



INSTITUTE OF PHYSICS
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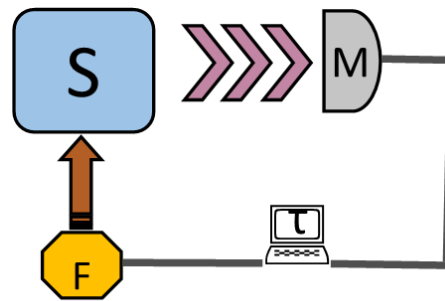
Dissipative Two-Mode Tavis-Cummings Model with Time-Delayed Feedback Control

PRA 92, 063832 (2015)

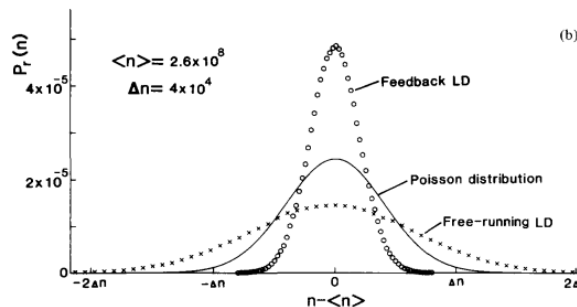
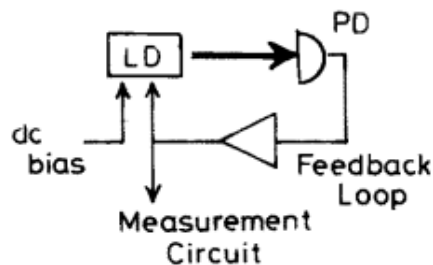
WASSILIJ KOPYLOV, MILAN RADONJIĆ,
TOBIAS BRANDES, ANTUN BALAZŽ
AND AXEL PELSTER

Why Feedback?

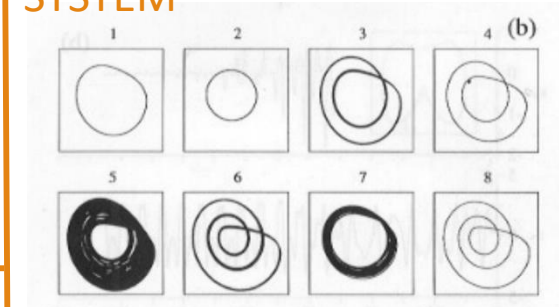
W. Just, A. Pelster et al. , Phil. Trans. R. Soc. A **368** (2010)



SUB-POISSONIAN STATISTICS IN A FEEDBACK SEMICONDUCTOR LASER



PYRAGAS CHAOS CONTROL IN A ROSSLER SYSTEM



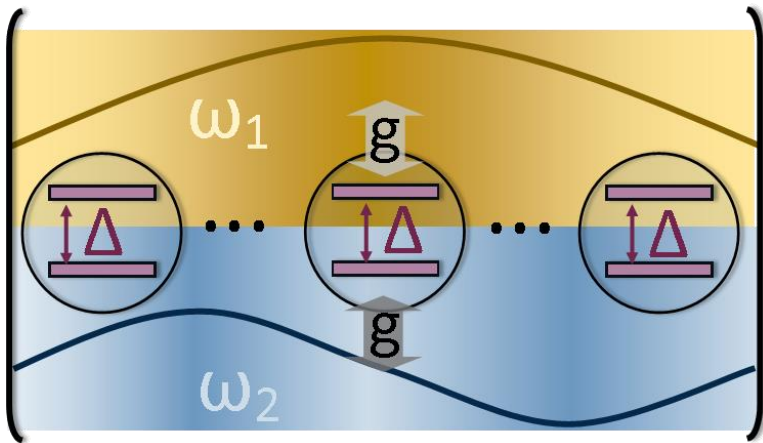
$$\text{Control} \sim (X(t - \tau) - X(t))$$

Non-invasive

Stabilizes fixed points, unstable limit cycles

From Tavis-Cummings to a Two-mode-Laser

$$H = \sum_{m=1}^2 \omega_m a_m^\dagger a_m + \frac{\Delta}{2} J^z + g \sum_{m=1}^2 (a_m J^+ + a_m^\dagger J^-)$$



ω_i : Frequencies of the cavity modes

Δ : Atomic Frequency

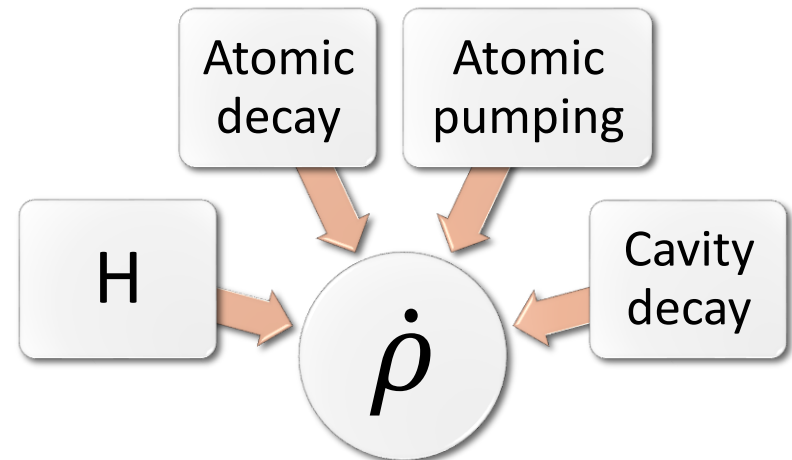
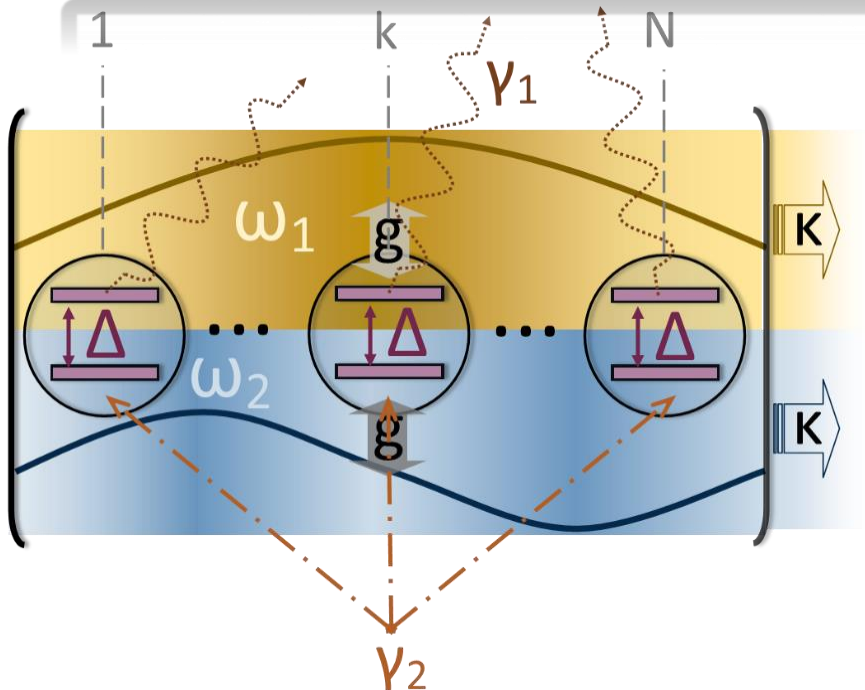
g : Coupling

$J^\pm = \sum_i \sigma_i^\pm$ collective spin

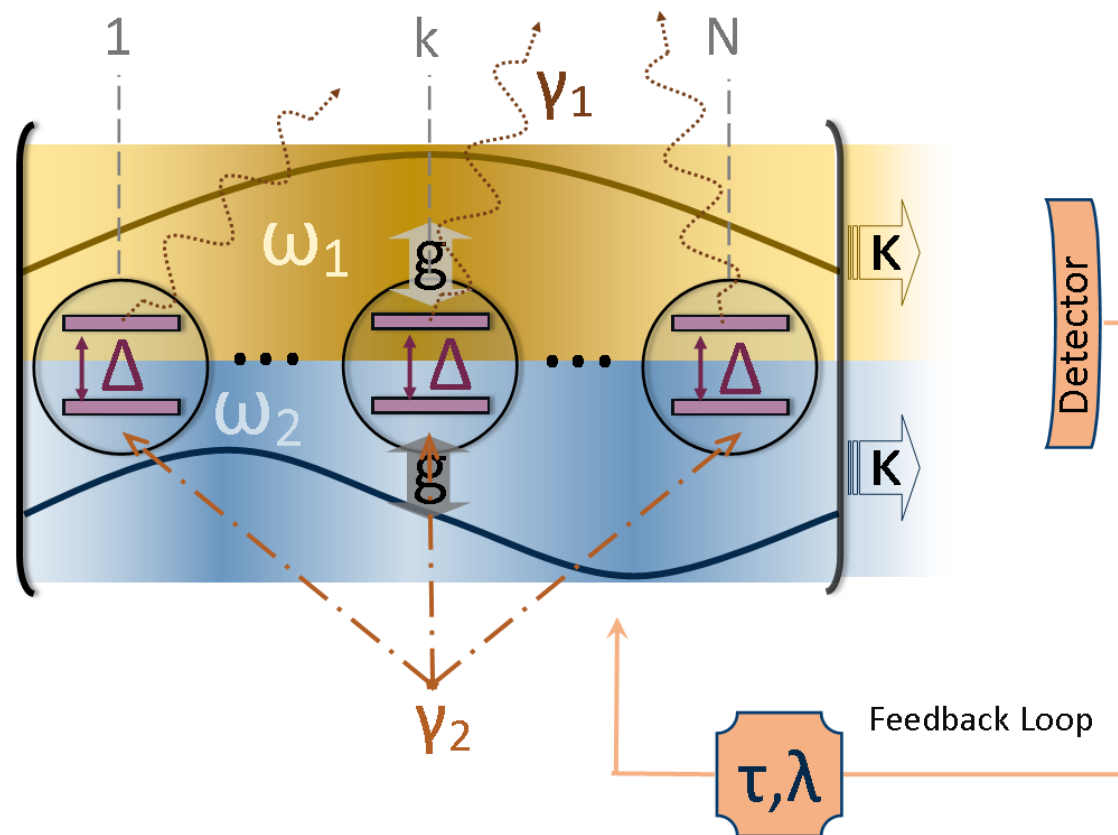
$J^z = \frac{1}{2} \sum_i \sigma_i^z$

From Tavis-Cummings to a Two-mode-Laser

$$\frac{d\hat{\rho}(t)}{dt} = -i[\hat{H}, \hat{\rho}] - \kappa\mathcal{L}[\hat{a}_1]\hat{\rho} - \kappa\mathcal{L}[\hat{a}_2]\hat{\rho} - \frac{\gamma_2}{2} \sum_{k=1}^N \mathcal{L}[\hat{\sigma}_k^+]\hat{\rho} - \frac{\gamma_1}{2} \sum_{k=1}^N \mathcal{L}[\hat{\sigma}_k^-]\hat{\rho},$$



5



Lindblad Master Eq.

- $\hat{\rho} = -i[\hat{H}, \hat{\rho}] - \frac{\kappa}{2} D[\hat{O}_k] \hat{\rho}$
- Lindbladian $D[\hat{O}_k] \hat{\rho} = \hat{O}_k^\dagger \hat{O}_k \hat{\rho} + \hat{\rho} \hat{O}_k^\dagger \hat{O}_k - 2\hat{O}_k \hat{\rho} \hat{O}_k^\dagger$
- \hat{O}_i - System operator

Mean-Field and Equations of Motion

- $\partial_t \langle \hat{O} \rangle = \text{Tr}(\hat{O} \hat{\rho})$
- Factorized operator $\langle \hat{O}_1 \hat{O}_2 \rangle \approx \langle \hat{O}_1 \rangle \langle \hat{O}_2 \rangle$, valid in thermodynamic limit
- Semiclassical equations: $\partial_t \vec{v}(t) = \vec{f}(\vec{v})$,
 $\vec{v} = (\langle \hat{O}_1 \rangle, \langle \hat{O}_2 \rangle \dots)$

Fixed Points

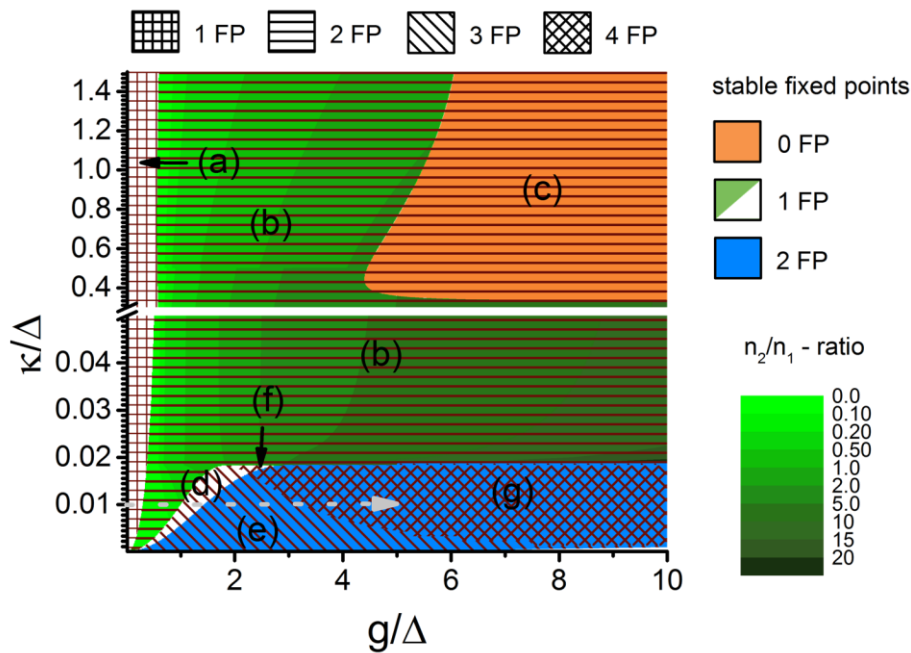
- $\vec{f}(\vec{v}^0) = 0 \Rightarrow \vec{v}^0 = (\langle \hat{O}_1^0 \rangle, \langle \hat{O}_2^0 \rangle, \dots)$

Introduce Pyragas Feedback into the Equations

Stability and Feedback

- Linear stability analysis: $\vec{v} = \vec{v}_0 + \delta \vec{v}$
- $\det(\Lambda \cdot \mathbf{1} - \mathbf{B} - (\mathbf{B} - \mathbf{A}) \cdot e^{-\Lambda \tau}) = 0$
- Numeric: Beyond linear stability

Stability Diagram



- Meanfield Eqs.

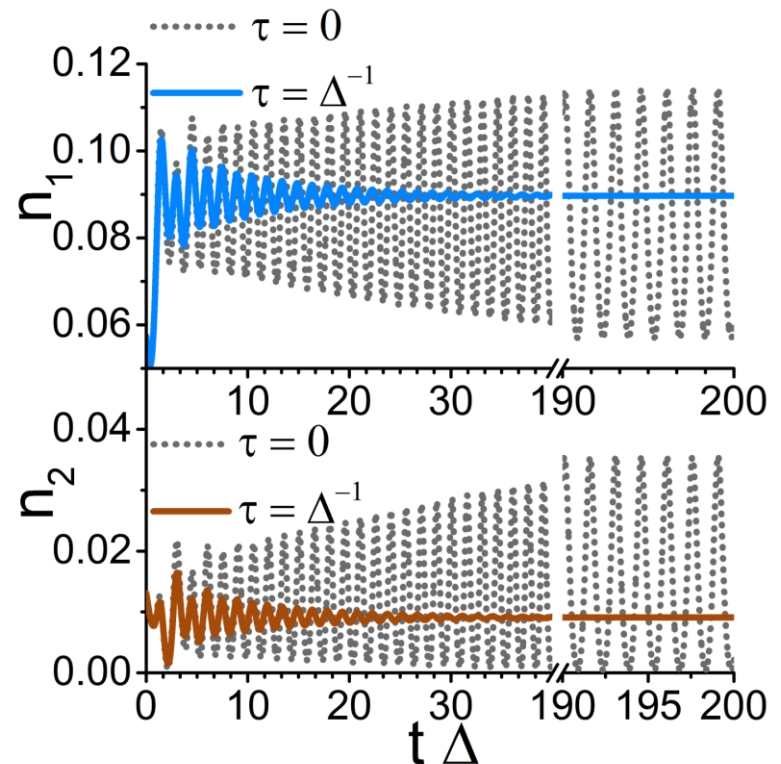
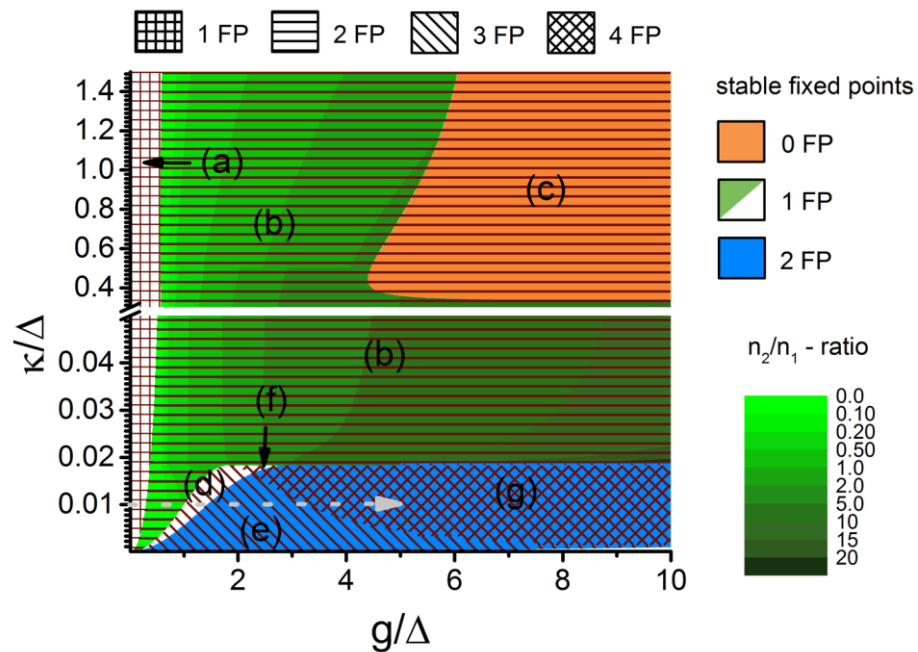
$$\partial_t \vec{v}(t) = \vec{f}(\vec{v})$$

$$\vec{v} = (\langle \hat{a}_1 \rangle, \langle \hat{a}_2 \rangle, \langle \hat{J}_+ \rangle, \langle \hat{J}_- \rangle, \langle \hat{J}_z \rangle)$$

- Steady state: $\partial_t \vec{v} = 0 \rightarrow$ exact solution

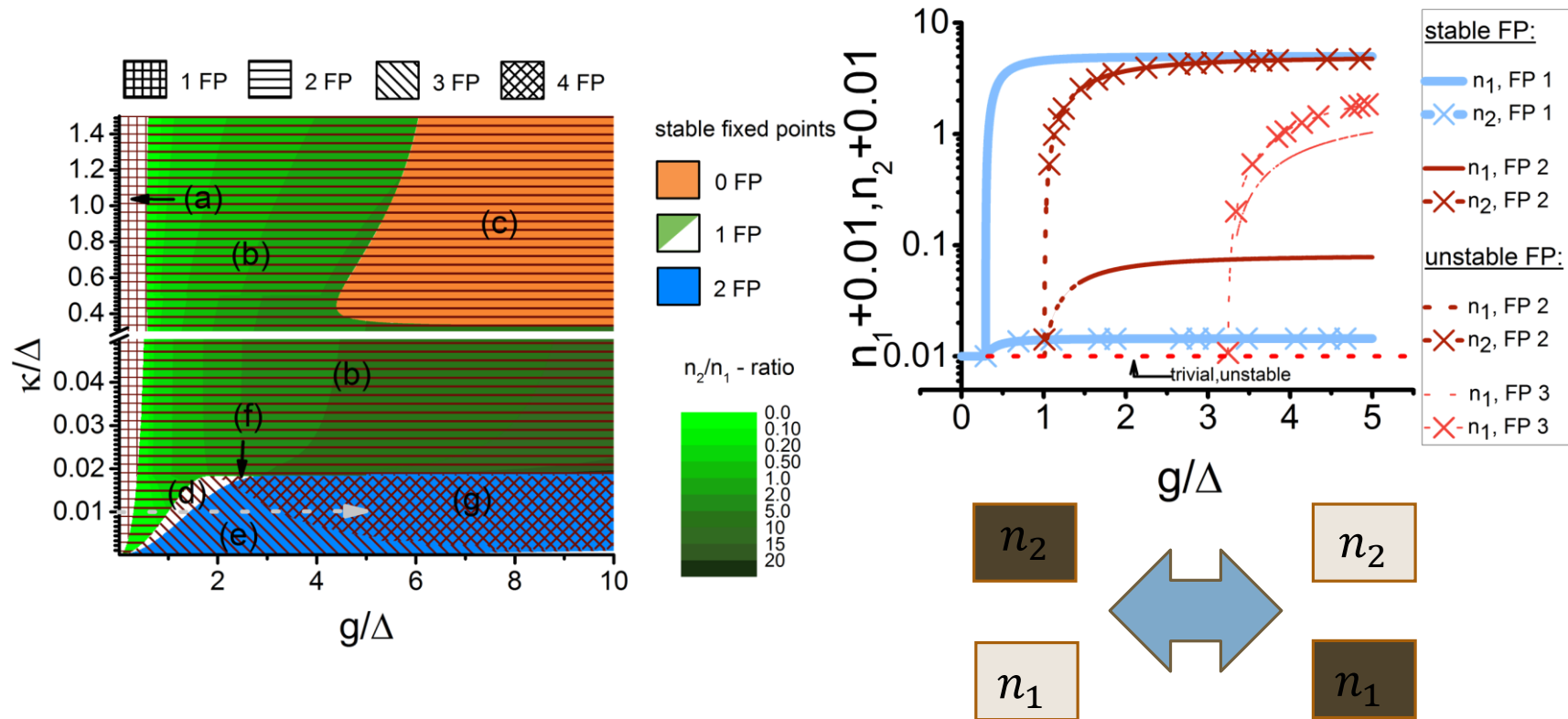
Area	(a)	(b)	(c)	(d)	(e)	(f)	(g)
Fixed Points	1	2	2	3	3	4	4
Stable	1	1	0	1	2	1	2

Pyragas Feedback: Stabilization (Area (c))



$$\dot{J}_Z = \cdots + \lambda(J_Z[t - \tau] - J_Z[t])$$

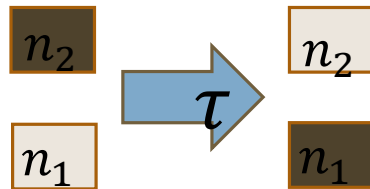
Pyragas Feedback: Selecting the Lasing Mode (Area (e))



Pyragas Feedback: Selecting the Lasing Mode (Area (e))

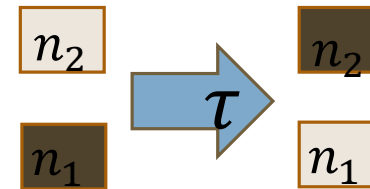
SWITCHING TO GROUND STATE
OCCUPATION

$$\omega_1 \rightarrow (\omega_1 + \lambda(n_2(t - \tau) - n_2(t)))$$



SWITCHING TO EXCITED STATE
OCCUPATION

$$\dot{a}_1 = \dots + \lambda(a_1(t - \tau) - a_1(t))$$



Summary&Outlook

TC-Laser

Various Pyragas Schemas

Outlook

Stabilization

Choice of the
lasing mode

Connection to
Thermalization
[*]

Control of
Thermalization

Beyond
Meanfield
Pyragas

W. Kopylov et al. PRA **92** (2015)

* P. Kirton and J. Keeling, PRL **111** (2013)