

Modern methods in experimental physics: site-directed spin labeling EPR

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discussed

Solution Set 1

Problem 1: On- and off-resonance pulses

On-resonance pulses

- The magnetization nutates on a cone around $\vec{B}_1 = \frac{\omega_1}{\gamma} \vec{e}_y$. The nutation frequency is $\omega_1/2\pi = 100$ kHz. With the pulse length of $\tau = 2.5 \mu\text{s}$ we obtain the flip angle $\beta_1 = \omega_1\tau = 90^\circ = 1.57$ rad.
- The resulting magnetization vector $\vec{M}(\tau)$ in the rotating frame after the $2.5 \mu\text{s}$ pulse is obtained by calculating the nutation of $\vec{M}(0)$ at $t=\tau$ using: $\vec{M}(\tau) = \mathbf{R}_y(\beta_1)\vec{M}(0)$ with the flip angle $\beta_1 = \omega_1\tau = \pi/2$ rad.

$$\mathbf{R}_y(\beta_1) = \begin{pmatrix} \cos \beta_1 & 0 & \sin \beta_1 \\ 0 & 1 & 0 \\ -\sin \beta_1 & 0 & \cos \beta_1 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ -1 & 0 & 0 \end{pmatrix} \quad (1)$$

giving $\vec{M}(\tau) = \mathbf{R}_y(\beta_1)\vec{M}(0) = (1, 0, 0)$

Off-resonance pulses

- The magnetization nutates on a cone around the effective field vector given by $\vec{B}_{\text{eff}} = \vec{B}_1 + \Delta\vec{B} = \frac{\omega_1}{\gamma} \vec{e}_y + \frac{\Omega}{\gamma} \vec{e}_z$. The nutation frequency is $\omega_{\text{eff}}/2\pi = \sqrt{\omega_1^2 + \Omega^2}/2\pi = 141.4$ kHz. With the pulse length of $\tau = 2.5 \mu\text{s}$ we then obtain the effective flip angle $\beta_{\text{eff}} = \omega_{\text{eff}}\tau = 127^\circ = 2.22$ rad.

The angle between the effective magnetic field and the z-axis is $\theta = \arctan(|\omega_1/\Omega|) = \pi/4 = 45^\circ$.

For the on-resonance case, the nutation frequency is $\omega_1/2\pi = 100.0$ kHz. With the pulse length of $\tau = 2.5 \mu\text{s}$ we would obtain a $\pi/2$ pulse, characterized by a flip angle $\beta = \omega_1\tau = 90^\circ = 1.57$ rad.

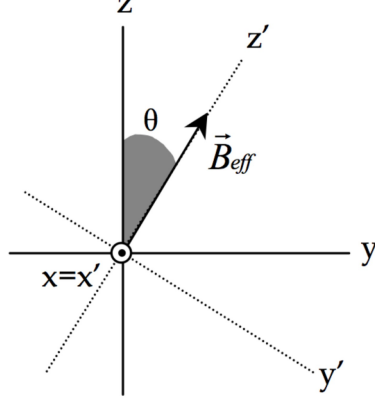


Figure 1: Orientation of the effective magnetic field in the two coordinate frames.

- b. For the off-resonance case, to make the calculations easier we study the precession of the magnetization vector in a tilted frame in such a way that its z-axis is defined by the effective field (see figure 2). The rotation matrix transforming vectors to this new coordinate system is

$$\mathbf{R}_x(\theta) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta & -\sin \theta \\ 0 & \sin \theta & \cos \theta \end{pmatrix} \quad (2)$$

The thermal equilibrium magnetization $\vec{M}(0)$ in this tilted frame is

$$\vec{M}'(0) = \mathbf{R}_x(\theta)\vec{M}(0) = (0, -\sin \theta, \cos \theta). \quad (3)$$

The matrix expressing the nutation around the effective field is

$$\mathbf{R}_z(\beta_{\text{eff}}) = \begin{pmatrix} \cos \beta_{\text{eff}} & -\sin \beta_{\text{eff}} & 0 \\ \sin \beta_{\text{eff}} & \cos \beta_{\text{eff}} & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (4)$$

The magnetization (in the tilted frame) after the pulse is

$$\vec{M}'(\tau) = \mathbf{R}_z(\beta_{\text{eff}})\vec{M}'(0) = (\sin \theta \sin \beta_{\text{eff}}, -\sin \theta \cos \beta_{\text{eff}}, \cos \theta). \quad (5)$$

and in the rotating frame (xyz)

$$\begin{aligned} \vec{M}(\tau) &= \mathbf{R}_x(-\theta)\vec{M}'(\tau) \\ &= (\sin \theta \sin \beta_{\text{eff}}, -\sin \theta \cos \beta_{\text{eff}} \cos \theta + \sin \theta \cos \theta, \sin^2 \theta \cos \beta_{\text{eff}} + \cos^2 \theta). \end{aligned} \quad (6)$$

- c. During the calculation we considered that the norm of the equilibrium magnetization $|\vec{M}(0)|$ is 1, thus the angle α between $\vec{M}(\tau)$ and $\vec{M}(0)$ is simply

$$\alpha = \arccos(\vec{M}(\tau)\vec{M}(0)) = \arccos(\sin^2 \theta \cos \beta_{\text{eff}} + \cos^2 \theta) \quad (7)$$

We finally get $\alpha = 78.6^\circ$.