

Stability diagrams and dynamics of quantum degenerate Fermi gases of polar molecules

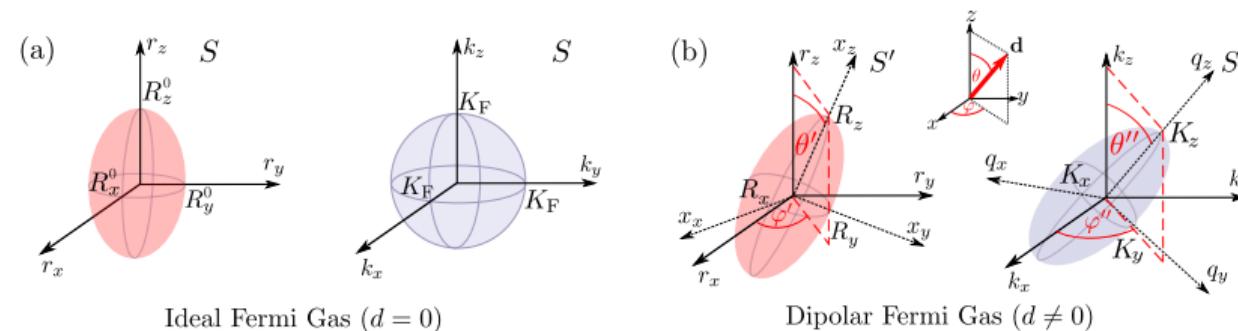
Vladimir Veljić¹, Axel Pelster², Antun Balaž¹

¹Scientific Computing Laboratory, Center for the Study of Complex Systems,
Institute of Physics Belgrade, University of Belgrade, Serbia

²Physics Department and Research Center OPTIMAS,
Technical University of Kaiserslautern, Germany

Motivation

- Observation of the Fermi surface (FS) deformation
F. Ferlaino group, Science **345**, 1484 (2014)
- Realization of the degenerate Fermi gas of polar molecules
J. Ye group, Science **363**, 853 (2019)



Veljić, Balaž, Pelster, Phys. Rev. A **95**, 053635 (2017)

Veljić et al., New J. Phys. **20**, 093016 (2018)

Wigner function approach

- Variational phase-space approach for the Wigner function

$$\nu(\mathbf{r}, \mathbf{k}) = \int d^3 r' e^{-i\mathbf{k} \cdot \mathbf{r}'} \rho \left(\mathbf{r} + \frac{1}{2}\mathbf{r}', \mathbf{r} - \frac{1}{2}\mathbf{r}' \right)$$

using the Hartree-Fock approximation

- Ansatz for the Wigner function at zero temperature:

$$\nu^0(\mathbf{r}, \mathbf{k}) = \Theta \left(1 - \sum_{i,j} x_i \mathbb{A}_{ij} x_j - \sum_{i,j} q_i \mathbb{B}_{ij} q_j \right)$$

$$\mathbb{A}' = \begin{pmatrix} 1/R_x'^2 & 0 & 0 \\ 0 & 1/R_y'^2 & 0 \\ 0 & 0 & 1/R_z'^2 \end{pmatrix} \quad \text{and} \quad \mathbb{B}'' = \begin{pmatrix} 1/K_x''^2 & 0 & 0 \\ 0 & 1/K_y''^2 & 0 \\ 0 & 0 & 1/K_z''^2 \end{pmatrix}$$

Energy and Thomas-Fermi radii and momenta

- Total energy of the system:

$$\begin{aligned} E_{\text{tot}} = & \frac{N}{8} \left(\sum_i \frac{\hbar^2 K_i^2}{2M} + \sum_{i,j} \frac{M \omega_i^2 \mathbb{R}_{ij}'^2 R_j^2}{2} \right) - \frac{6N^2 c_0}{R_x R_y R_z} \\ & \times \left[F_A \left(\frac{R_x}{R_z}, \frac{R_y}{R_z}, \theta, \varphi, \theta', \varphi' \right) - F_A \left(\frac{K_z}{K_x}, \frac{K_z}{K_y}, \theta, \varphi, \theta'', \varphi'' \right) \right] \end{aligned}$$

- Generalized anisotropy function:

$$F_A(x, y, \theta, \varphi, \tilde{\theta}, \tilde{\phi}) = \left(\sum_i \mathbb{R}_{iz} \tilde{\mathbb{R}}_{ix} \right)^2 f \left(\frac{y}{x}, \frac{1}{x} \right) + \left(\sum_i \mathbb{R}_{iz} \tilde{\mathbb{R}}_{iy} \right)^2 f \left(\frac{x}{y}, \frac{1}{y} \right) + \left(\sum_i \mathbb{R}_{iz} \tilde{\mathbb{R}}_{iz} \right)^2 f(x, y)$$

- Particle number conservation: $N = \frac{1}{48} R_x R_y R_z K_x K_y K_z$
- Fermi surface remains cylindrically symmetric ellipsoid: $K_x = K_y$

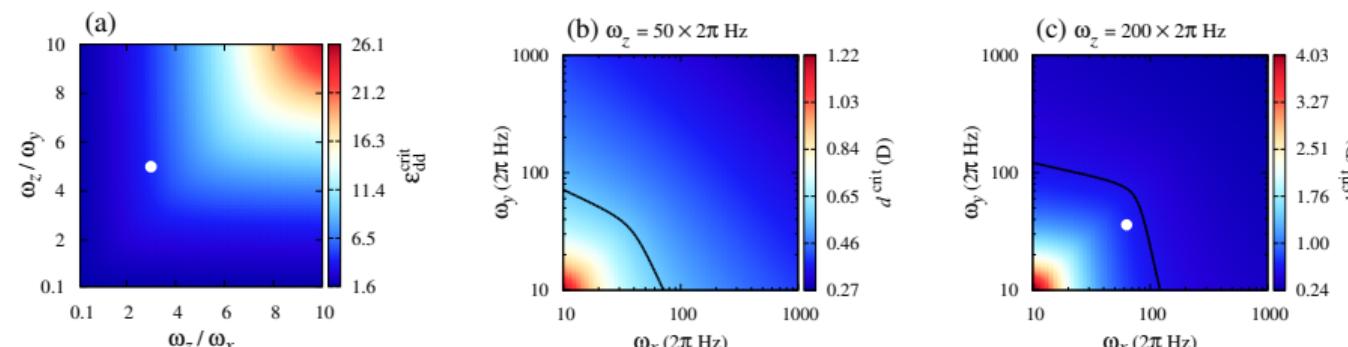
Veljić, Pelster, Balaž, arXiv:1902.09518

Stability diagrams

- Non-dimensional form of equations: species-independent!
- Relative DDI strength:

$$\varepsilon_{dd} = \frac{d^2}{4\pi\varepsilon_0} \sqrt{\frac{M^3}{h^5}} (\omega_x \omega_y \omega_z N)^{1/6}$$

- Stability only depends on the trap aspect ratios and the dipoles' orientation



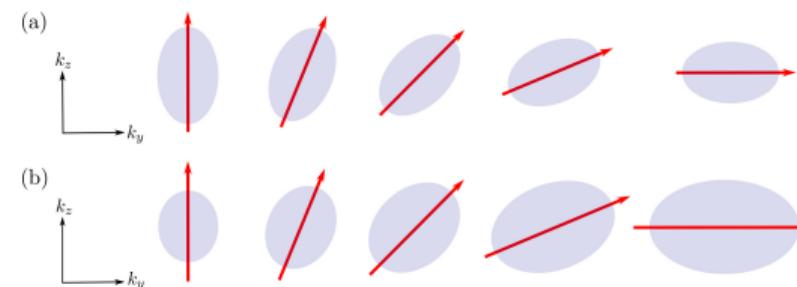
Veljić, Pelster, Balaž, arXiv:1902.09518

Fermi surface deformation

- Fermi surface deformation:

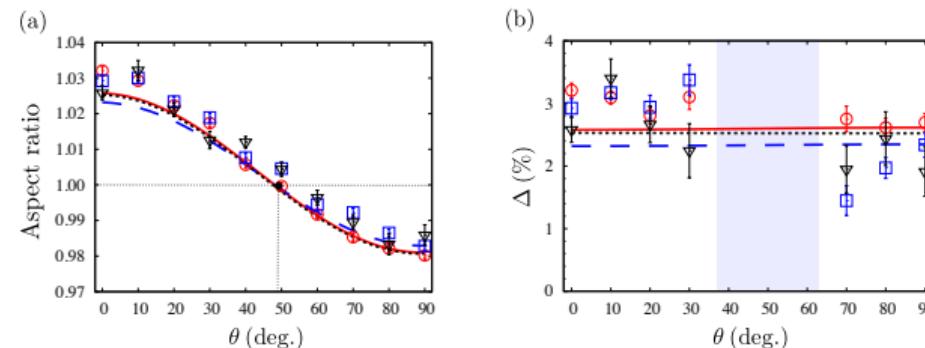
$$\Delta = \frac{K_z}{K_x} - 1$$

- Fermi surface orientation along the dipoles' direction: $\theta'' = \theta$, $\varphi'' = \varphi$
- For polar molecules effect of FS deformation is more dominant and complex
- Dipolar atoms \rightarrow rigid FS; Polar molecules \rightarrow soft FS



Fermi surface deformation - Er

- Comparison of theoretical results for A_K and Δ with the experimental results for $A_R(t = 12 \text{ ms})$ for ^{167}Er
- Ballistic expansion: $A_K(0) = \lim_{t \rightarrow \infty} A_R(t)$

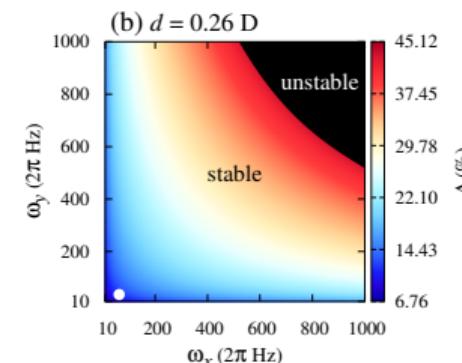
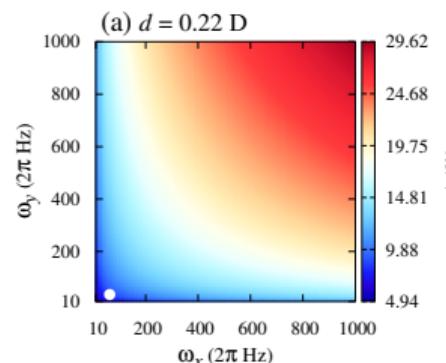


- $\varphi = 14^\circ$, $\gamma = 28^\circ$
- $N = 6.6 \cdot 10^4$, (579, 91, 611) Hz
- $N = 6.3 \cdot 10^4$, (428, 91, 459) Hz
- $N = 6.1 \cdot 10^4$, (408, 212, 349) Hz

Veljić et al., New J. Phys. **20**, 093016 (2018)

Fermi surface deformation - KRb

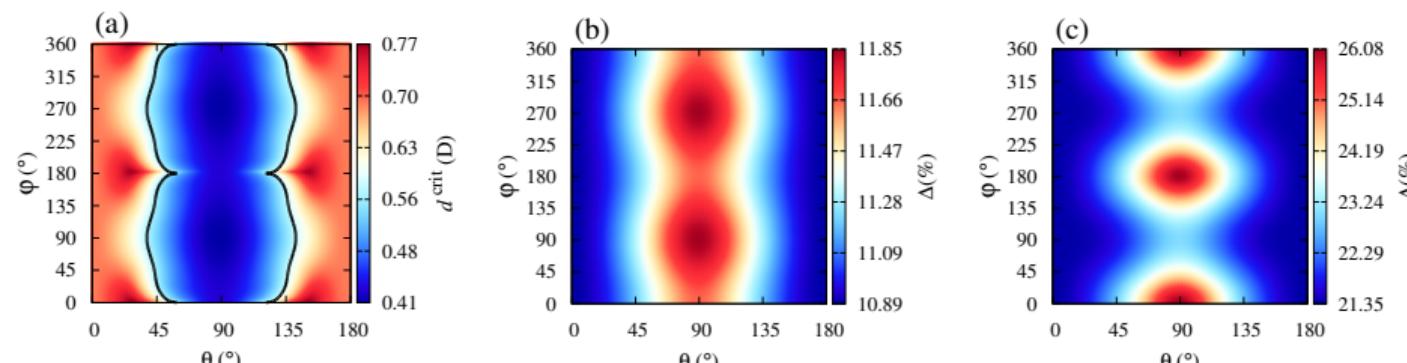
- FS deformation is much larger in gases of polar molecules
- $N = 3 \cdot 10^4$, $\omega_z = 2\pi \times 200$ Hz, $\theta = \varphi = 0$
- Even small changes in the dipolar moment strength can significantly affect the systems' stability



Veljić, Pelster, Balaž, arXiv:1902.09518

Fermi surface deformation - KRb

- FS deformation and its angular distribution can be tuned by changing the trap frequencies
- $N = 3 \cdot 10^4$, $2\pi \times (63, 36, 200)$ Hz; $2\pi \times (50, 500, 900)$ Hz



Veljić, Pelster, Balaž, arXiv:1902.09518

Conclusions

- General Hartree-Fock theory for the ground state of polarized trapped fermions with arbitrary \mathbf{d} at $T = 0$
- Species-independent, universal stability diagrams
- Molecular cloud shape and FS deformation strongly depend on the dipoles' orientation
- FS not only rigidly follows the dipoles' orientation, but changes its shape
- DDI has a significant role during the TOF
- Our new theory enables the study of the interplay between the FS deformation and superfluid pairing
- Outlook: 2D and 1D geometry