Primordial Models for Dissipative Bose-Einstein Condensates

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March 12, 2012

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Experiment

- 87 Rb-BEC of 10^5 particles at T = 80 nK
- Anisotropic harmonic trap $\Omega_{||} = 2\pi \cdot 170 \text{ Hz},$ $\Omega_{\perp} = 2\pi \cdot 13 \text{ Hz}$
- Focussed electron beam of Gaussian shape with width $w = \frac{FWHM}{\sqrt{8 \ln 2}} = 100 \text{ nm}$
- H. Ott et *al.*, Nature Phys. **4**, 949 (2008)



Complex potential

Gross-Pitaevskii-equation with complex potential

$$i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{r}, t) = \left[-\frac{\hbar^2}{2M} \Delta + V_{\mathrm{R}}(\mathbf{r}) + iV_{\mathrm{I}}(\mathbf{r}) + g |\Psi(\mathbf{r}, t)|^2 \right] \Psi(\mathbf{r}, t)$$

• Real part: harmonic trapping $V_{\rm R}(\mathbf{r}) = \frac{1}{2}M \left[\Omega_{\perp}^2(x^2 + y^2) + \Omega_{||}^2 z^2\right]$

• Imaginary part: Gaussian beam $V_{\rm I}({f r}) = - C \exp\left(-rac{x^2+y^2}{2w^2}
ight)$

•
$$C = \frac{\hbar}{2} \frac{\sigma_{\text{tot}}}{e} \frac{I}{2\pi w^2} = 1, 8 \cdot 10^{-30} \text{ J}, \quad \sigma_{\text{tot}} = 100 a_0, \quad I = 20 \text{ nA}$$

• Non-stationary probability density $\rho(x, t) = \exp \left[\frac{2}{\hbar}E_{\rm I}(t - t_0)\right]\rho(x, t_0)$

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Complex potential square well Quantization conditions

Complex potential square well

- Simplications for primordial model
- One-dimensional BEC, extension $L=R_{\perp}=$ 640 $\mu{
 m m}$
- No interaction g = 0 (Schrödinger equation)
- \bullet Square-well approximation of $V_{\rm R}$ and $V_{\rm I}$
- Separation into area 1, 2 and 3

$$V_{
m R}(x) = \left\{egin{array}{cc} 0 & , & |x| < L \ & \ \infty & , & |x| > L \end{array}
ight.$$



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Complex potential square well Quantization conditions

Quantization conditions

• Dimensionless variables (extension L normalized to $\frac{\pi}{2}$)

$$\varepsilon = \frac{E}{\frac{\hbar^2 \pi^2}{2M(2L)^2}}, \qquad c = \frac{C}{\frac{\hbar^2 \pi^2}{2M(2L)^2}}, \qquad \chi = \frac{\pi}{2} \frac{x}{L}, \qquad \omega = \frac{\pi}{2} \frac{w}{L}$$

- Experiment: $\omega = 0.0002$, c = 80000
- Quantization conditions for symmetric and antisymmetric states

$$0 = \sqrt{\varepsilon^{s}} \cot\left[\left(\omega - \frac{\pi}{2}\right)\sqrt{\varepsilon^{s}}\right] + \sqrt{\varepsilon^{s} + ic} \tan\left(\omega\sqrt{\varepsilon^{s} + ic}\right)$$
$$0 = \sqrt{\varepsilon^{a}} \cot\left[\left(\omega - \frac{\pi}{2}\right)\sqrt{\varepsilon^{a}}\right] - \sqrt{\varepsilon^{a} + ic} \cot\left(\omega\sqrt{\varepsilon^{a} + ic}\right)$$

•
$$\lim_{c \to 0} \varepsilon = \lim_{\omega \to 0} \varepsilon = m^2$$
; $\lim_{\omega \to \frac{\pi}{2}} \varepsilon = m^2 - ic$, $m \in \mathbb{N}$

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Energies Densities Limit of large dissipation

General solutions of quantization conditions

- ε_{I} non-positive (damping)
- Correct limit $\varepsilon(c=0) = m^2$
- $\lim_{c \to \infty} \varepsilon_{\mathrm{I}} = 0 \to k$ -states $\lim_{c \to \infty} \varepsilon_{\mathrm{I}} \stackrel{\text{linear}}{=} -\infty \to n$ -states
- *n*-states most damped
- Corresponds to $\omega \to 0$, $\omega \to \frac{\pi}{2}$
- $m \in \mathbb{N}$ describes limit c = 0, $n, k \in \mathbb{N}$ for $c \to \infty$



Figure: $\omega = 0.6$

• 2 adjoining k-states (different parity) fuse to one state (same k)

Energies Densities Limit of large dissipation

Limit of small and big waist



- $\omega \approx 0$: only k-states; antisymmetric states $\varepsilon^a \approx m^2$ for all c; ε^s tends to ε^a ; $\varepsilon^s_{\rm L}$ reveals deep minimum
- $\omega \approx \frac{\pi}{2}$: only *n*-states; all states $\varepsilon \approx m^2 ic$ for all *c*

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Energies Densities Limit of large dissipation

Interchange of states



Energies Densities Limit of large dissipation

Densities



- Lowest energy levels for $\omega=0.6$
- *n*-states tend to inside (area 2),
 k-states to outside (area 1, 3)
- Corresponds to $\omega = 0, \omega = \frac{\pi}{2}$ (areas vanish: states disappear)
- *k*-states: no central maximum
 → no symmetric states
- Three independent wells develop
- Area 1 and 3 equivalent

 → always 2 k-states fuse

Energies Densities Limit of large dissipation

Limit of large dissipation

n-states: tend to center, strongly damped
 k-states: tend to borders, less damped

- Ansatz: $\varepsilon_{\mathrm{I}}^k = 0$ for k-states and $\varepsilon_{\mathrm{I}}^n = -c$ for n-states for $c \to \infty$
- Limit for real part via quantization conditions (in original variables)

$$E_{\rm R}^n = \frac{\hbar^2 \pi^2}{2M(2w)^2} n^2 \ , \qquad E_{\rm R}^k = \frac{\hbar^2 \pi^2}{2M(L-w)^2} k^2$$

- Energies of real square well potentials with the width of each area
- Two k-states have same energy (area 1 and 3 have equal width)

Possible improvements

- Taking more dimensions, only two areas $(|\mathbf{r}| < w \text{ and } |\mathbf{r}| > w)$
- Considering interaction $g \neq 0$, derive Gross-Pitaevskii-Equation with complex potential via variational ansatz
- Harmonic trap $V_{\mathrm{R}} = \frac{1}{2} M \Omega_{\perp}^2 x^2$ instead of square well
- Next order Taylor approximation of $V_{\rm I}$, harmonic approximation

$$V_{\rm I}(x) = C\left(1 - rac{x^2}{w^2}
ight)$$

- \bullet Actually no upper border for ω
- Nearly same results for harmonic approximation



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Possible improvements

Possible assignments of numerical solutions



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Possible improvements

Non-hermitian dynamics

- Non-hermitian Hamilton operator $H = -\frac{\hbar^2}{2M} \frac{d^2}{dx^2} + V_R(x) + iV_I(x)$
- Complex energy eigenvalues $E=E_{\mathrm{R}}+iE_{\mathrm{I}}$
- Time evolution exhibits non-stationary probability density

$$\rho(x,t) = \exp\left[\frac{2}{\hbar}E_{\mathrm{I}}(t-t_0)\right]\rho(x,t_0)$$

• Continuity equation includes drain of probability since $V_{\mathrm{I}}(x) \leq 0$

$$rac{\partial}{\partial t}
ho(x,t)+rac{\partial}{\partial x}j(x,t)=rac{2}{\hbar}V_{\mathrm{I}}(x)
ho(x,t)$$