right\}\Psi(\mathbf{x}) = \mu\,\Psi(\mathbf{x})$$

Nonperturbative Approach:

Gaussian approximation for random fields $V(\mathbf{x})$ and $\Psi(\mathbf{x})$ $\frac{\overline{V(\mathbf{x})} = 0, \quad \overline{V(\mathbf{x})V(\mathbf{x}')} = R(\mathbf{x} - \mathbf{x}'), \quad \overline{\Psi(\mathbf{x})} = \sqrt{n_0} \\
\overline{\Psi(\mathbf{x})\Psi(\mathbf{x}')} = G_{\Psi\Psi}(\mathbf{x} - \mathbf{x}'), \quad \overline{V(\mathbf{x})\Psi(\mathbf{x}')} = G_{V\Psi}(\mathbf{x} - \mathbf{x}') \\
\Rightarrow \text{Self-consistency equations} \\
\text{Result:} \\
n_0(r) = n\left(1 - \frac{r^2}{2} - \frac{r}{2}\sqrt{r^2 + 4}\right) \\
r_c = \frac{1}{\sqrt{2}}, r = \sqrt{\frac{\pi}{2an}} R\left(\frac{M}{2\pi\hbar^2}\right)^{3/2}$

Krumnow, von Hase, and Pelster (to be published)

6.3 Disordered Bosons in Lattice

Bose-Hubbard Hamilton Operator:

$$\begin{split} \hat{H}_{\rm BH} &= -t \sum_{\langle \mathbf{i}, \mathbf{j} \rangle} \hat{a}_{\mathbf{i}}^{\dagger} \hat{a}_{\mathbf{j}} + \sum_{\mathbf{i}} \left[\frac{U}{2} \hat{n}_{\mathbf{i}} (\hat{n}_{\mathbf{i}} - 1) + (\boldsymbol{\epsilon}_{\mathbf{i}} - \mu) \, \hat{n}_{\mathbf{i}} \right] \\ - &= \prod_{\mathbf{i}} \int_{-\infty}^{+\infty} \bullet p(\boldsymbol{\epsilon}_{\mathbf{i}}) \, d\boldsymbol{\epsilon}_{\mathbf{i}} \ , \qquad p(\boldsymbol{\epsilon}_{\mathbf{i}}) = \begin{cases} 1/\Delta & ; \ \boldsymbol{\epsilon}_{\mathbf{i}} \in [-\Delta/2, +\Delta/2] \\ 0 & ; \ \text{otherwise} \end{cases}$$

Mean-Field Phase Diagram:



Krutitsky, Pelster, and Graham, NJP 8, 187 (2006)

Stochastic Mean-Field Theory: Bissbort and Hofstetter, EPL 86, 50007 (2009)