Bonds in motion – the time-dependent electron localization function



T. Burnus, M. A. L. Marques & E. K. U. Gross Freie Universität Berlin, Germany



# Classification of bonds

How can one give a rigorous mathematical meaning to chemical concepts such as

- single, double and triple bonds
- lone pairs

The electron density  $n_{\sigma}(\mathbf{r}, t)$  is not useful and the one-particle orbitals that stem from a density functional theory calculation are generally quite delocalised over several atoms, ambiguous and not very intuitive.

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# Definitions

$$D_{\sigma}(\boldsymbol{r},\boldsymbol{r}') := \sum_{\sigma_3,...,\sigma_N} \int \mathrm{d}^3 \boldsymbol{r}_3 \cdots \int \mathrm{d}^3 \boldsymbol{r}_N \left| \Psi(\boldsymbol{r}\sigma,\boldsymbol{r}'\sigma,\boldsymbol{r}_3\sigma_3,\ldots,\boldsymbol{r}_N\sigma_N,t) \right|^2$$

is the diagonal of the two-body density matrix and gives the probability of finding an electron at r and another with the same spin at r'.

$$P_{\sigma}(oldsymbol{r},oldsymbol{r}'):=rac{D_{\sigma}(oldsymbol{r},oldsymbol{r}')}{
ho_{\sigma}(oldsymbol{r})}$$

is the conditional probability of finding an electron with spin  $\sigma$  at r, knowing with certainty that there is an electron with the same spin at r'.

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#### Coordinate transformation

Since we are only interested in the probability near the reference point r, we make a spherical average, setting r' := r + s

$$p_{\sigma}(\boldsymbol{r},s) = \frac{1}{4\pi} \int_0^{\pi} \mathrm{d}\theta \,\sin\theta \int_0^{2\pi} \mathrm{d}\phi \, P_{\sigma}(\boldsymbol{r},\boldsymbol{r}+(s,\theta,\phi))$$

and Taylor expand  $p_{\sigma}(\boldsymbol{r},s)$ 

$$p_{\sigma}(\boldsymbol{r},s) = \underbrace{p_{\sigma}(\boldsymbol{r},0)}_{0} + \underbrace{\frac{\mathrm{d}p_{\sigma}(\boldsymbol{r},s)}{\mathrm{d}s}}_{0} + \frac{1}{3}C_{\sigma}(\boldsymbol{r}) \, \boldsymbol{s}^{2}.$$

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# The meaning of $C_{\sigma}(\boldsymbol{r})$

 $C_{\sigma}(\mathbf{r})$  is a measure of electron localization. It gives the probability of finding an electron with the same spin in vicinity of the reference electron.

- If  $C_{\sigma}(\mathbf{r})$  is small then this reference electron must be localized.
- $C_{\sigma}(\mathbf{r}) \geq 0$  (since it  $p_{\sigma}(\mathbf{r}, s)$  is a probability) and not bounded from above.

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### Definition of the ELF

The electron localization function is defined as [Becke, Edgecombe, *JCP* **92**, 5397 (1990)]

$$ext{ELF} = rac{1}{1 + \left( C_\sigma(oldsymbol{r}) / C_\sigma^{ ext{uni}}(oldsymbol{r}) 
ight)^2},$$

where  $C_{\sigma}^{\text{uni}}(\boldsymbol{r}) = \frac{3}{5}(6\pi)^{2/3} n_{\sigma}^{5/3}(\boldsymbol{r})$  is the kinetic energy density of the uniform electron gas. With this definition, the ELF is dimensionless and  $0 \leq \text{ELF} \leq 1$ .

high localization  $\Leftrightarrow C_{\sigma}(\mathbf{r}) \approx 0 \qquad \Leftrightarrow \text{ELF} \approx 1$ completely delocalized  $\Leftrightarrow C_{\sigma}(\mathbf{r}) \approx C_{\sigma}^{\text{uni}}(\mathbf{r}) \Leftrightarrow \text{ELF} \approx 1/2$ 

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#### Derivation of the ELF

When the wave function is a slater determinant one obtains

• in the static case

$$C_{\sigma}(oldsymbol{r}) = \sum_{i=1}^{N_{\sigma}} \left| 
abla arphi_{i\sigma}(oldsymbol{r}) 
ight|^2 - rac{1}{4} rac{\left( 
abla n_{\sigma}(oldsymbol{r}) 
ight)^2}{n_{\sigma}(oldsymbol{r})}$$

• in the time-dependent case

$$C_{\sigma}(\boldsymbol{r},t) = \sum_{i=1}^{N_{\sigma}} \left| \nabla \varphi_{i\sigma}(\boldsymbol{r},t) \right|^{2} - \frac{1}{4} \frac{\left( \nabla n_{\sigma}(\boldsymbol{r},t) \right)^{2}}{n_{\sigma}(\boldsymbol{r},t)} - \frac{\left( \boldsymbol{j}_{\sigma}(\boldsymbol{r},t) \right)^{2}}{n_{\sigma}(\boldsymbol{r},t)}$$

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# ELF Examples (static case)



[A. Savin, R. Nesper, S. Wengert and T. F. Fässler, Angew. Chemie Int. Ed. 36, 1808 (1997)]

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# ELF Examples (static case)



Scattering of a high energetic proton on ethene (ethylene)



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### **TD-ELF** Examples

#### Ethyne (actylene) in a strong laser field



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## **TD-ELF Examples**

#### Scattering of a high energetic proton on ethene (ethylene)



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